

霜田治朗

相対論的現象で探る宇宙の進化V (2025/02/24-27) 部分的Collaborators: Y. Ohira, K. Asano, S. Inutsuka



Supernova Remnants:e.g., Tycho's SNR (SN 1572, Ia)



Supernova Remnants: Temporal Evolution

~ Late Sedov phase

SNR DEM L71 @LMC

(X-ray), age~a few kyr

(w/ clear Reverse shock)

G1.9+0.3: Youngest SNR in our Galaxy (age~140 yr, Bamba & Williams 22) Free Expansion Phase



~ Snowplow phase (T<0.1 keV or Vsh < 300 km/s)

SNR Cygnus Loop (UV), age~10 kyr

G70.0-21.5

Blue OIII, Red H α

age ~ 10-100 kyr?



Supernova Remnants: Emission Mechanism Examples



特に高エネルギー宇宙物学的 なものが抜粋されている (X-ray ~ Gamma-ray)。

Supernova Remnants: Emission Mechanism Examples







被加速粒子の量は予言しない. →独立な検証が必要だが、特に理論的な理 解はほとんど進んでいない.







Diffusive Shock Acceleration (DSA) →衝撃波を往復して粒子が加速していく

→Vshが小さくなると、衝撃波が拡散する 宇宙線に追いつけない。

→宇宙線が上流に逃走して加速終了」加速 により到達可能な最高エネルギーを決める。

「いつ、どれだけ逃走しているのか?」は 分かっていない。















~10 kpc ④ 銀河系からの逃走

the analogy to stellar winds, but in particular the CR component of the Interstellar Medium (ISM) (s does indeed secularly escape from the Galaxy to space, leading to a picture schematically given in Fi

At least those CR particles which exist in the ne of the solar system are observed to be well scattered i.e. to be efficiently coupled to the interstellar gas a field. The gaseous halo above a few kpc is probabl has a higher degree of ionization than regions of 1 thus coupling should be even better up there, ex in the isolated HVC's. Since the CR pressure is of the gas pressure, it does not only act on the solar neighbourhood, but it must have in addition dynamical effect on the overall gas and field confi Parker 1968, 1969 for reviews of the early thin matter). Because CRs do not cool radiatively, w their pressure gradient to act on the gas all along



加速現場からの逃走

~10 kpc ← 毎 ~ 10 kpc ← 毎 ~ 6 の 逃走

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Roles of Galactic Cosmic Rays

Evolution of the Protoplanetary disk

CR ionization rate controls the size/life-time of protoplanetary disk via the angular momentum transportation by B-field. \rightarrow Planet formation

FR [Msun/yr]

Lookback time [Gyr]

The long-term SFR is regulated by the galactic wind driven by CRs.

 \rightarrow Galaxy formation

"Puzzling" Star Formation History (the metal amount)

@ disk

SFR ~ 3 Mo/yr

Gas mass ~ 10^9 Mo (Metallicity Zo ~ $0.01 \rightarrow$ Metal mass ~ 10^7 Mo) Salpeter IMF \rightarrow Massive Star FR ~ 0.1 Mo/yr

Total Metal Mass Ejected by SNe

 $\rightarrow \sim$ (SFR) x (Massive Star fraction) x (CO core mass fraction) x (14 Gyr) ~ (3 Mo/yr) x (0.1) x (3 Mo/8 Mo) x (14 Gyr) ~ 1.6 x 10⁹ Mo

~99 % of metals should be removed from the disk!

- $^{2} \rightarrow$ Persistent Outflow is required!
 - (SJ & Inutsuka 22, SJ, Inutsuka, & Nagashima 24, SJ & Asano 24)

Fermi Bubble & eROSITA Bubble

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Breitschwerdt+91

Outflow: $T \sim 0.1$ keV (~virial temp. of the MW) \rightarrow eROSITA bubble is consistent with this expectation. Should be launched at the disk (removing the metals) \rightarrow Formation mechanism is unclear...

Galactic Wind Scenario (SJ&Asano 24)

Assumption: ~10 % of SNe energy is consumed for launching the wind. This assumption should be tested. We consider a possibility: *ISM heated by CRs around the CR sources*.

*Hadronic γ -ray scenario $p_{cr}+p_{gas} \rightarrow 2\gamma$, ν

逃走宇宙線との闘争

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ISM中の輸送・伝搬

~10 pc-1 kpc

■ 偉い先行研究 「拡散するん ちゃう?知ら んけど、」 ~10 kpc ← 毎 銀河系からの逃走

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JWST bubbles: Hints for 10 pc-1 kpc physics/phenomena?

Implications: Star Formation

Inutsuka+15

- ・天の川銀河の体積
 Vg~(0.3 kpc)x π (10 kpc)²~10² kpc³
- ・1つのシェルが占める体積
 Vs ~ (0.1 kpc)^{^3}~10⁻³ kpc³
- ・天の川銀河全体を穴だらけにするのに必 要なSNの個数

Nsn ~ Vg/Vs ~ 10^5

- ・シェルが~0.1 kpc拡がるのにかかる時間
 Ts ~ (0.1 kpc)/(10 km/s) ~ 10 Myr
- ・必要なMSの形成率

MSFR ~ Nsn/Ts x Ms

~ 0.1 Mo/yr(Ms/10Mo)

・IMFからMSのmass fractionは~0.1 SFR ~ MSFR/0.1 ~ 1 Mo/yr

Implications: Star Formation

MWの星形成率(Haywood+14)

MWで提案されている星形成シナリオは,最近の観 測と(1=10くらいの気持ちで)整合的. *Bubble-Filament Paradigm* by Inutsuka

~10 kpc 銀河系からの逃走

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Old SNRs: Hints for 10 pc-1 kpc physics/phenomena?

Red: radio synchrotron Blue: [OIII]λ5007

推定半径>10 pcの*Old SNR*で, 最 近見つかった**[OIII]***λ***5007輝線.**

電波シンクロトロンのシェル (shock面)よりも*遠く*で明るい. →なんで??? 著者のFesen et al. 2024は, CRの関 与について言及しているけれども, 具体的なシナリオ・物理過程は不明. 要するに,単純な説明がない.

1st ionization potentials: O^{+1} →35 eV O^{+2} →55 eV O^{+3} →77 eV

Old SNRs: Hints for 10 pc-1 kpc physics/phenomena?

O+3→77 eV

逃走宇宙線との闘争

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Lookback linte 'Gyr

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- ・SNRから離れたところでISMが加熱されるなどが起こりうる.
- CR加熱率は色々提案されている(e.g., Acterberg+81, Zweibel 20, Yokoyama & Ohira 23)が、実際のところよく分かっていない(microphysicsがよく分からん).
- ・CRの有無で、ガスの運動や熱力学的状態がどう変わるかも整備されていない(流体力学).

The CR-hydrodynamics

$$\begin{aligned} \frac{\partial \rho}{\partial t} + \boldsymbol{\nabla} \left(\rho \boldsymbol{v} \right) &= 0\\ \rho \frac{d \boldsymbol{v}}{d t} &= - \boldsymbol{\nabla} \left(P_{\rm g} + P_{\rm cr} \right) \end{aligned}$$

 $P_{\rm g}$ is the pressure of thermal gas $P_{\rm cr}$ is the CR pressure

The energy equation is ...

$$dQ = d(E_{g} + E_{cr}) + (P_{g} + P_{cr})dV$$
$$dQ_{rad} + dQ_{conv} + dQ_{vis} + dQ_{cr} + \dots$$

The 1st law of thermodynamics should include the CRs

Radiation, (thermal) convection, viscosity, CRs energy interactions, …

$$P_{\rm g} = (\gamma_{\rm g} - 1)e_{\rm g} = 2e_{\rm g}/3$$

 $P_{\rm cr} = (\gamma_{\rm cr} - 1)e_{\rm cr} = e_{\rm cr}/3$

$$\begin{aligned} \frac{dP_{g}}{dt} - C_{g}^{2} \frac{d\rho}{dt} &= G(\rho, T, P_{cr}) \\ \frac{dP_{cr}}{dt} - C_{cr}^{2} \frac{d\rho}{dt} &= R(\rho, T, P_{cr}) \\ \frac{dP_{cr}}{dt} - C_{cr}^{2} \frac{d\rho}{dt} &= R(\rho, T, P_{cr}) \\ C_{cr} &= \sqrt{\frac{\gamma_{cr}P_{cr}}{\rho}} \\ R(\rho, T, P_{cr}) &= \nabla \cdot (D\nabla P_{cr}) - \gamma_{*}\mathcal{H}, \\ \gamma_{*} &= \frac{\gamma_{cr} - 1}{\gamma_{g} - 1} = \frac{1}{2}, \end{aligned}$$

$$\begin{aligned} C_{g} &= \sqrt{\frac{\gamma_{g}P_{g}}{\rho}}, \\ G(\rho, T, P_{cr}) &= \mathcal{L}^{*}(\rho, T) + \nabla \cdot (\mathcal{K}^{*}\nabla T) + \mathcal{H}, \\ \mathcal{L}^{*}(\rho, T) &= (\gamma_{g} - 1)\mathcal{L}_{rad,g}(\rho, T), \\ \mathcal{H}(\rho, T, P_{cr}) &= (\gamma_{g} - 1)\mathcal{L}_{een,cr}(\rho, T, P_{cr}), \end{aligned}$$

$$\begin{aligned} Energy Exchange effect \\ \frac{d(P_{g} + P_{cr})}{dt} - (C_{g}^{2} + C_{cr}^{2})\frac{d\rho}{dt} \sim (1 - \gamma_{*})\mathcal{H} \\ \frac{d(P_{g} + P_{cr})}{dt} - (C_{g}^{2} + C_{cr}^{2})\frac{d\rho}{dt} \sim (1 - \gamma_{*})\mathcal{H} \end{aligned}$$

 $+ \mathcal{H},$

$$P_{
m g} = (\gamma_{
m g} - 1)e_{
m g} = 2e_{
m g}/3$$
 The CR heating produces
 $P_{
m cr} = (\gamma_{
m cr} - 1)e_{
m cr} = e_{
m cr}/3$ additional total pressure!

$$\begin{aligned} \mathcal{P} &= P_{\rm g} + P_{\rm cr} &= \mathcal{P}(e_{\rm g} + |e_1|, e_{\rm cr} - |e_1|) \\ &= \mathcal{P}(e_{\rm g}, e_{\rm cr}) + \frac{|e_1|}{3} \end{aligned}$$

What happens??? →Test by Simple Numerical Model

The CR-hydrodynamics: Simple Numerical Model

1D spherical Uniform medium Put *CR bomb*

CR transport equation

$$\partial_t \mathcal{N} + \nabla \cdot (v \mathcal{N} - D_\gamma \nabla \mathcal{N}) - \partial_\gamma \left(\frac{\nabla \cdot v}{3} \gamma \mathcal{N} \right) = -|\mathcal{V}_A \partial_r \mathcal{N}|$$

CR Heating



CR initial condition

$$\mathcal{N} \propto \frac{\gamma^{-2.3}}{(r/R_s)^2} \exp \left[-\left(\frac{r}{R_s}-1\right)\left(\frac{\gamma}{\gamma_s}\right)^{-\epsilon}\right]$$

CR total energy = 10^{50} erg

The CR-hydrodynamics: Simple Numerical Model



The CR-hydrodynamics: Simple Numerical Model



$$\mathcal{P} = P_{g} + P_{cr} = \mathcal{P}(e_{g} + |e_{1}|, e_{cr} - |e_{1}|)$$
$$= \mathcal{P}(e_{g}, e_{cr}) + \frac{|e_{1}|}{3}$$

CR heating \rightarrow Increasing total pressure \rightarrow prefer to form hot & tenuous gas

Cooling rate decreases

Old SNRs: Hints for 10 pc-1 kpc physics/phenomena?



O+3→77 eV





ISM中の輸送・伝搬

~10 pc-1 kpc JWST - 7.7µm 1kpc Lookback linte 'Gyr

~10 kpc ← 毎 銀河系からの逃走

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10⁻¹⁵

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Microquasars

Ruoyu-Liu氏のスライドより



V4641 Sgr





HMXB & LMXB distribution (Yue+arxiv2412.13889)



Figure 1. Projection of the sample pulsars on to the Galactic plane. High-mass X-ray binaries (HMXB) and low-mass X-ray binaries are marked as red dots and blue pluses, respectively. The GC is the origin of the coordinate system and the coordinates of the Sun are (-8.5, 0.0). Other lines and symbols are described in the text.



Fermi & eROSITA Bubbles

The Galactic Evolution Scenario



Total mass of DM: ~ $10^{12} M_{sun}$ Total mass of stars: ~ 4-6 x $10^{10} M_{sun}$ Current SFR: ~ $1 M_{sun}/yr$ Total gas mass: ~ $10^9 M_{sun}$

Cf. Bland-Hawthorn & Gerhard 16, the Planck Collaboration 18

Ø From the current MW ...

- 1. The gas should be depleted within $\sim 1 \text{ Gyr}!$
- 2. Replenishment of gas is required.
- 3. Galactic halo (CGM) may be a dominant gas reservoir.

The Galactic Evolution Scenario



Z=0

@ disk

SFR ~ 3 Mo/yr

Gas mass ~ 10^9 Mo (Metallicity Zo ~ $0.01 \rightarrow$ Metal mass ~ 10^7 Mo) Salpeter IMF \rightarrow Massive Star Formation Rate ~ 0.1 Mo/yr

Total Metal Mass Ejected by SNe over Cosmic age

 \rightarrow ~ (SFR) x (Massive Star fraction) x (CO core mass fraction) x (14 Gyr) ~ (3 Mo/yr) x (0.1) x (3 Mo/8 Mo) x (14 Gyr) ~ 1.6 x 10⁹ Mo

~99 % of metals should be removed from the disk! → Persistent Outflow is required! (see, Shimoda, Inutsuka, & Nagashima 2024 for details)

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(SJ & Inutsuka 22, SJ, Inutsuka, & Nagashima

24, SJ & Asano 24)



Outflow Model (SJ & Inutsuka 2022) w/ Radiative cooling & CR diffusion

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and the ras to be fully ionized (by disk OB stars from below, and globular Gueters, et as well as QSO's from above) that diffusion becomes negligible compared to Alfvénic drift. The combined pressure gradients of gas, CRs and waves should then gradually accelerate the gas away from the disk with mass velocity u (Fig. 2).

The average field direction is also influenced by CR effects even in the lower disk (|z| < 1 kpc): the Parker instability (Parker, 1966) of an assumed equilibrium between the confining gravitational field and the disruptive CR and magnetic field pressures will lead to an inflation of field lines, which might allow CRs plus gas to escape into closed magnetic bubbles. This is topologically different from open field lines which extend to $|z| = \infty$, but nevertheless adds to any galactic mass loss.

Thus in detail we expect a complicated configuration of disk and halo. On average, however, and on spatial scales of

Essen

CR

(Shapiro & Field 76)

Radiative cooling $\rightarrow T < T_vir \rightarrow$ wind never launchin Heating by CRs \rightarrow Comparable with Radiative coolin

> CRs scattered by δB \rightarrow Momentum transferred to δB $\rightarrow \delta B$ grows \rightarrow dissipation of δB \rightarrow Thermal gas heated

 $\Phi_{H^{(x)}} = \frac{GM_{H0}}{R_b} [\Gamma = |V_A \nabla P_{cr}| (erg/cc/s)$ (e.g., Kulsrud 2005)

The existence vation all **Brend** schwerdt+91 extended the pioneering work by Ipavich 75, the Galaxy there of unservorses timated that the mass loss rate due to the wind is ~1

corresponding/tyr.

parameters are a_i

 $\mathcal{M}_i = (2.05 \times 10^{10}$

where $M_{H0} = 1$.

 $\sqrt{R_0^2+z^2}, x \equiv R/$

The halo poter

respectively.

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$$\frac{n^2 \Lambda}{Q_{\rm w}} \simeq 0.91 \left(\frac{n}{10^{-3} \,{\rm cm}^{-3}}\right)^{5/2} \left(\frac{B}{1 \,\mu{\rm G}}\right)^{-1} \left(\frac{P_{\rm cr}}{0.3 \,{\rm eV} \,{\rm cm}^{-3}}\right)^{-1} \times \left(\frac{H_{\rm cr}}{10 \,{\rm kpc}}\right) \left(\frac{\Lambda}{10^{-22} \,{\rm erg} \,{\rm cm}^3 \,{\rm s}^{-1}}\right).$$
(20)

The reasonable physical parameters result in the comparable heating rate!

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where $M_{H0} = 1$.

 $\sqrt{R_0^2+z^2}, x \equiv R/$

The halo poter





Consistent with both X-ray observations and theoretical model o (SJ & Inutsuka 2022)



The baryon accretion rate

 \rightarrow Use the results of the DM N-body simulation

(Rodriguez-Puebla+16)



*Metals are removed from the disk

What is essence?





Problems

1. The origin of the X-ray emissions at/around the Galactic disk.

→Numerical simulations imply its existence, but actual physical processes are still under debated.

2. The existence of the Cosmological accretion gas onto the disk.

 \rightarrow If it has a virial temperature of ~ 10⁶ K, the observations are difficult (FUV ~ soft X-rays are obscured).

3. The resultant disk & CGM conditions (including CRs, metal, etc.) are still unclear.

We introduce one of the consequences of the Galactic wind (SJ & Asano 2024)

Fermi Bubble & eROSITA Bubble

 Predehl+20

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Breitschwerdt+91

Using observational data of current MW, we will check whether our scenario is consistent with the observed bubbles.

*Hadronic γ -ray scenario $p_{gas} \rightarrow 2\gamma$, ν



Fig. 3. The gas maps of the Galaxy (*left panels*) in direct comparison with the fitted model (*right panels*). The upper left panel shows the molecular hydrogen distribution (Dame et al. 2001), while the atomic hydrogen distribution (Kalberla et al. 2005) is presented in the bottom left panel. Both maps are given for the Galactic latitude range of -30 to 30 degrees.

Parameter	Units	Value
$\rho_{\rm H_2}(0,0)$	cm ⁻³	4.06
$h_{\rm H_2}$	kpc	2.57
Z_{H_2}	kpc	0.08
$\hat{\rho_{\mathrm{HI}}(0,0)}$	cm^{-3}	0.32
$h_{ m HI}$	kpc	18.24
$z_{ m HI}$	kpc	0.52
R_t	kpc	2.75

Ь

$$D_{H_2}(R, z) = \rho_{H_2}(0, 0) \exp\left(-\frac{R}{h_{H_2}} - \frac{|z|}{z_{H_2}}\right),$$
$$D_{HI}(R, z) = \begin{cases} \rho_{HI}(0, 0) \exp\left(-\frac{R}{h_{HI}} - \frac{|z|}{z_{HI}}\right), & \rho > R_t\\ 0, & \rho < R_t \end{cases}$$



Gas distribution @the disk

Total Mass at R < 30 kpc $M_{\rm HI} \simeq 3.9 \times 10^9 \ M_{\odot}$ $M_{\rm H_2} \simeq 1.3 \times 10^9 \ M_{\odot}$ Enclosed SFR (M_{\Box} yr⁻¹)



Star Formation Rate

$$\begin{split} \dot{\Sigma}_{\rm sf} &= \frac{\epsilon_{\rm sf}}{\tau_{\rm sf}} \Sigma_{\rm H_2}, \\ \tau_{\rm sf}/\epsilon_{\rm sf} &= 0.5 \; \rm Gyr \end{split}$$

Total Star Formation Rate: 2.7 Mo/yr





$$p_{cr} + p_{gas} \rightarrow 2\gamma$$
, v

Out boundary = 30 kpc

$$\frac{d\boldsymbol{v}_{\mathrm{w}}}{dt} = -\frac{\boldsymbol{\nabla}P_{\mathrm{cr}}}{\rho_{\mathrm{w}}} - \boldsymbol{g}$$

Wind parcels are prepared at each R and $z = z_h$.

Their trajectories are computed like test

particles. the analogy to stellar winds, but in particular the CR component of the Interstellar Medium (ISM) (s does indeed secularly escape from the Galaxy to Λ space, leading to a picture schematically given in Fi

 \succ

At least those CR particles which exist in the ne Hot gasfldyesolar system are observed to be well scattered

i.e. to be efficiently coupled to the interstellar gas $z_h = 2$ kpeed. The gaseous halo above a few kpc is probabl has a higher degree of ionization than regions of 1 midplane coupling should be even bett Breitschwerdt#91

Model Computation

$$\overset{@V_{\rm cr}(r,\gamma)}{@P} = \dot{N}_{\rm cr,s}(r,\gamma) + D_{\rm cr}(\gamma)r^2 N_{\rm cr}(r,\gamma)$$

Steady-state

$$N_{cr}(r,\gamma) \stackrel{?}{=} \frac{\langle \overline{\varphi}(\gamma) \rangle}{4 \stackrel{?}{=} H^2} \stackrel{\checkmark}{=} \frac{N_{cr,s}(r^0,\gamma)}{|r-r^0|} d^3 r^4$$
the $\langle \overline{\varphi} \rangle = 1 \text{ Myr} \stackrel{\checkmark}{=} \frac{\gamma}{2} \stackrel{\forall \overline{U}_{cr} = 0.6}{(D_{cr} \approx 10^{28} \text{ cm}^2 \text{ s}^{-1}(\gamma/2)^{0.6})}$
(5)
(5) $V_{cr,s} = N_0 \stackrel{\checkmark}{=} \frac{\gamma}{2}$

*10% of SN energy→CR acceleration

We estimate the wind dynamics assuming less efficient CR heating.



*density is smoothed by Gaussian with the width of $\sim z_h$

Results

Initial rotation velocity: $v_{\phi} = 0.5 v_{rot}(R, z_h)$

Dissipation of Angular Momentum happens at the disk-halo interface?

the analogy to stellar winds, but in particular the CR component of the Interstellar Medium (ISM) (s does indeed secularly escape from the Galaxy to space, leading to a picture schematically given in Fi At least those CR particles which exist in the ne Hot gasfletyesolar system are observed to be well scattered i.e. to be efficiently coupled to the interstellar gas a $z_h = 2$ kiped. The gaseous halo above a few kpc is probabl has a higher degree of ionization than regions of H midplane coupling should be even better up there av Breitschwerdt+91 *density is smoothed by Gaussian with the width of $\sim z_h$





Results

Initial rotation velocity: $v_{\phi} = 0.5 v_{rot}(R, z_h)$

Dissipation of Angular Momentum happens at the disk-halo interface?

he analogy to stellar winds, but in particular the 10⁻¹⁶ CR component of the Interstellar Medium (ISM) (s does indeed secularly escape from the Galaxy to space, leading to a picture schematically given in Fi At least those CR particles which exist in the ne Hot gasfletyesolar system are observed to be well scattered i.e. to be efficiently coupled to the interstellar gas a $z_h = 2$ kipeld. The gaseous halo above a few kpc is probable has a higher degree of ionization than regions of H midplane coupling should be even better up there av Breitschwerdt+91 *density is smoothed by Gaussian with the width of $\sim z_h$



Results

*Gamma-ray from the disk is consistent with the observation (Strong+04).







Results: vs. Fermi Bubble

*Hadronic γ-ray

 The intensity of 0.1-0.5e-6 is consistent.
 The flat surface brightness profiles are well reproduced.




Results: vs. eROSITA Bubble

*Thermal X-ray emissivity = $n^2\Lambda/4\pi$ Λ =10²³ erg cm³ s⁻¹ is assumed (Shimoda & Inutsuka 22)

- 1. The flat surface brightness profiles are well reproduced.
- 2. The intensity depends on the assumed emissivity.



R (kpc)



Prospects for New Motivations

- Reacceleration at the galactic halo
 Can the turbulent halo re-accelerate CRs upto...?
 High-energy Gamma-rays & Neutrinos can be emitted from
 usual/starburst galaxies?
- \rightarrow I expect as a new candidate of the source of UHECRs .
- The Angular Momentum should be redistributed.
 The gas w/ smaller AM will return to the disk.
 The resultant AM will be recorded into the formed stars.
 The motions of stars in the disk will reflect the turbulent Halo.
 Possible *Correlations* among the CRs, Gamma-rays, Neutrinos, and Stellar Dynamics!



https://www.eurekalert.org/news-releases/6

Turbulent Halo affects the UHECRs?

the analogy to stellar winds, but in particular the CR component of the Interstellar Medium (ISM) (§ does indeed secularly escape from the Galaxy of e of f space, leading to a sicture schematically given in Fi At least those CR particles which exist in the ne of the solar system are observed to be well scattered i.e. to be efficiently coupled to the interstellar gas ε field. The gaseous halo above a few kpc is probabl has migher degree of ionization than regions of l thus coupling should be even better up there, ex in the isolated HVC's. Since the CR pressure is the gas plessure, it does not only act on the solar neighbourhood, but it must have in addition dynamical effect on the overall gas and field confi Parker 1968, 1969 for process of he parly thin matter). Because CRs do not cool raditively, w their pressure gradient to act on the gas all along

Cted's

R ~ 1 kpc ($B/1 \mu G$)⁻¹($E_{cr}/1 EeV$) Examine B-filed @ Halo!

9er

Galaxy Walls within 307 Mpc (Wikipedia)



Hints: Case of NGC 4631 (starburst?, D~7 Mpc)

SFR ~ 1 - 5 Mo/yr (Rand+92, Hunter+86, Kennicutt+83)







Carretti+13

The external galaxies exhibit many information of the Halo. I still don't know the diversity...