

Particle Acceleration and Magnetic-field Amplification by Magnetic Turbulence in the West Hot Spot of Pictor A Revealed with Mid-to-Far Infrared Observations - Isobe et al. 2017 ApJ, 850, 193 - Isobe et al. 2020 ApJ, 899, 17

N. ISOBE (ISAS/JAXA)

M. KINO (KOGAKUIN U.), Y. SUNADA, M. TASHIRO (SAITAMA U.), S. KOYAMA (ASIAA), T. NAKAGAWA (ISASA/JAXA), C. PEARSON (RAL SPACE), H. NAGAI (NAOJ)

Hot Spots of Radio Galaxies





- Compact regions with a high radio flux at the end of the AGN jets
- ✓ Strong jet-terminal shocks
- Non-thermal emission from relativistic electrons
 - Synchrotron radio emission
 - Inverse-Compton X-ray emission
- ✓ Standard picture
 - Diffusive shock acceleration (Fermi-I acceleration; Begelman et al. 1984)

✓One of the most promising candidates for ultra-high energy cosmic rays (E>10¹⁸ eV; Hillas 1984, Kotera & Olinto,2011)

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MHD simulation for the turbulence behind the shock.



Turbulence in Hot Spots

It is theoretically/numerically predicted that the shocks inevitably induce magnetic turbulence, due to plasma instabilities and/or inhomogeneities.

Possible roles of the turbulence.

- Magnetic-field amplification (e.g., Inoue et al. 2009, 2011, Sano et al. 2012)
- Particle acceleration (so-called stochastic acceleration or Fermi-II; Schlickeiser1984)
- The magnetic-field amplification by the turbulence is invoked to explain the strong magnetic field in the X-ray shell/rim of supernova remnants
 - ✓ ~100µG (Bamba et al.2003;Vink & Laming2003)

No concrete observational evidence has not yet obtained for the turbulence acceleration in the hot spots

 except for a few observational indications (Orienti et al. 2017) and theoretical applications (e.g., Asano et al. 2009, 2014, 2015, 2018)

Impotence of Infrared Observation to Search for Turbulence Acceleration

- ✓ The turbulence acceleration is expected to be a slow process.
 - Its synchrotron flux negligible in the radio band.
- ✓ The spectrum of the turbulence acceleration is significantly harder than that of the diffusive shock acceleration.
 - Spectrum of magnetic turbulences
 - $\checkmark \ |\delta B^2|_k \propto k^{-q} \ (1 < q < 2)$
 - $\checkmark k$: Wave number of turbulences
 - Electron spectrum
 - $\checkmark N(\varepsilon_e) \propto \varepsilon_e^{-p} \propto \gamma^{-p}$

$$\checkmark p = q - 1$$

- Synchrotron spectrum $\checkmark E_v \propto v^{-\alpha}$
 - ✓ Energy index: $\alpha = \frac{p-1}{2}$

(e.g., Asano et al. 2014)



2/3

-1/6

1

0

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BLACK HOLE ASTROPHYSICS WITH VLBI

Electron *p*

Synchrotron α

2

0.5

The West Hot Spot of the Radio Galaxy Pictor A



- One of the brightest and most extensively studied hot spots, thanks to its proximity (z = 0.035, $D_L = 151.9$ Mpc)
- Optical emission is detected (e.g., Thomson et al.1995)

Far/Mid Infrared Observation of the Pictor A West Hot Spot



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Spectral energy distribution of the Pictor A west hot spot



Discovery of the Far/Mid infrared excess



Discovery of the Far/Mid infrared excess



Origin of the Far/Mid Infrared excess



Acceleration process for the Far/Mid Infrared excess

The spectrum of the Far/Mid IR excess is difficult to be explained by the diffusive shock acceleration in the hot spots.

- $\sim \alpha = 0.22 \pm 0.06$ for the Far/Mid IR excess below the break
- $< \alpha > 0.5$ for the diffusive shock acceleration.

The turbulence acceleration is a promising process for generating a hard spectrum.

- The spectrum of the Far/Mid IR excess is slightly softer than that of the turbulence acceleration
 - $\checkmark \alpha = 0$ for the hard sphere limit (i.e., the softest case).
- The spectral softening by such as particle escape is required.

The origin of the Far/Mid infrared excess is the substructures within the hot spot where the electrons are energized via the turbulence acceleration.



Estimation of the magnetic field

The broken PL model for the Far/Mid excess satisfies the radiative cooling break under the continuous energy injection.

 $\checkmark \Delta \alpha = 0.5$ (e.g., Carilli et al. 1991)

Electron Lorentz factor at the cooling break: $\gamma_b = \frac{6\pi v m_e c^2}{\sigma_T B^2 R}$ (1)

 $\checkmark t_{\rm rad} = t_{\rm ad}$ (e.g., Inoue and Takahara 1996)

✓ Radiative cooling timescale: $t_{rad} = \frac{3m_e c}{4u_B \sigma_T \gamma}$ (neglecting the inverse Compton cooling)

Adiabatic loss time scale:
$$t_{ad} = \frac{R}{v}$$

✓ Synchrotron critical frequency at the break: $v_{\rm b} = \frac{3\gamma_{\rm b}^2 eB \sin \theta}{4\pi m_e c}$ (2)

Magnetic field estimated from the cooling break

$$\checkmark B^3 \simeq \frac{27\pi e m_e v^2 c}{\sigma_T^2} R^{-2} v_b \frac{-1}{(3)}$$

$$\nu_b = 1.6^{+3.0}_{-1.0} \times 10^{12} \text{ Hz}$$

R: region size (Adopting the size of the substructures)

 \checkmark v: downstream flow velocity (v = 0.3c is assumed; e.g., Kino & Takahara 2004)

 \checkmark

Estimation of the magnetic field



Consistent to the picture of magnetic field amplification by the turbulence.

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Summary and conclusion

 Far/Mid infrared observations are a possible probe for the turbulence acceleration in the hot spots of radio galaxies.

Thanks to the spectral hardness of the turbulence acceleration, its synchrotron spectrum is
possible dominates that of the diffusive shock acceleration, in the Far/Mid infrared band.

✓ By making most of the Herschel and WISE data, we searched for the signature of the turbulence acceleration in the west hot spot of Pictor A.

The Far/Mid infrared excess synchrotron emission is detected over the radio and optical synchrotron spectrum from the west hot spot of Pictor A.

✓ The Far/Mid excess is successfully described by a broken PL model.

- The radio flux is consistent to that of the substructures within this hot spot.
- ✓ The spectrum below the break ($\alpha = 0.22 \pm 0.06$) is harder than that of the diffusive shock acceleration ($\alpha > 0.5$).

The Far/Mid excess is suggested to originate in the 10-pc scale substructures within the hot spot, where the magnetic turbulence plays two important roles.

- In order to interpret the hard spectrum of the Far/Mid infrared excess, the turbulence acceleration is required.
- The strong magnetic field of the excess estimated from the cooling break condition is consistent to the magnetic-field amplification due to the turbulence.