



J-KAREN P レーザーを用いた 誘導コンプトン散乱実験の 初期報告

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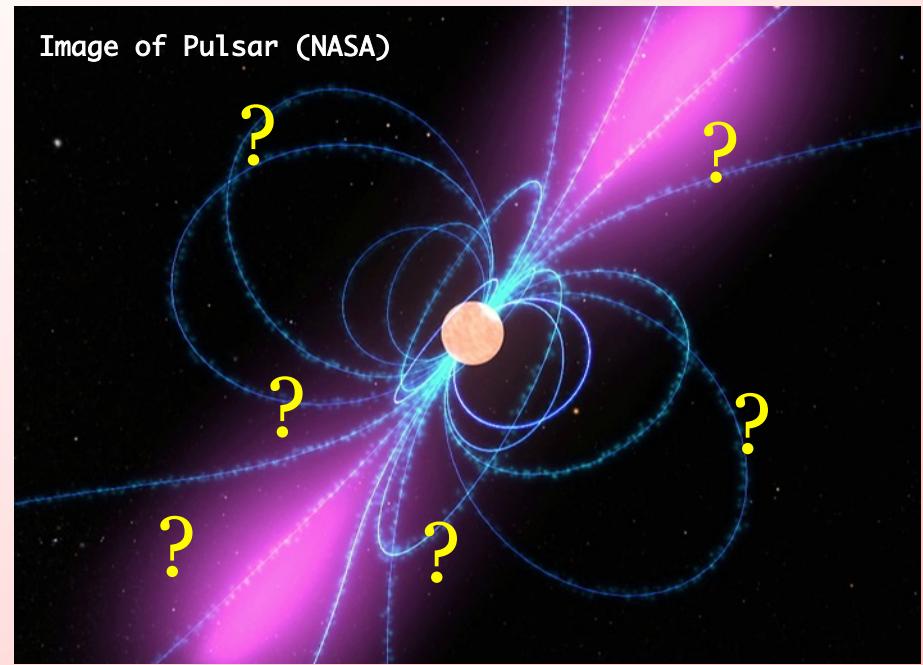
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山崎了(青学)、坂和洋一(阪大)



Image of Pulsar (NASA)

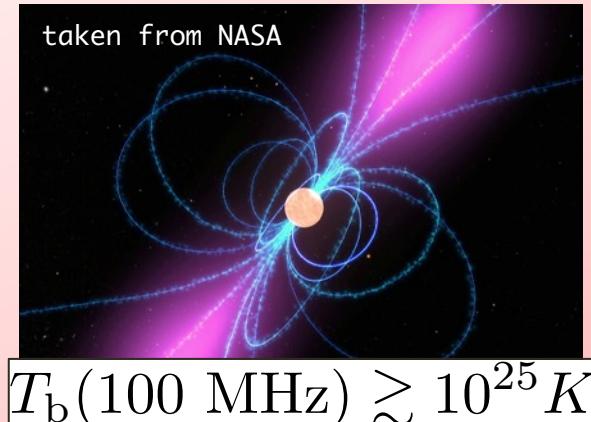


Induced Compton Scattering

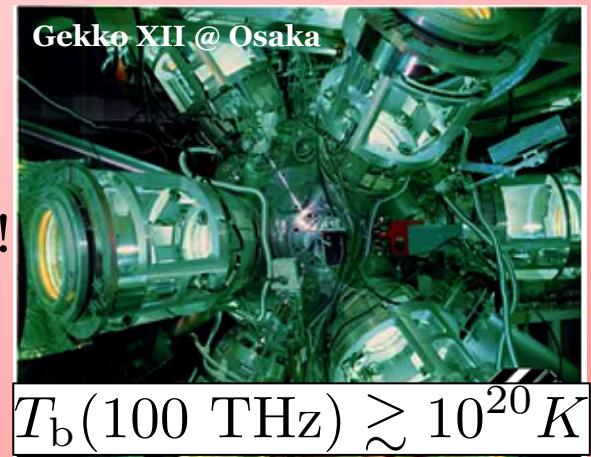
Induced Compton Scattering

Interaction between rarefied plasma & bright radiation

- Rarefied plasma ($\lambda < \lambda_D$, $\omega > \omega_{pe}$)
 - Scattering of photons by an electron
 - Cross section is given by Klein-Nishina formula
- Bright radiation ($k_B T_b \gg m_e c^2$)
 - $k_B T_b(\nu) \equiv h\nu n_{ph}(\nu) \equiv E/(\Delta t \Delta \nu)$
 - $n_{ph} > 2$ is possible for Boson \Leftrightarrow induced process rather than exclusion one! $n_{ph} \sim 10^{27}$ for pulsar!!

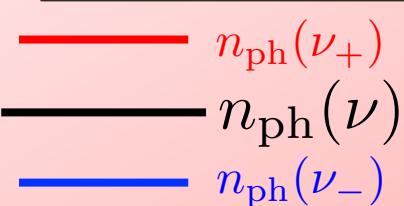


Pulsed radio emission

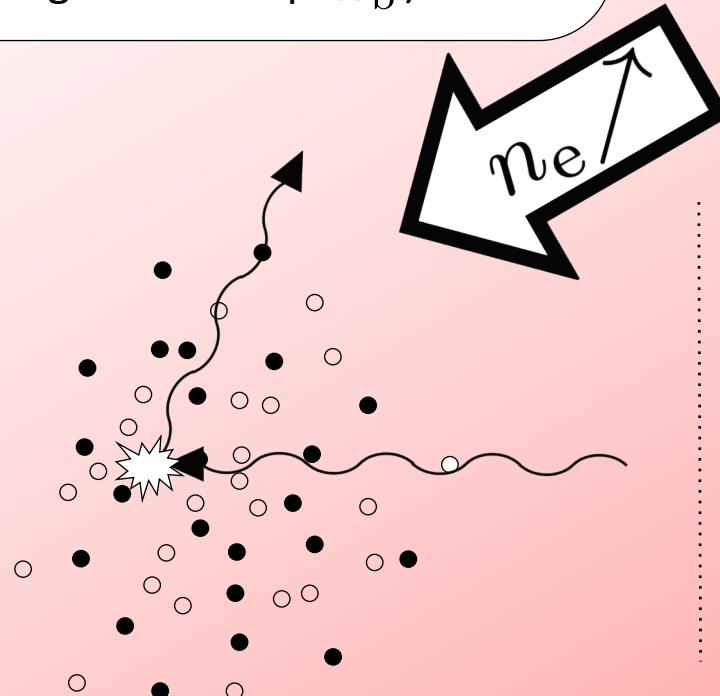
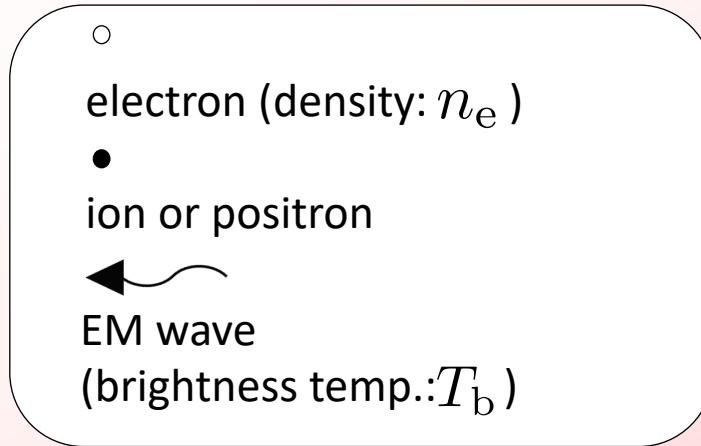


Laser facilities

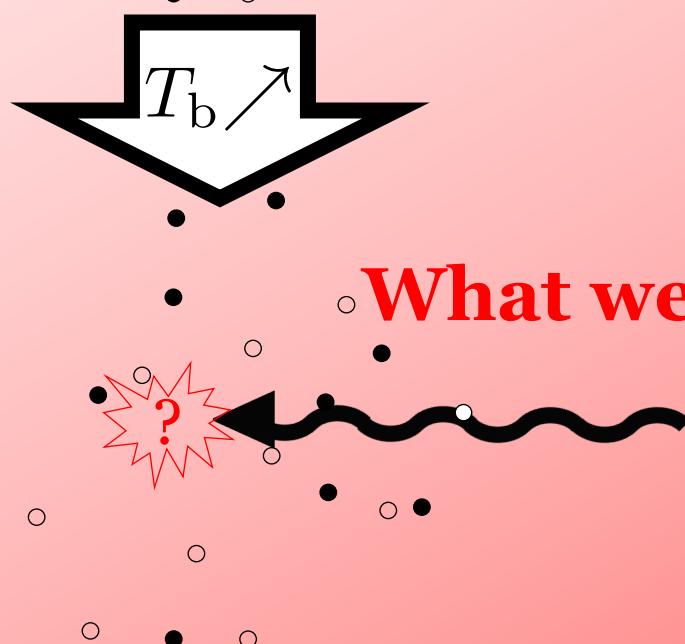
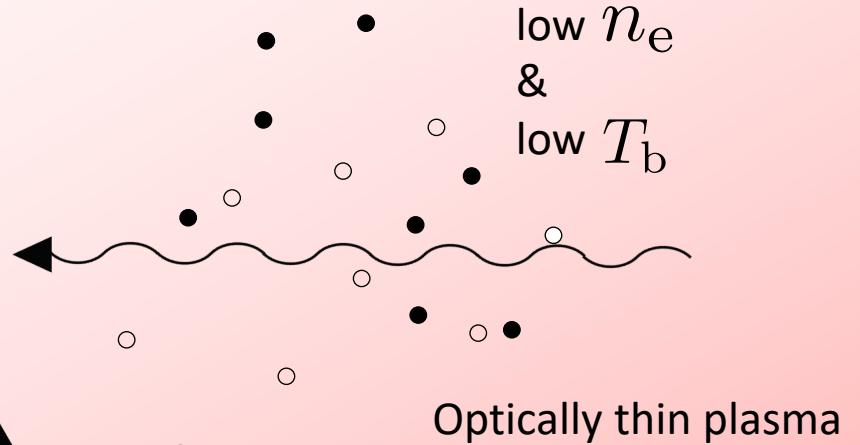
$$\frac{dn_{ph}(\nu)}{dt} \propto n_{ph}(\nu_+) \underbrace{(1 + n_{ph}(\nu))}_{\text{spontaneous + induced terms}} - n_{ph}(\nu) \underbrace{(1 + n_{ph}(\nu_-))}_{\text{spontaneous + induced terms}}$$



Spontaneous vs. Induced



Thomson scattering

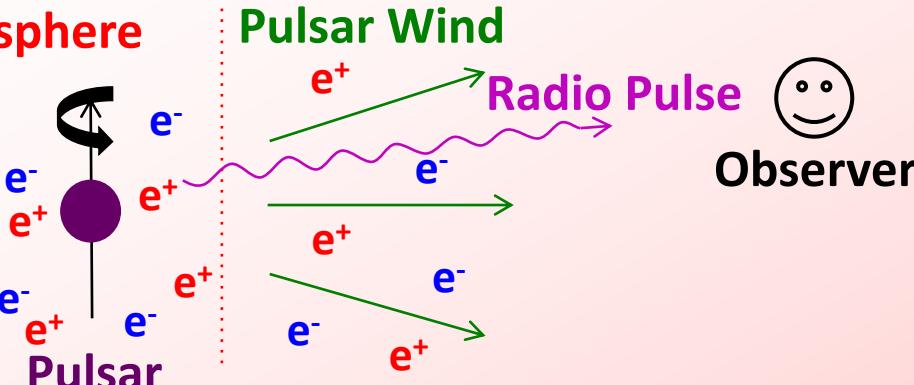


Induced Compton scattering

What we see?

ICS & PSR

Magneto-sphere

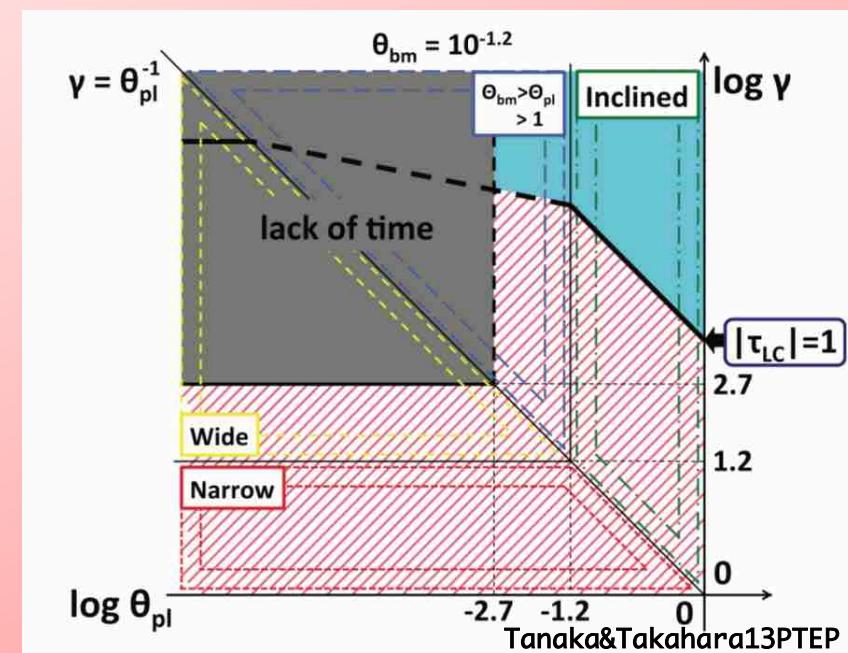


e.g., Tanaka&Takahara13PTEP

ICS has potential to constrain
• \dot{N} : particle number flux
• σ : magnetization
• γ : bulk Lorentz factor
of the pulsar wind

$$L_{\text{spin}} = \dot{N}(1+\sigma)\gamma m_e c^2$$

Interaction between radio pulse (laser) and pulsar wind (plasma).
=> Characteristic signature?



Basic Equation

Kinetic Equation for Photon

Compton scattering off photons $n_{\text{ph}}(\mathbf{k})$ by plasmas $f(\mathbf{p})$.

$$\left(\frac{\partial}{\partial t} + c\boldsymbol{\Omega} \cdot \nabla \right) n(\mathbf{k}) = cn_{\text{pl}} \int d^3\mathbf{p} f(\mathbf{p}) \int d^3\mathbf{k}_1$$

Boltzmann-Uehling-Uhlenbeck Equation

$$\times [\sigma_{\text{KN}}(\mathbf{k}_1, \mathbf{k}, \mathbf{p}) n(\mathbf{k}_1) (1 + \underline{n(\mathbf{k})}) - \sigma_{\text{KN}}(\mathbf{k}, \mathbf{k}_1, \mathbf{p}) n(\mathbf{k}) (1 + \underline{n(\mathbf{k}_1)})]$$

induced term

Cross-section for Compton scattering
(Klein-Nishina formula)

$$\sigma_{\text{KN}}(\mathbf{k}_i, \mathbf{k}_f, \mathbf{p}) = \frac{r_e^2}{2\gamma^2 k_i k_f} \delta \left(k_f D_f - k_i D_i + \frac{k_i k_f (1 - \mu)}{\gamma} \right)$$
$$\times \left[1 + \left(1 - \frac{1 - \mu}{\gamma^2 D_i D_f} \right)^2 + \frac{k_i k_f (1 - \mu)^2}{\gamma^2 D_i D_f} \right]$$

Compton effect:
not symmetric
about $k_i \Leftrightarrow k_f$

$$D_{i,f} = 1 - \boldsymbol{\beta} \cdot \boldsymbol{\Omega}_{i,f} \quad \mu = \boldsymbol{\Omega}_i \cdot \boldsymbol{\Omega}_f$$

Kompaneets equation

uniform + isotropic + 1st order in $h\nu \ll m_e c^2, k_B T_e \ll m_e c^2$

Kompaneets 1957

$$\frac{\partial n(x)}{\partial y} = \frac{1}{x^2} \frac{\partial}{\partial x} x^4 \left(n(x) + n^2(x) + \frac{\partial n(x)}{\partial x} \right)$$

$x \equiv \frac{h\nu}{k_B T_{\text{pl}}}, y \equiv \frac{k_B T_{\text{pl}}}{m_e c^2} n_{\text{pl}} \sigma_T c t$

- Photon number conservation
- Bose-Einstein distribution as equilibrium solution
- No Thomson scatt. (0th order) because of isotropy.
- 1st term = Compton effect
(energy loss for photon)
- 2nd term = Induced Compton
(energy loss for photon)
- 3rd term = Inverse Compton
(energy gain for photon)

$$\tau_{\text{Comp}} \approx \sigma_T \ln n_{\text{pl}} \times \frac{h\nu}{m_e c^2}$$

$$\tau_{\text{ind}} \approx \sigma_T \ln n_{\text{pl}} \times \frac{k_B T_b(\nu)}{m_e c^2}$$

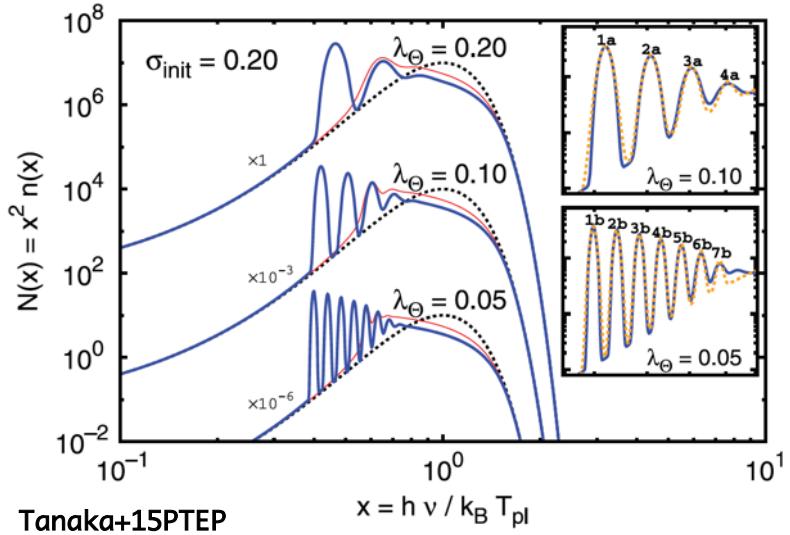
$$\tau_{\text{IC}} \approx \sigma_T \ln n_{\text{pl}} \times \frac{k_B T_e}{m_e c^2} \approx y$$

The case for $n_{\text{ph}} \gg 1$

$$n\Theta \times \min(1, x) \gg 1$$

$$\frac{\partial g}{\partial y} - 2g \frac{\partial g}{\partial x} = 0 \quad \xrightarrow{\text{2nd order}}$$

$$\frac{\partial g}{\partial y} - 2g \frac{\partial g}{\partial x} + \frac{17\Theta}{5}g \frac{\partial}{\partial x}g = \frac{14\Theta}{5}(xg) \frac{\partial^3}{\partial x^3}(xg)$$



$$\frac{\partial g}{\partial y} - 2g \frac{\partial g}{\partial x} \approx \frac{14\Theta}{5}(xg) \frac{\partial^3}{\partial x^3}(xg)$$

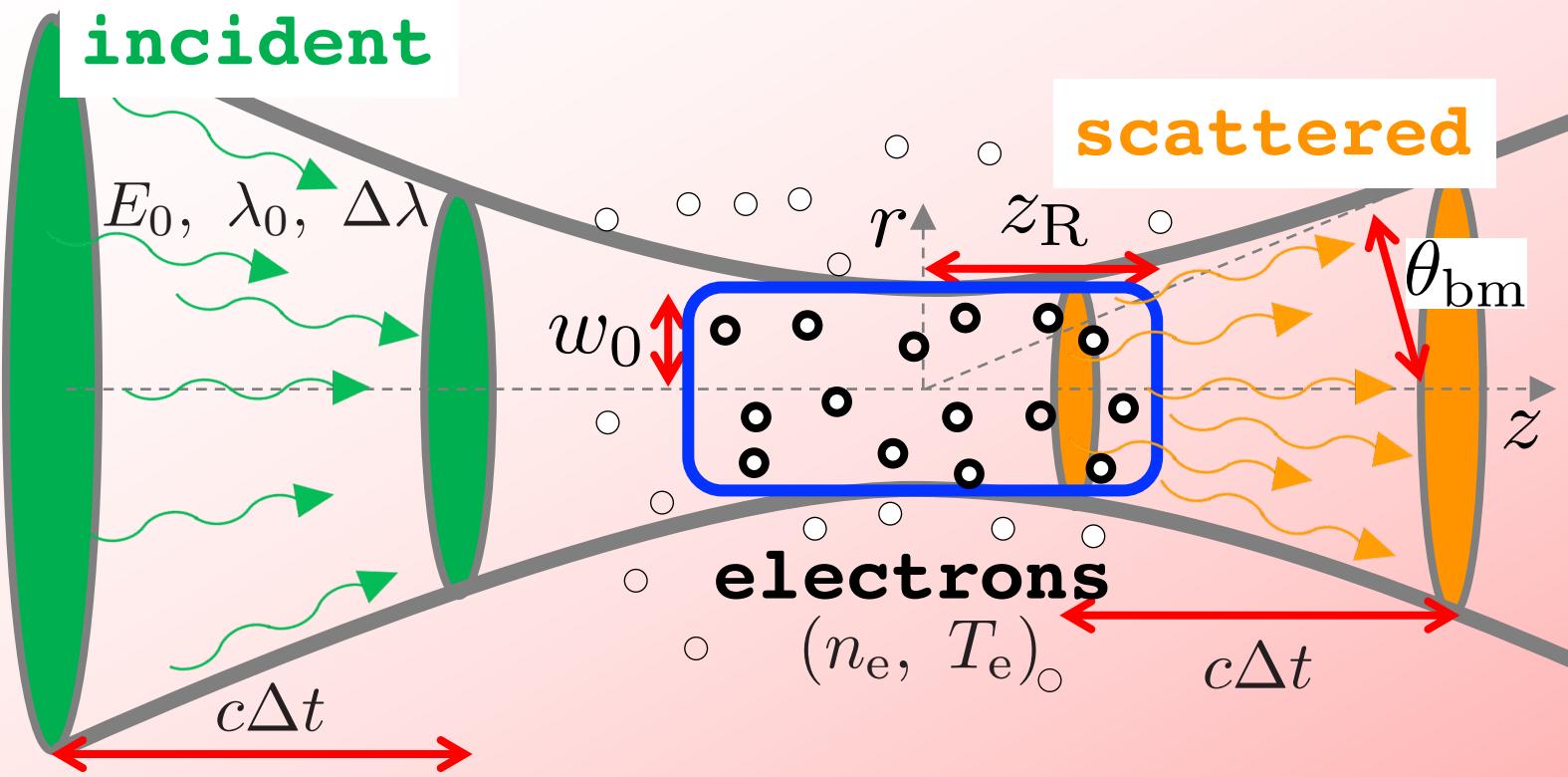
order of $(hv)\Theta/(m_e c^2)^2$

$$x = \frac{h\nu}{k_B T_e},$$

$$\Theta = \frac{k_B T_e}{m_e c^2},$$

$$g(x) = x^2 n(x)$$

Solitary structures in spectrum



Predictions in Experiments

Four Constraints on Plasma

- Typical optical depth for ICS (blue)

$$\tau_{\text{ICS}} > 0.1 \quad \rightarrow$$

$$n_e > 0.1 \times \frac{16\pi^3 m_e c^2 \Delta t \Delta \nu w_0^2}{6\sigma_T E \lambda^3}$$

- Doppler width (λ_Θ) of solitary waves (green)

$$\frac{\lambda_\Theta}{2} < \frac{\Delta \nu}{\nu} \quad \rightarrow$$

$$k_B T_e < \frac{3m_e (w_0 \Delta \nu)^2}{2}$$

- No screening effect (e.g., Galeev & Syunyaev 1973) (red)

$$\Delta \nu > \nu_{pe} \quad \rightarrow$$

$$n_e < \frac{\pi m_e \Delta \nu^2}{e^2}$$

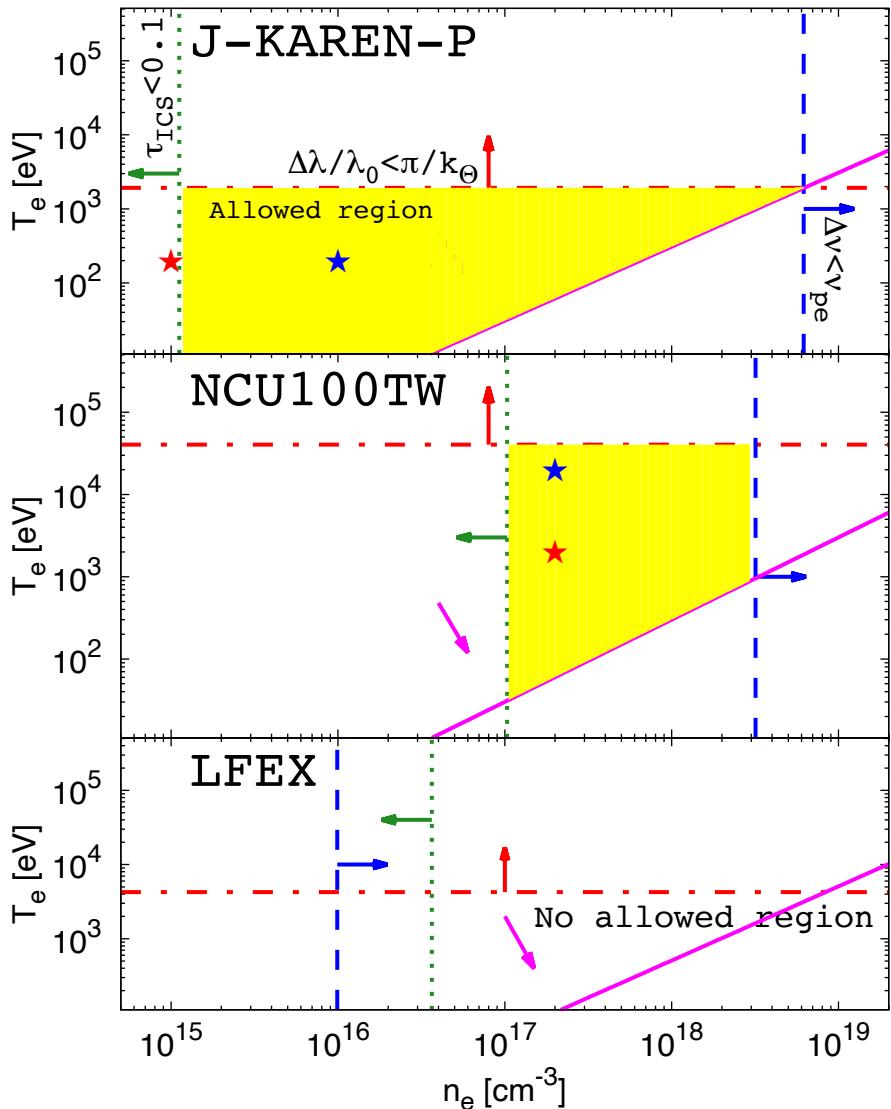
- Non-collective scattering (yellow)

$$\lambda < \lambda_D \quad \rightarrow$$

$$k_B T_e > 4\pi e^2 \lambda^2 n_e$$

Laser Facilities

Tanaka+20PTEP



Allowed plasma parameters are found in yellow region.

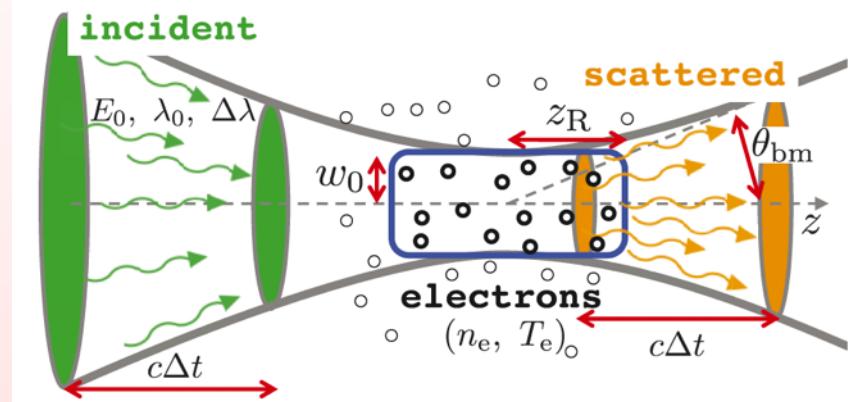
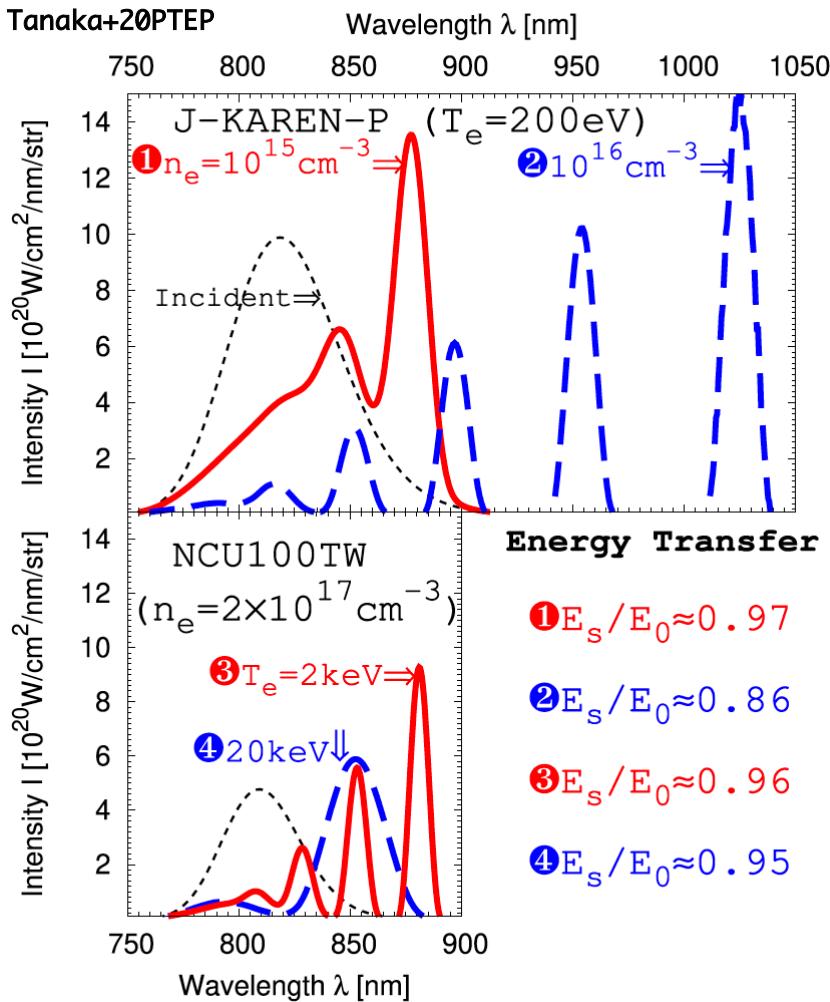
High-power short pulse laser is favored for ICS experiments rather than high (total) energy laser.

We can draw the same plot for other facilities of the given parameters!

- E : total energy
- Δt : pulse width
- Δv : band width
- w_o : minimum waist
- λ : central wavelength

Predictions

Spectra of transmitted (scattered) light

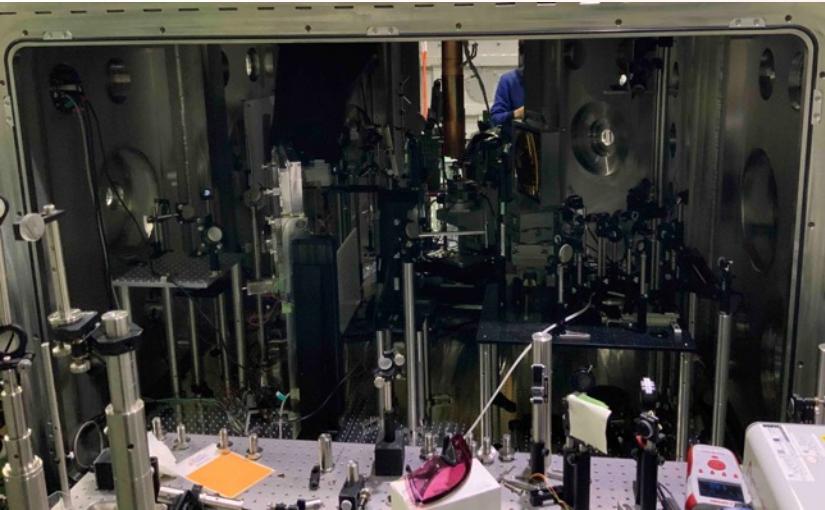


| Parameter | J-KAREN-P | NCU100TW | LFEX |
|---|-----------------------|-----------------------|-----------------------|
| E_0 [J] | 10 | 3.3 | 400 |
| λ_0 [nm] | 820 | 810 | 1053 |
| $\Delta\lambda$ [nm] | 50 | 35 | 3.3 |
| Δt [fs] | 30 | 30 | 1500 |
| w_0 [μm] | 0.67 | 4.3 | 50 |
| $k_B T_b/m_e c^2$ | 1.8×10^{14} | 8.4×10^{13} | 3.6×10^{15} |
| $\Delta\lambda/\lambda_0$ | 6.1×10^{-2} | 4.3×10^{-2} | 3.1×10^{-3} |
| θ_{bm} | 3.9×10^{-1} | 6.0×10^{-2} | 1.3×10^{-2} |
| z_R [μm] | 1.7 | 72 | 1.9×10^3 |
| τ_{ICS}/n_e [cm^3] | 9.0×10^{-17} | 9.7×10^{-19} | 2.7×10^{-18} |
| $\tau_D/(n_e \Theta)$ [cm^3] | 9.0×10^{-18} | 2.3×10^{-21} | 3.3×10^{-22} |
| τ_{Th}/n_e [cm^3] | 2.3×10^{-28} | 9.5×10^{-27} | 2.5×10^{-25} |

Only $\pi w_0^2 2z_R n_e \sim 10^4$ electrons at Rayleigh region

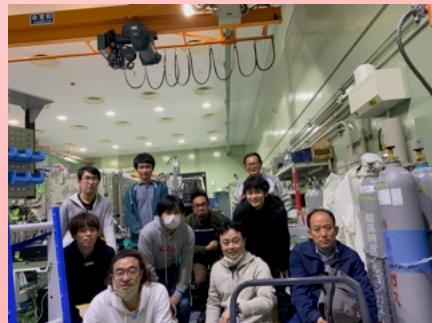
They would attain $\sim \text{PeV}$ which is rad. reaction limited

$$\frac{dE}{dx} \approx \frac{3}{2} \frac{m_e c^2}{r_e} \approx 2 \times 10^{14} \text{ MeV/m} = 0.2 \text{ PeV}/\mu\text{m}$$



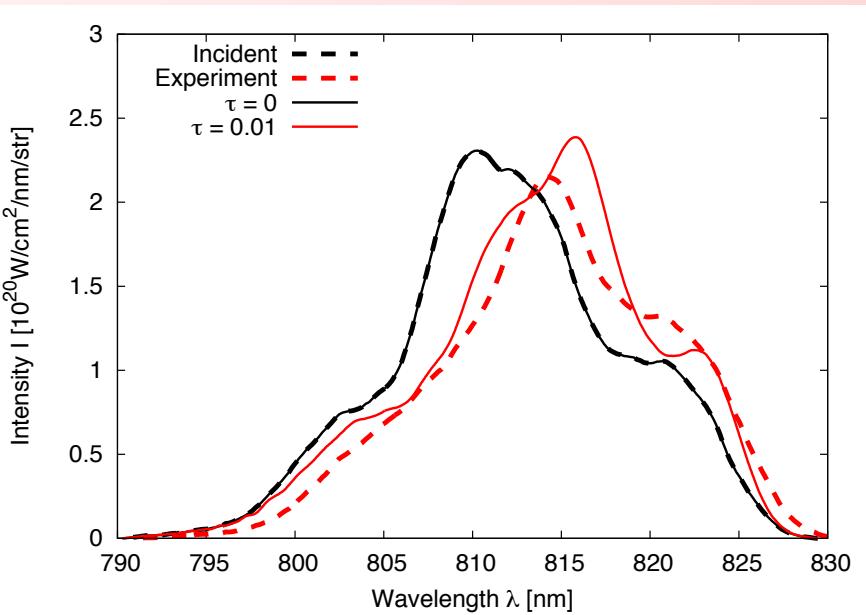
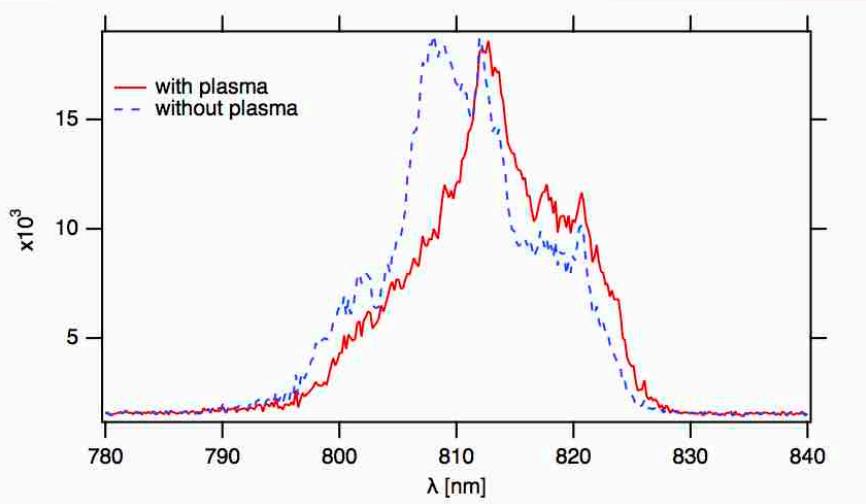
J-KAREN Experiment

2020/12/2, 3

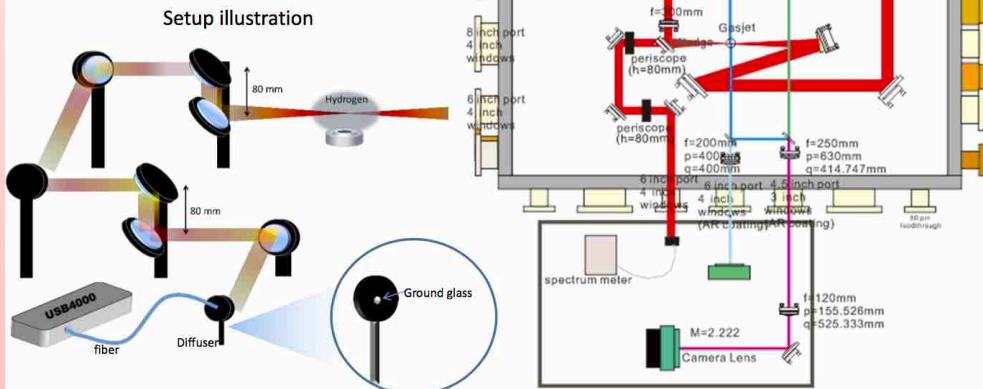


Initial Experiment

July 2017, National Central University @ Taiwan



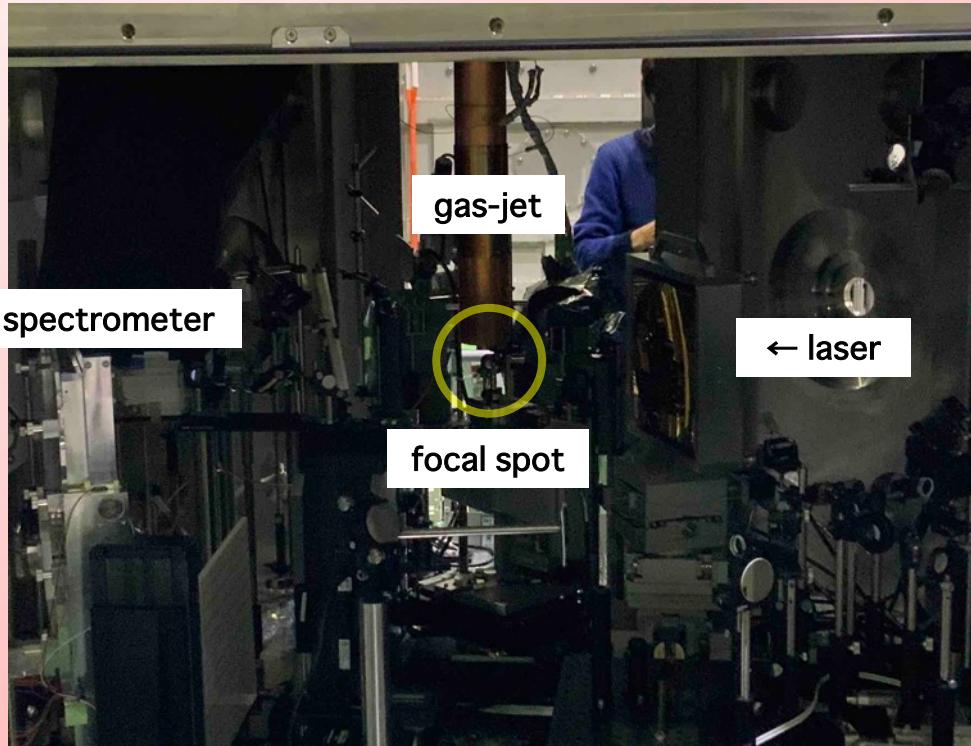
Experimental Setup



Produced by Kuramitsu-san

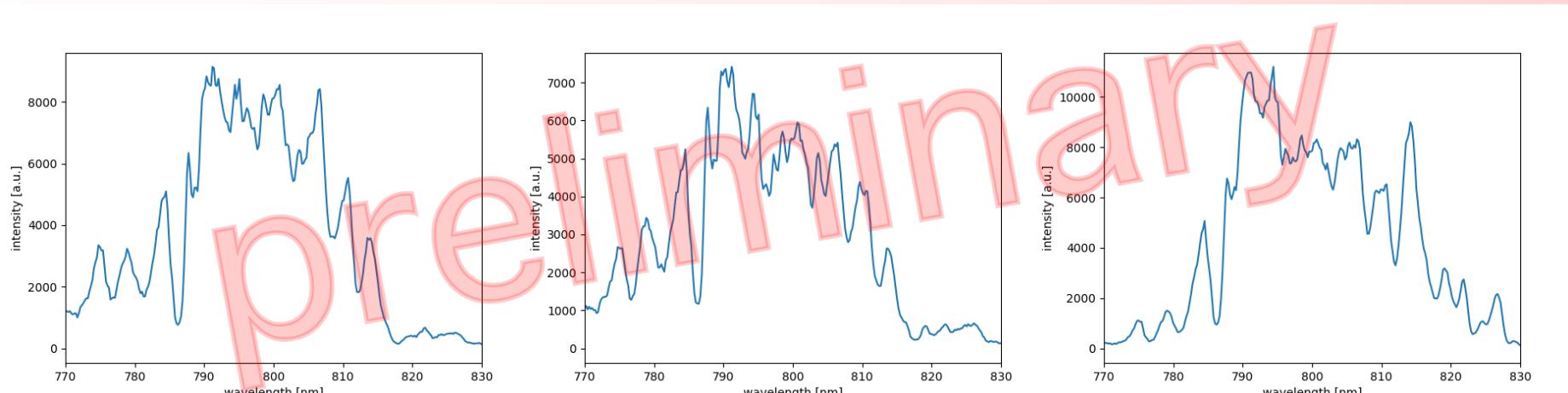
J-KAREN experiment

- What do we try to observe?
 - spectrum of scattered light (redshifted compared with incident one)
 - no side- & back-scattering, no change of polarization
 - dependence on electron density (optical depth)
 - dependence on electron temperature
 - acceleration of electrons (radiation reaction limited)

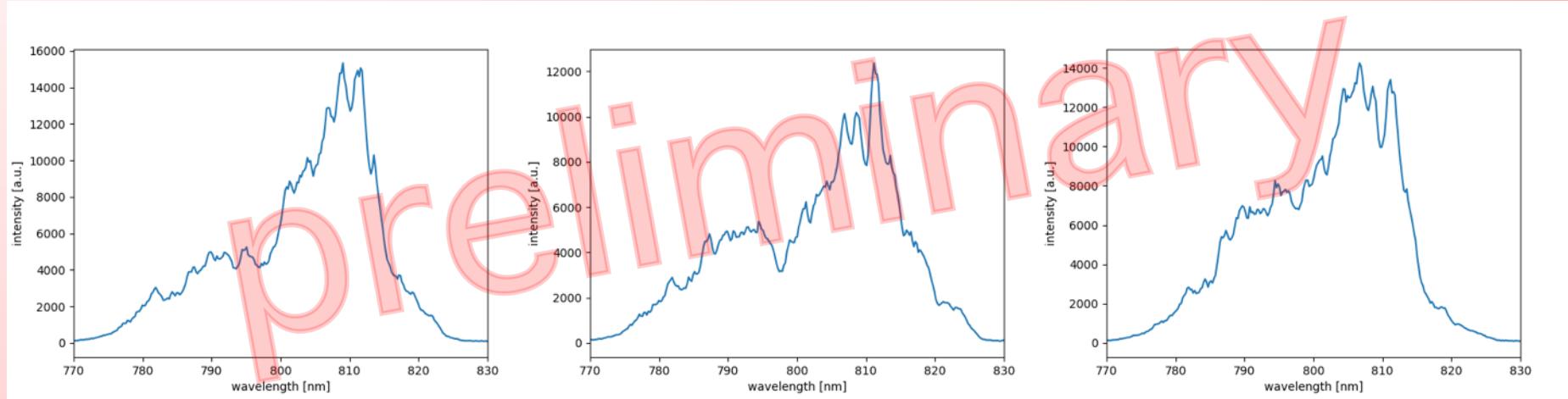


Spectrometer

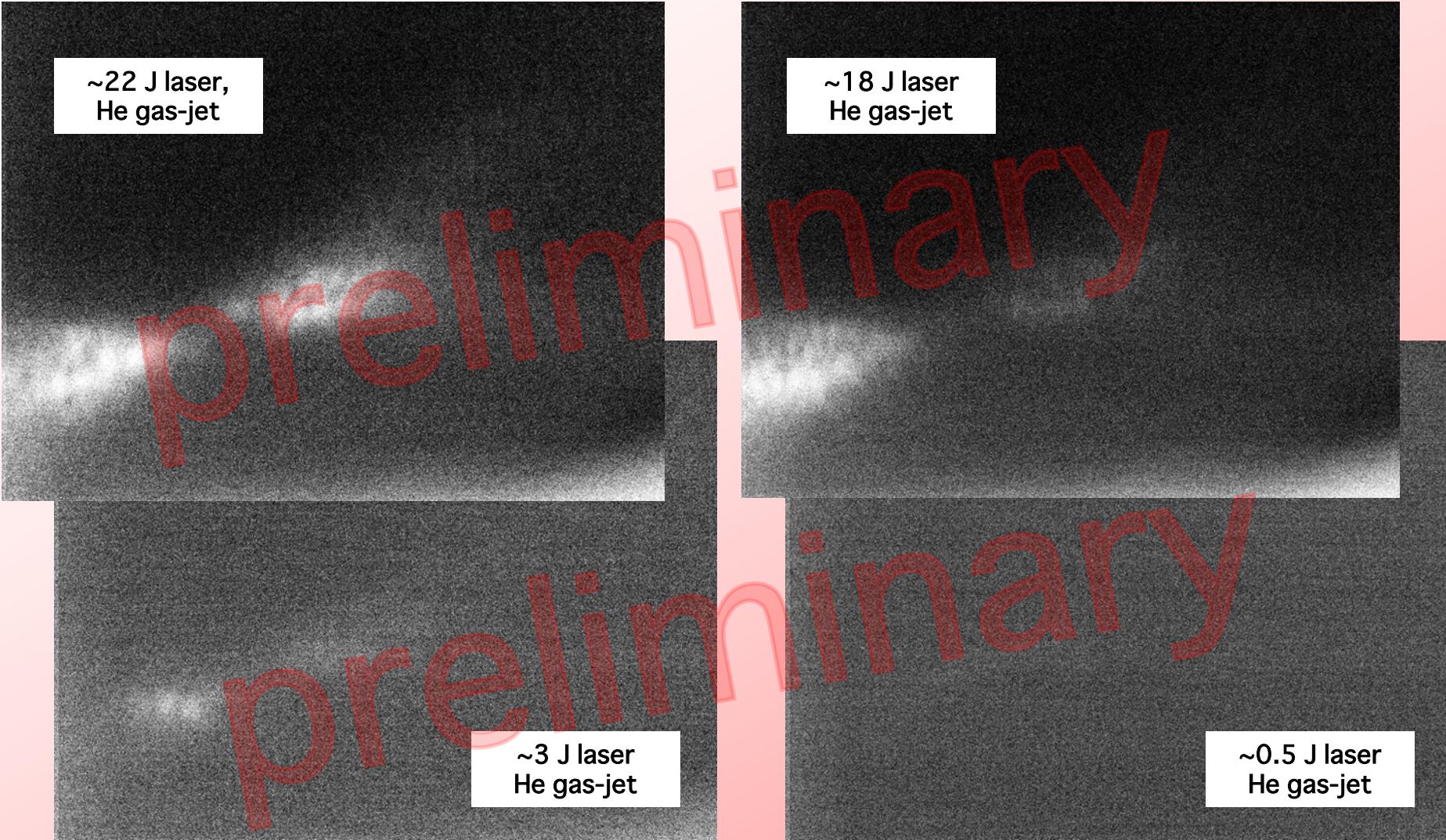
He gas-jet



gas-jet off



Imaging (Thomson scattering)



Summary

- J-KAREN Pレーザーを用いた誘導コンプトン散乱の実験を行った。
- Spectrometer: 透過光スペクトルを解析中だが、誘導コンプトン散乱による明確なスペクトル変調はまだ確認できていない。
- Imaging (Thomson scattering): 集光点付近にプラズマがない可能性がある。
- 2021年1月に分光器とガス密度の校正実験を行う。
- 系統的に解析を進めて次の実験に繋げる!

Supported by

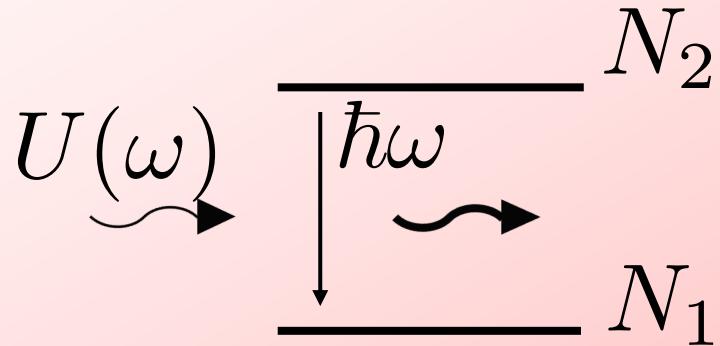
JSPS (17K18270), Osaka Univ. (2020B2-TANAKA), Aoyama Gakuin Univ. (Early Eagle Program)

Various light scattering processes

- Geometric optics limit: $R > \lambda$, R = size of scatterer, λ = light wavelength
- Mie scattering: $R \sim \lambda$
- Rayleigh scattering: $R < \lambda$
- Brillouin scattering: Phonon, ion acoustic wave
- Raman scattering: Plasmon, vibration and/or rotational level of a molecule
- Compton scattering: an electron $\gamma + e \rightarrow \gamma + e$
 - Thomson scattering: classical limit of Compton ($h = 0$, elastic)
 - Induced Compton scattering: induced counterpart, **redshift**
 - Nonlinear Compton scattering: multiple photon absorption, **blueshift**
$$N\gamma + e \rightarrow \gamma + e$$
 - Double Compton scattering: photon splitting, **redshift**
$$\gamma + e \rightarrow 2\gamma + e$$

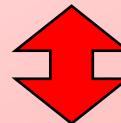
Induced effect in laser vs. ICS

- number in state i : N_i
- pumping rate to state i : Φ_i
- decay rate from state i : γ_i
- Einstein's B coefficient: B
- energy level between state 1 & 2: $\hbar\omega$
- photon energy density at ω : $U(\omega)$



$$\frac{dN_2}{dt} = \Phi_2 - \gamma_2 N_2 - (N_2 - N_1) BU(\omega)$$

$$\frac{dN_1}{dt} = \Phi_1 - \gamma_1 N_1 + (N_2 - N_1) BU(\omega)$$



$$\boxed{\frac{dn_{\text{ph}}(\nu)}{dt} \propto n_{\text{ph}}(\nu_+)(1 + n_{\text{ph}}(\nu)) - n_{\text{ph}}(\nu)(1 + n_{\text{ph}}(\nu_-))}$$

\rightarrow
$$\frac{dn_{\text{ph}}(\nu)}{dt} \propto n_{\text{ph}}(\nu_+) - n_{\text{ph}}(\nu) + (n_{\text{ph}}(\nu_+) - n_{\text{ph}}(\nu_-)) n_{\text{ph}}(\nu)$$