

# Remnants of first stars for gravitational wave sources

Tomoya Kinugawa

Institute for Cosmic Ray Research

The University of Tokyo

collaborators: T. Nakamura, K. Inayoshi, K. Hotokezaka, D. Nakauchi

A. Miyamoto, N. Kanda, A. Tanikawa, T. Yoshida



# The beginning of Gravitational wave astronomy

- Gravitational wave detectors

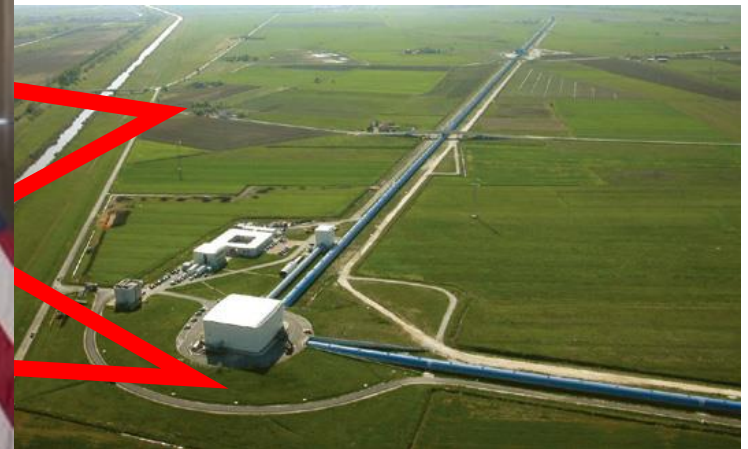
KAGRA



©KAGRA



Advanced VIRGO



©VIRGO



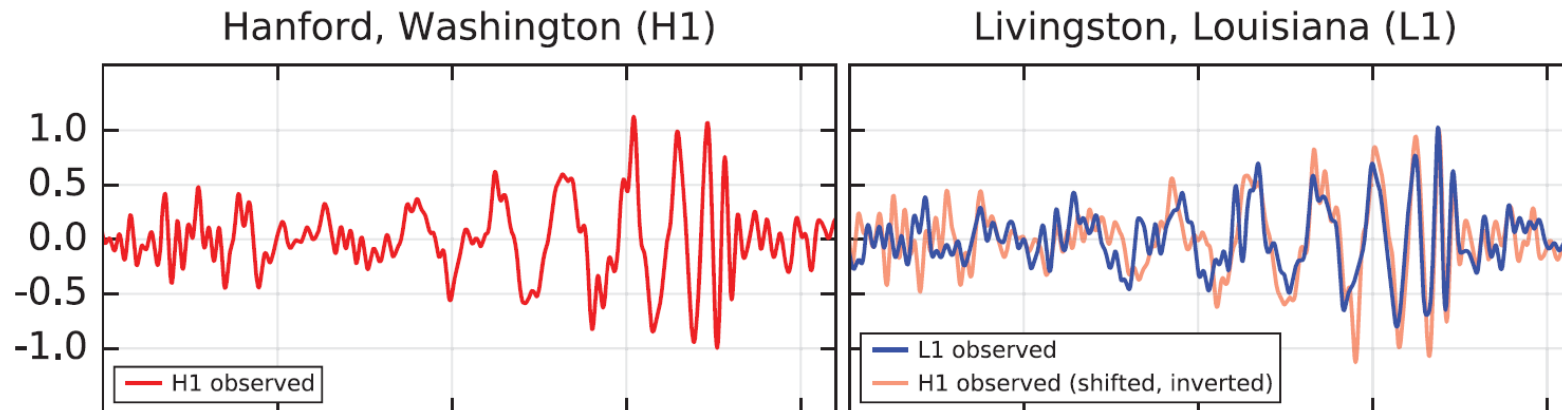
# Observation of Gravitational Waves from a Binary Black Hole Merger

B. P. Abbott *et al.*\*

(LIGO Scientific Collaboration and Virgo Collaboration)

(Received 21 January 2016; published 11 February 2016)

On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed a transient gravitational-wave signal. The signal sweeps upwards in frequency from 35 to 250 Hz with a peak gravitational-wave strain of  $1.0 \times 10^{-21}$ . It matches the waveform predicted by general relativity for the inspiral and merger of a pair of black holes and the ringdown of the resulting single black hole. The signal was observed with a matched-filter signal-to-noise ratio of 24 and a false alarm rate estimated to be less than 1 event per 203 000 years, equivalent to a significance greater than  $5.1\sigma$ . The source lies at a luminosity distance of  $410^{+160}_{-180}$  Mpc corresponding to a redshift  $z = 0.09^{+0.03}_{-0.04}$ . In the source frame, the initial black hole masses are  $36^{+5}_{-4} M_{\odot}$  and  $29^{+4}_{-4} M_{\odot}$ , and the final black hole mass is  $62^{+4}_{-4} M_{\odot}$ , with  $3.0^{+0.5}_{-0.5} M_{\odot} c^2$  radiated in gravitational waves. All uncertainties define 90% credible intervals. These observations demonstrate the existence of binary stellar-mass black hole systems. This is the first direct detection of gravitational waves and the first observation of a binary black hole merger.



# Masses of GW events

- GW events show that there are many massive BHs ( $\gtrsim 30 M_{\text{sun}}$ ).
- 7/10 BBHs are massive BBHs
- On the other hand, the typical mass of BHs in X-ray binaries is  $\sim 10 M_{\text{sun}}$ .

Event	$m_1/M_{\odot}$	$m_2/M_{\odot}$
GW150914	$35.6^{+4.8}_{-3.0}$	$30.6^{+3.0}_{-4.4}$
GW151012	$23.3^{+14.0}_{-5.5}$	$13.6^{+4.1}_{-4.8}$
GW151226	$13.7^{+8.8}_{-3.2}$	$7.7^{+2.2}_{-2.6}$
GW170104	$31.0^{+7.2}_{-5.6}$	$20.1^{+4.9}_{-4.5}$
GW170608	$10.9^{+5.3}_{-1.7}$	$7.6^{+1.3}_{-2.1}$
GW170729	$50.6^{+16.6}_{-10.2}$	$34.3^{+9.1}_{-10.1}$
GW170809	$35.2^{+8.3}_{-6.0}$	$23.8^{+5.2}_{-5.1}$
GW170814	$30.7^{+5.7}_{-3.0}$	$25.3^{+2.9}_{-4.1}$
GW170817	$1.46^{+0.12}_{-0.10}$	$1.27^{+0.09}_{-0.09}$
GW170818	$35.5^{+7.5}_{-4.7}$	$26.8^{+4.3}_{-5.2}$
GW170823	$39.6^{+10.0}_{-6.6}$	$29.4^{+6.3}_{-7.1}$

# Origin of massive BBHs

7/10 GW BBHs are massive BBHs

In order to explain the origin of such massive BBHs

Many theories exist such as

- 1)Pop II BBH
- 2)Pop III BBH **No metal field binaries**
- 3)Primordial Binary BH
- 4)N body origin from Globular Cluster
- .....

# Pop III binary population synthesis

We simulate  $10^6$  Pop III-binary evolutions and estimate how many binaries become compact binary which merges within Hubble time.

× 84 models (*Kinugawa et al. 2014, 2016*)

Initial stellar parameters are decided by Monte Carlo method with initial distribution functions

- Initial parameter (M1, M2, a, e) distribution in our standard model

M1 : Flat ( $10 M_{\odot} < M < 100 M_{\odot}$ )

$q = M2/M1$  :  $P(q) = \text{const.}$  ( $0 < q < 1$ )

a :  $P(a) \propto 1/a$  ( $a_{\min} < a < 10^6 R_{\odot}$ )

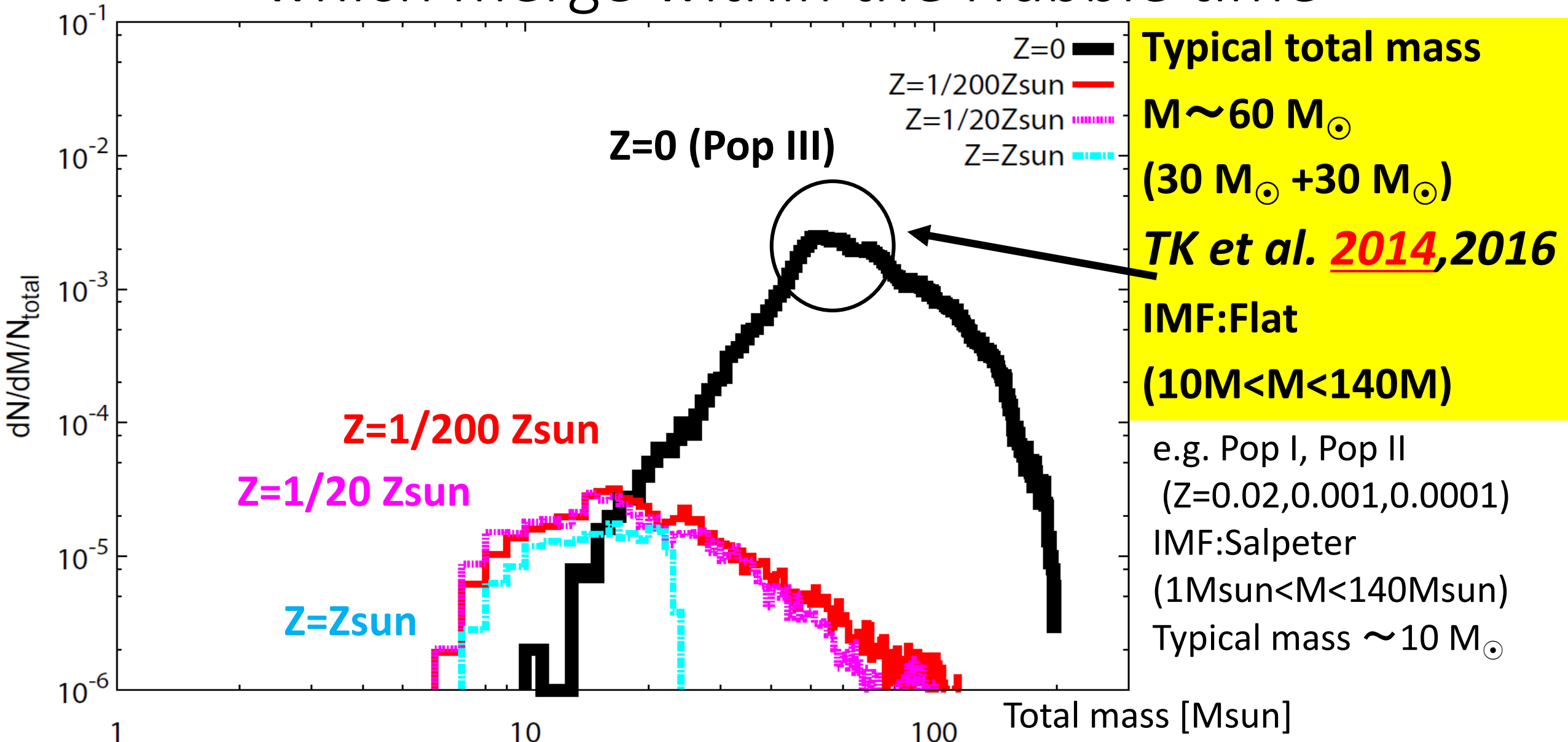
e :  $P(e) \propto e$  ( $0 < e < 1$ )



The same distribution functions  
adopted for Pop I population  
synthesis

- de Souza SFR

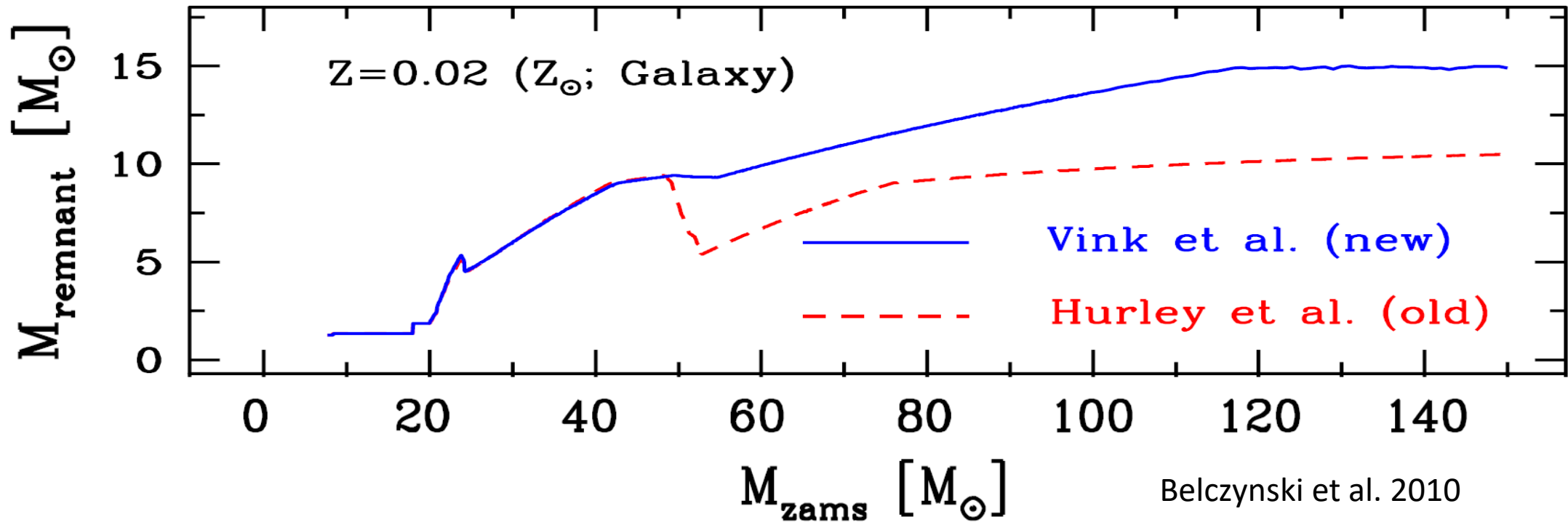
# Total mass distribution of BBH which merge within the Hubble time





# Wind mass loss & IMF

- If the progenitor of BH is Pop I (=Solar metal stars)
- Typical mass is small ( $\text{IMF} \propto M^{-2.35}$ ,  $0.1 M_{\text{sun}} < M < 100 M_{\text{sun}}$ )
- Stars lose a lot of mass due to the strong stellar wind



- The orbit become wide due to stellar wind mass loss



# Wind mass loss & IMF

- If the progenitor is low metal,

- Pop II (Metal < 0.1 Solar Metal)  
Typical mass is same as Pop I  
But, weak wind mass loss

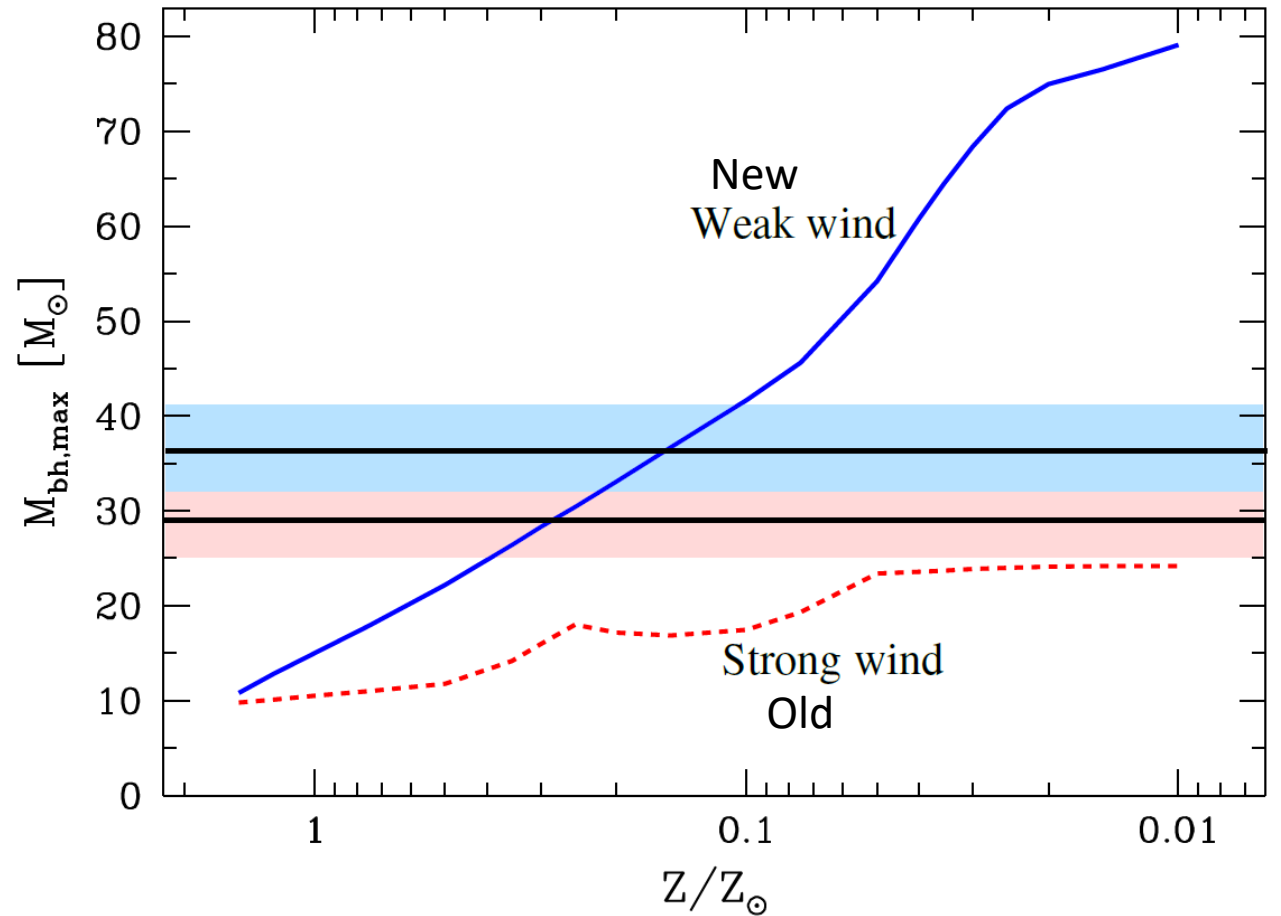
- Pop III (No metal)

Pop III stars are *the first stars* after the Big Bang.

Typical mass is more massive than Pop I, II

$M_{\text{pop III}} \sim 10\text{-}100 M_{\text{sun}}$

No wind mass loss due to no metal.

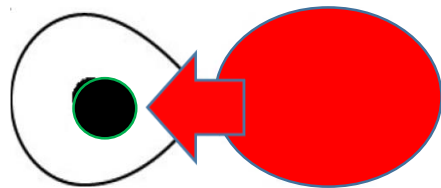


Initial:  $8 M_{\text{sun}} < M < 150 M_{\text{sun}}$

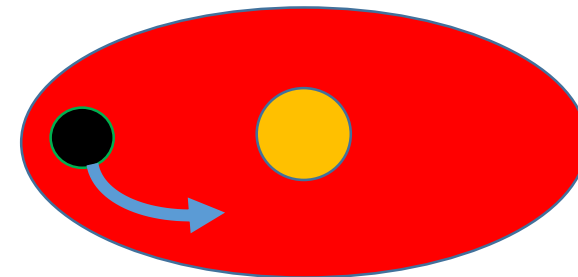
Single stellar evolution  
with 2 stellar wind models.  
(Belczynski et al. 2010,  
Abbot et al. 2016)

# Binary interaction changes progenitor mass

- Mass transfer
- Common envelope



Mass transfer

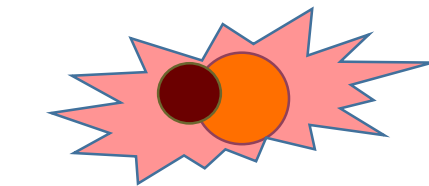


Common envelope

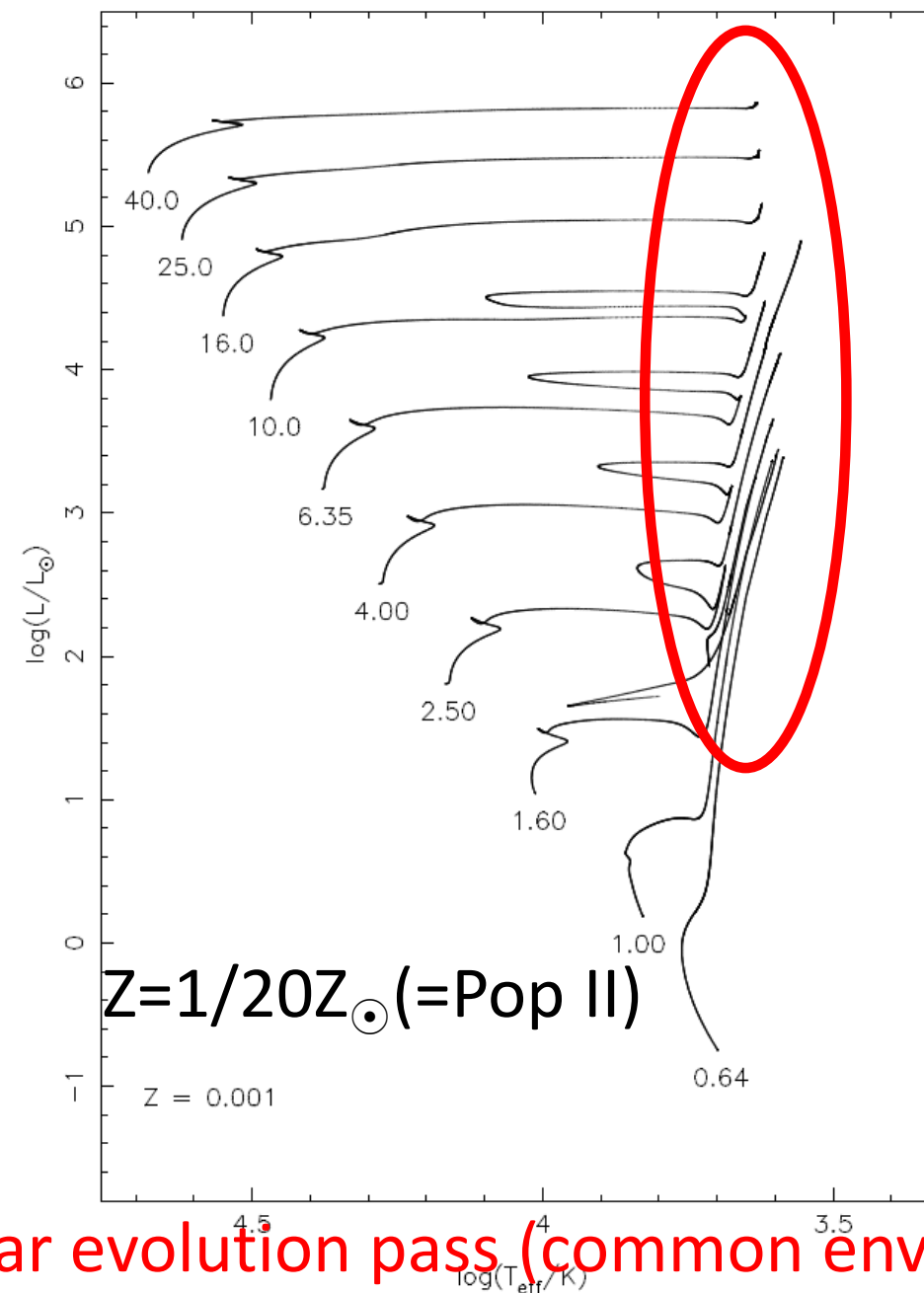
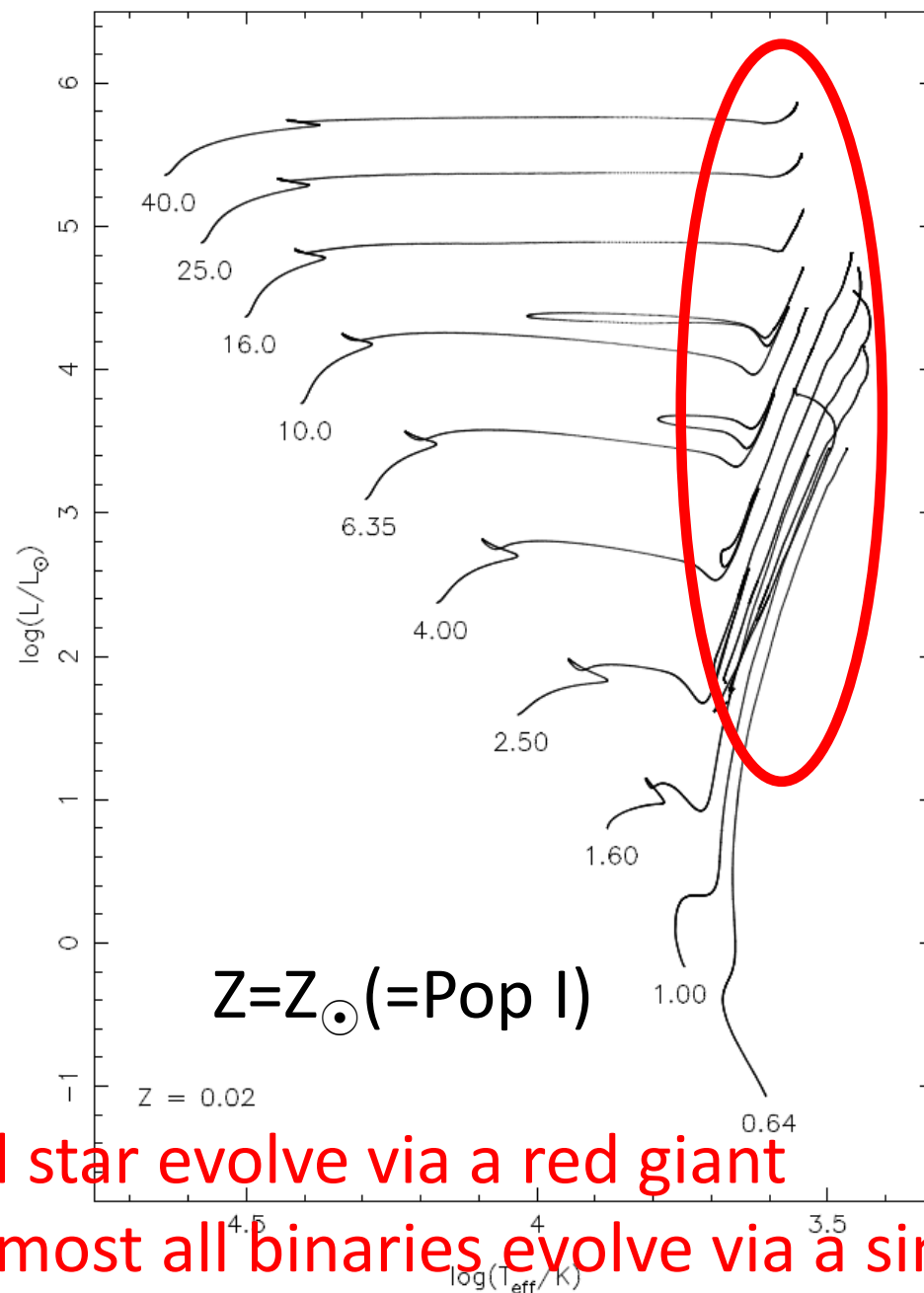
Red Giants tend to  
become CE



Close binary



or merge



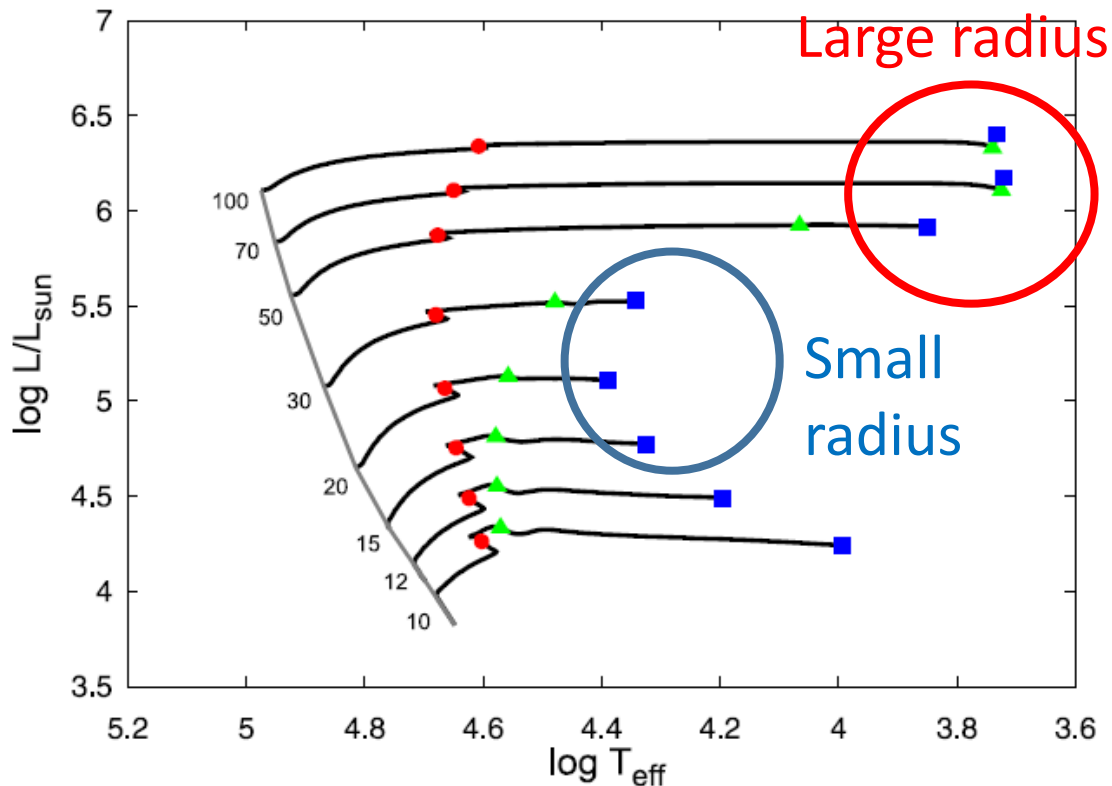
All stars evolve via a red giant

Almost all binaries evolve via a similar evolution pass (common envelope)

**Figure 1.** Selected OVS evolution tracks for  $Z = 0.02$ , for masses 0.64, 1.0, 1.6, 2.5, 4.0, 6.35, 10, 16, 25 and  $40 M_{\odot}$ .

**Figure 2.** Same as Fig. 1 for  $Z = 0.001$ . The  $1.0 M_{\odot}$  post He flash track has been omitted for clarity.

# Why Pop III binaries become 30Msun BH-BH

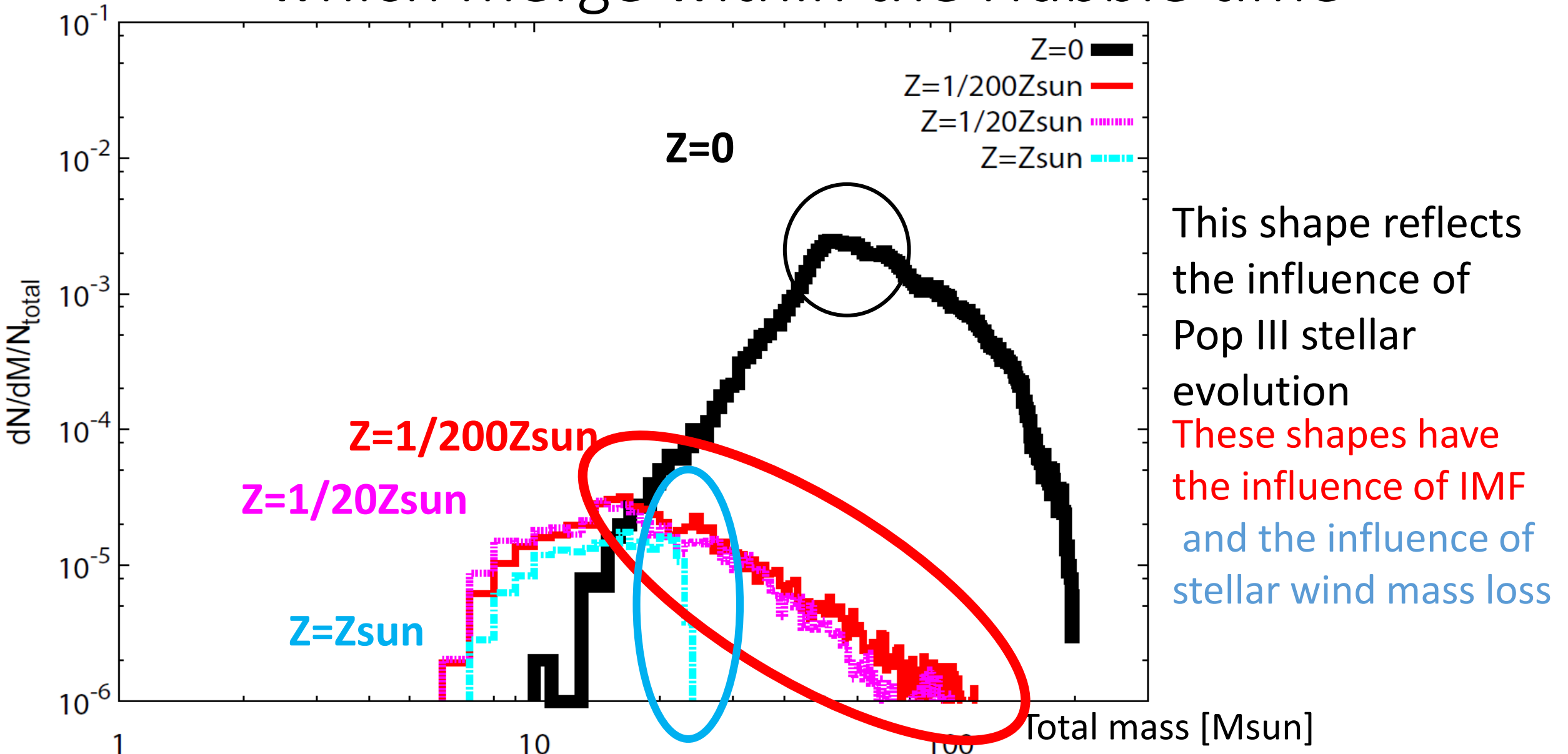


Marigo et al. 2001

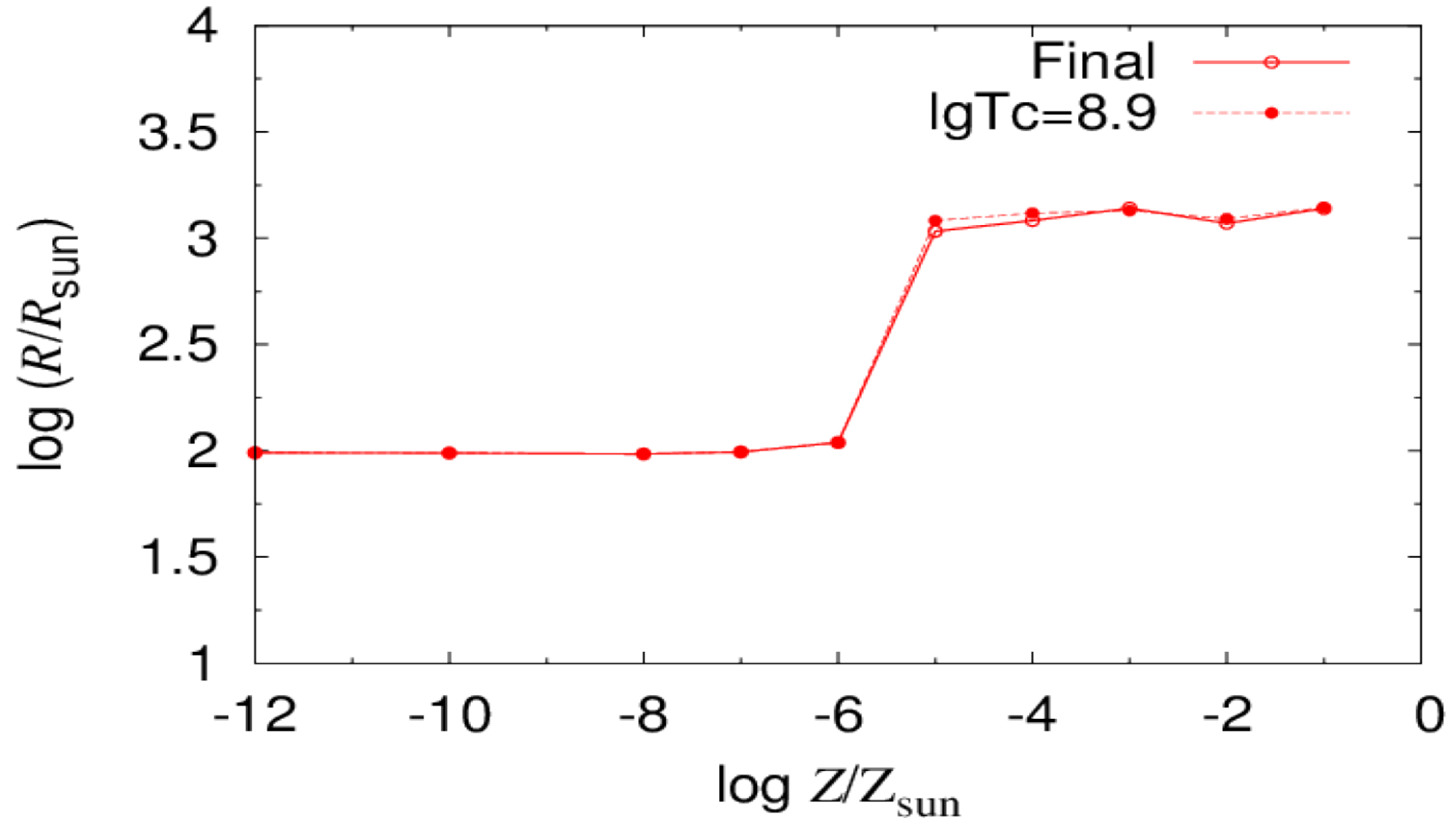
- $M > 50 M_{\text{sun}}$  **red giant**
  - Mass transfer is unstable
  - **common envelope**
  - **$1/3 \sim 1/2$  of initial mass**  
**( $\sim 25\text{-}30 M_{\text{sun}}$ )**
- $M < 50 M_{\text{sun}}$  **blue giant**
  - Mass transfer is stable
  - mass loss is not so effective
  - **$2/3 \sim 1$  of initial mass** ( $25\text{-}30 M_{\text{sun}}$ )



# Total mass distribution of BBH which merge within the Hubble time



# Evolution Transition from Pop III to Pop II



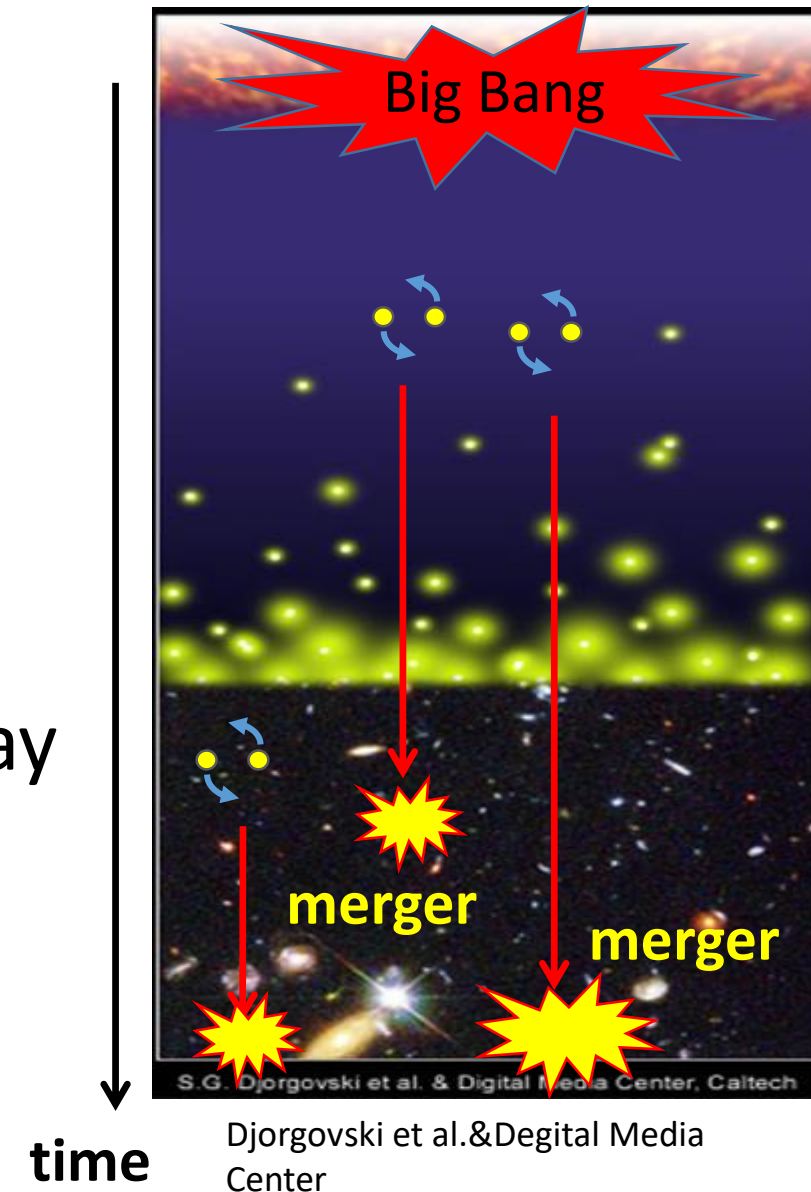
Maximum radius of 30Msun (Tanikawa, Yoshida, TK et al. 2019)

# Pop III BBH remnants for gravitational wave

- Pop III stars were born and died at  $z \sim 10$  ( $\sim 13.3$  Gyrs ago).
- The typical merger time of compact binaries  $\sim 10^{8-10}$  yr
- $dN/dt \propto t^{-1}$  (Kinugawa et al. 2014, Inayoshi et al. 2017)
- We might see Pop III BBH at the present day.
- Predicted Pop III merger rate at the present day  $\sim 30$  /yr/Gpc<sup>3</sup> (Kinugawa et al. 2014, 2016)

BBH merger rate estimated by LIGO

$$9.7-101 \text{ [yr}^{-1} \text{ Gpc}^{-3}] \quad (1811.12907)$$



# Pop III BBH?

ASTROPHYSICAL IMPLICATIONS

2014, D.

On the possibility that BBH binaries formed from the first generation of stars (Pop III) could be observed as gravitational wave sources. The predicted merger times for these binaries are astonishingly long.

merger times to occur in the nearby universe (Kinugawa et al. 2014). This is in contrast to the predicted mass properties



BLACK-HOLE MERGER GW150914

It has been proposed that the case of stellar black-hole binaries formed from the first generation of stars (Pop III) stars; see (Kinugawa et al. 2014). The predictions for these binaries are, since we have no direct observations of first-generation black holes, that the properties of these binaries will be different from binary black holes formed from the merger of two stellar black holes. It is estimated that a considerable number of these binaries will have total masses agree with the observed masses of the first black-hole merger GW150914. They can have sufficiently large masses to be observed as gravitational wave sources.



# However....

After GW150914, there are 1 bad news and 1 objection for Pop III BBH scenario

## 1.Bad news

~**1/3 decreasing** expected **Pop III SFR**

Because of constraints by Planck  $\tau_e$

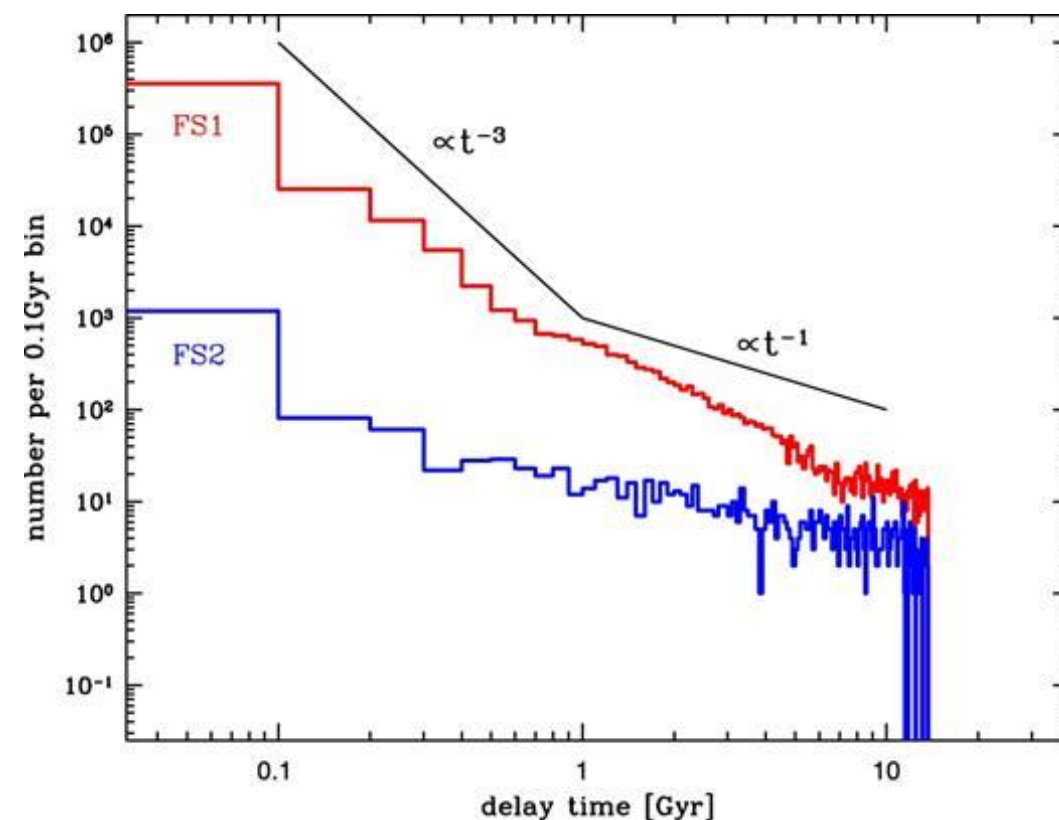
(Visbal et al.2015, Hartwig et al.2016, Inayoshi et al.2016)

## 2.Objection

Chris Belczynski tried to calculate  
Pop III BBH merger rate.

In his calculation, almost all Pop III  
BBHs ***merge at the early universe***

Belczynski et al. 2017



# The star formation rate of Pop III

In order to calculate merger rate,  
we need to know

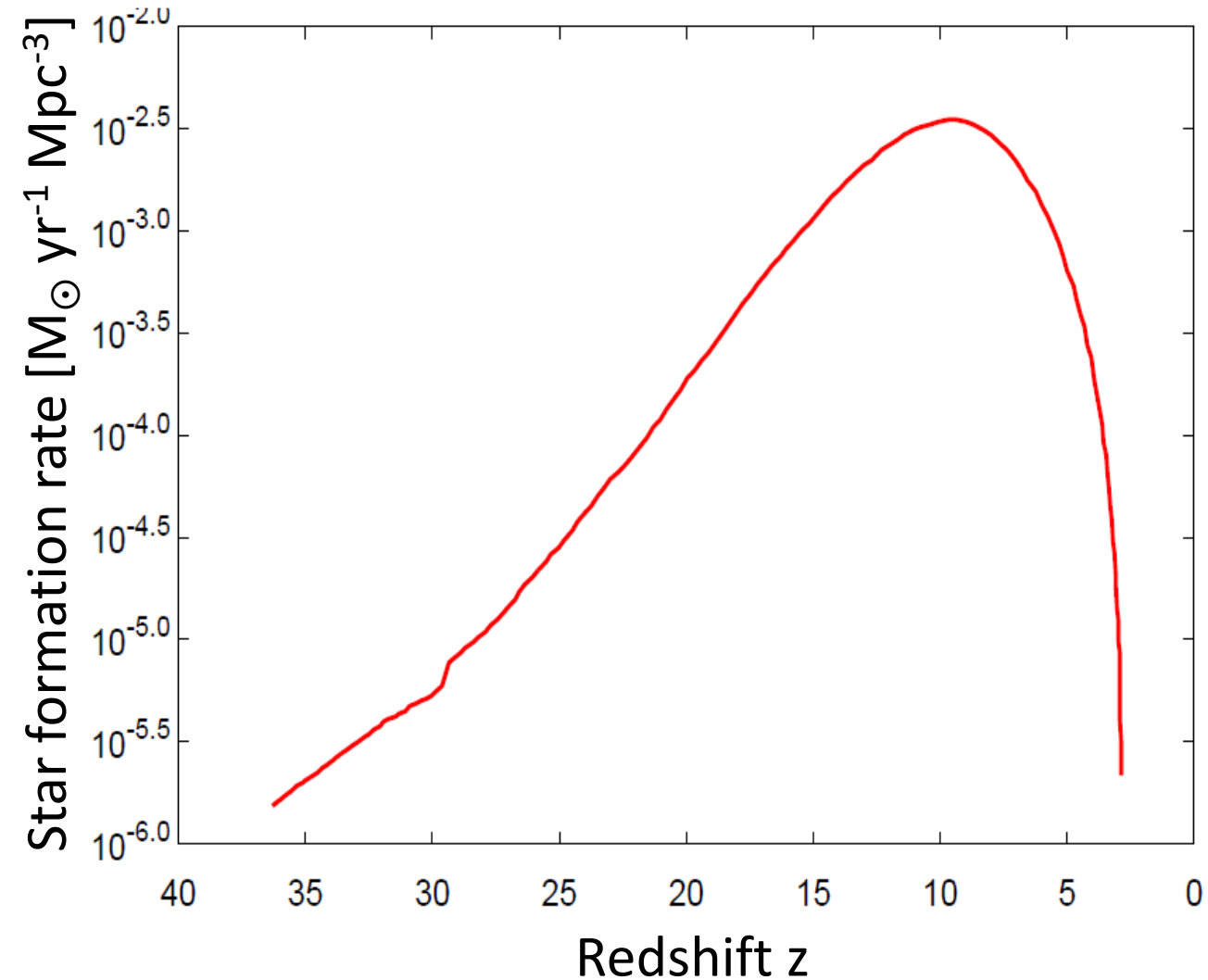
- When were Pop III stars born?
- How many were Pop III stars born?

⇒ Star formation rate

We adopt the Pop III SFR

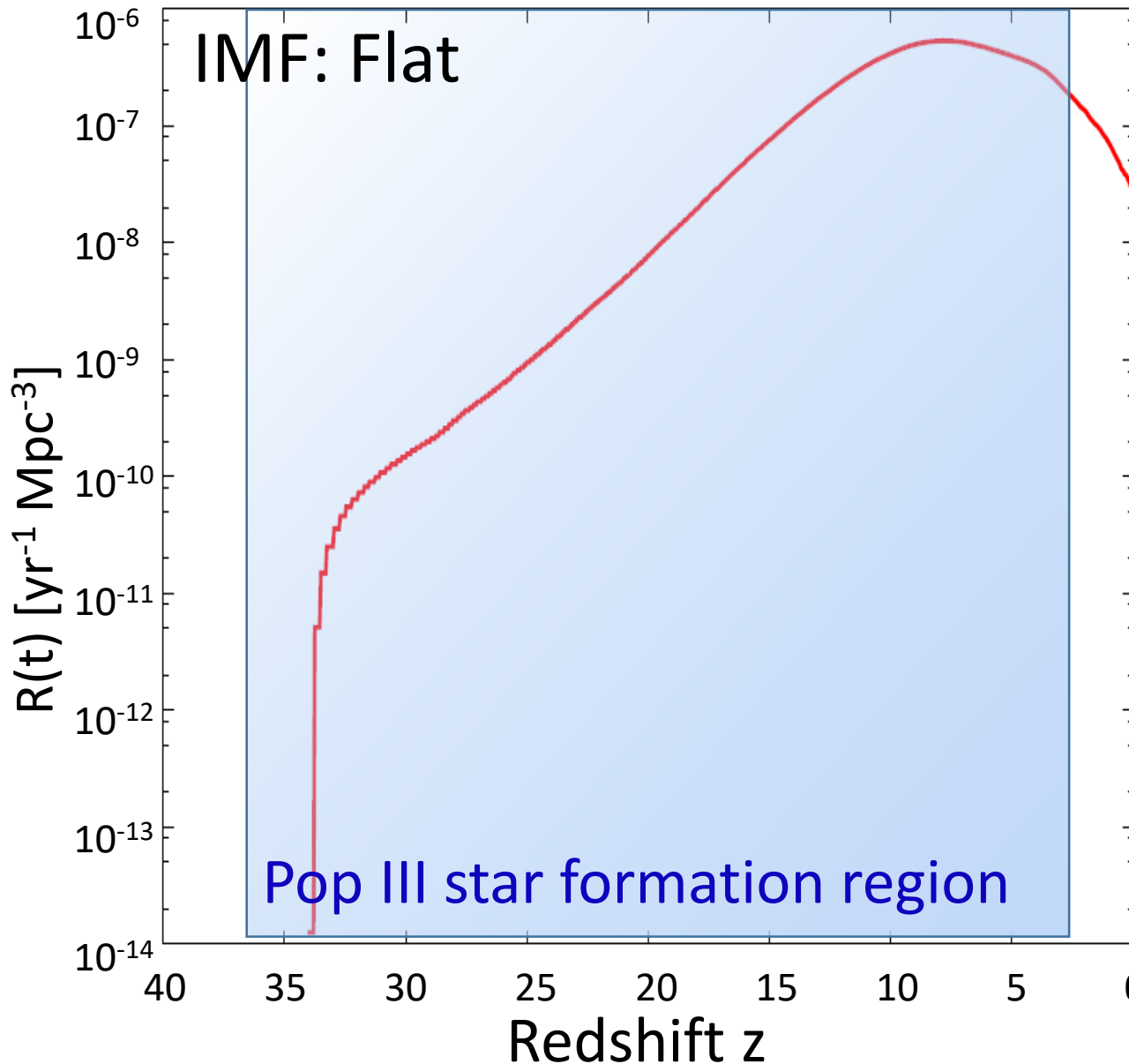
by de Souza et al. 2011

$$SFR_{peak} \sim 10^{-2.5} [M_{\odot} \text{ yr}^{-1} \text{ Mpc}^{-3}]$$
$$\rightarrow 10^{-3}$$



(de Souza et al. 2011)

# The Pop III BH-BH merger rate density



Pop III BHBH merger rate at the present day  
In our standard model

(Kinugawa et al. 2014,2016)

**$R \sim 30$  [ $\text{yr}^{-1} \text{Gpc}^{-3}$ ]**

**(Kinugawa et al. 2014,2016)**

**SFR decreased factor 1/3**

**$\rightarrow \sim 10$  [ $\text{yr}^{-1} \text{Gpc}^{-3}$ ]**

BBH merger rate estimated by LIGO

**$R = 9.7-101$  [ $\text{yr}^{-1} \text{Gpc}^{-3}$ ]** (1811.12907)

**Merger rate of massive BBH ( $\sim 30 M_{\text{sun}}$ )**

**$R \sim \text{several}$  [ $\text{yr}^{-1} \text{Gpc}^{-3}$ ]** (1811.12940)

# However....

After GW150914, there are 1 bad news and 1 objection for Pop III BBH scenario

## 1.Bad news

$\lesssim 1/3$  decreasing expected Pop III SFR

Because of constraints by Planck  $\tau_e$

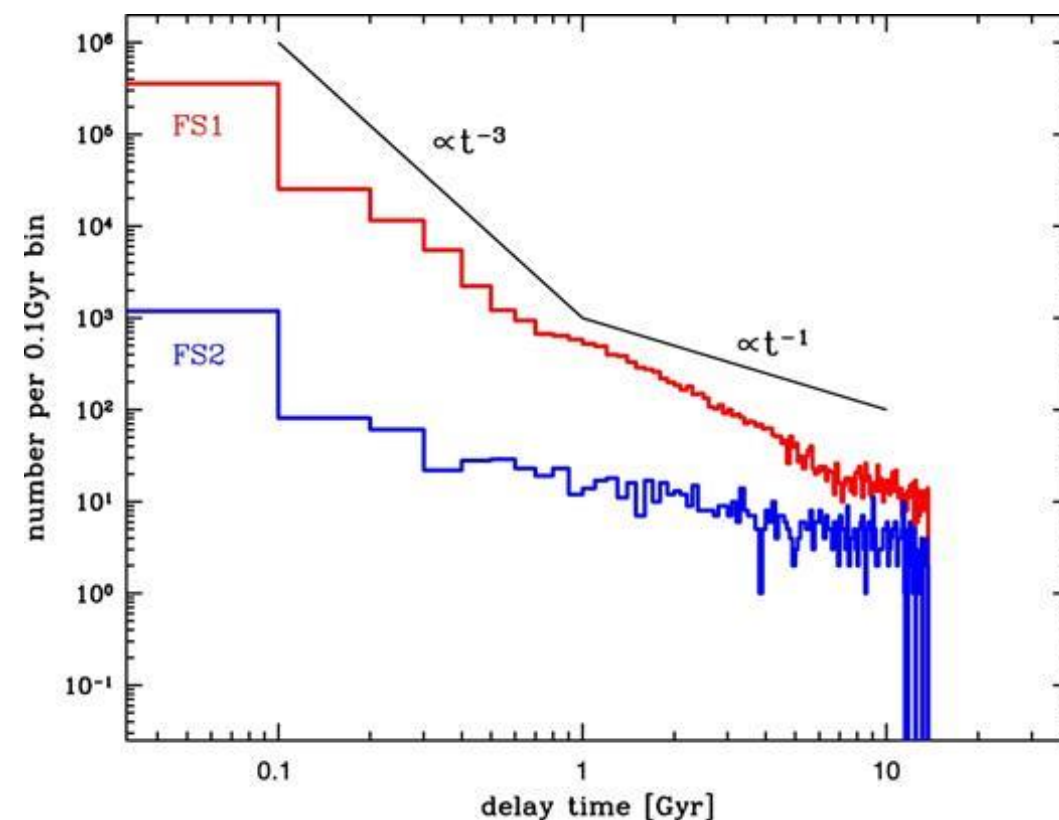
(Visbal et al.2015, Hartwig et al.2016, Inayoshi et al.2016)

## 2.Objection

Chris Belczynski tried to calculate Pop III BBH merger rate.

In his calculation, almost all Pop III BBHs *merge at the early universe*

Belczynski et al. 2017

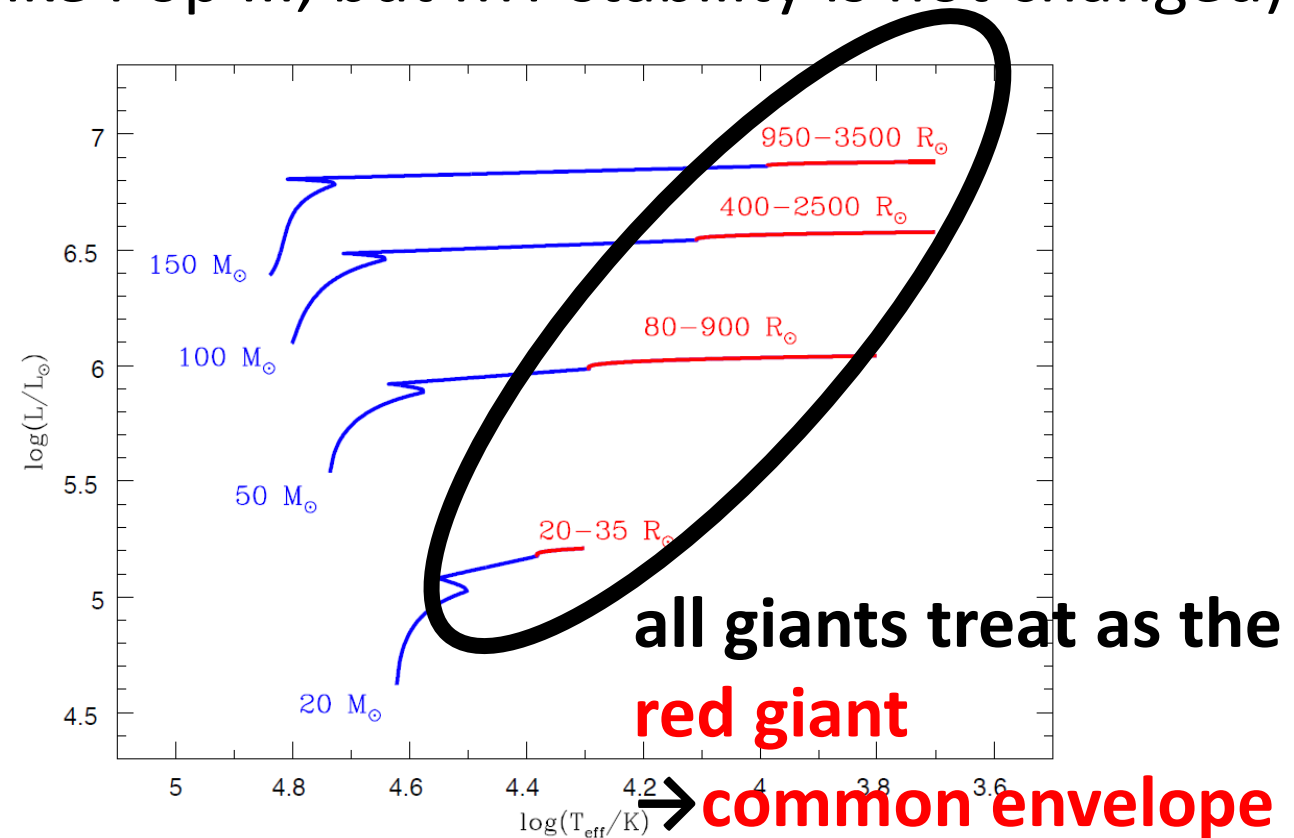
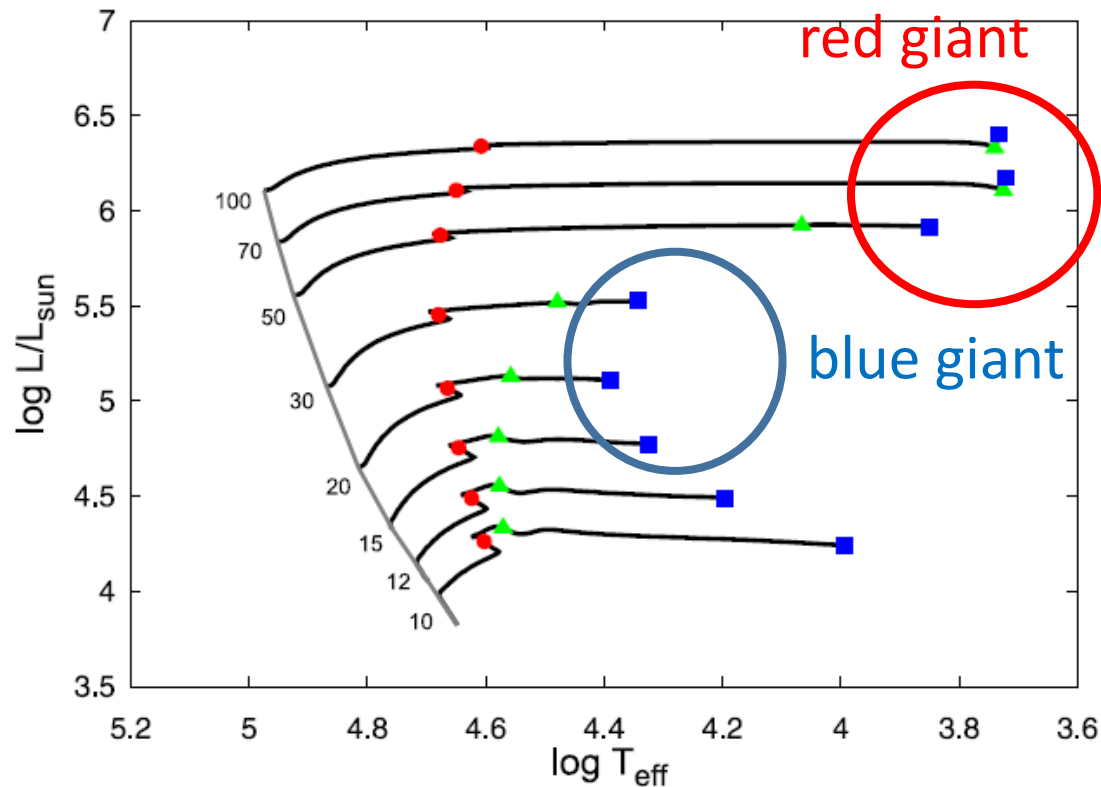




# Difference between K14 and Belczynski's Pop III calc.

- Kinugawa 2014: use Pop III stellar evolution model (Marigo et al.2001)
- Belczynski 2017: use modified  $Z=0.005Z_{\text{sun}}$  model.

(HR and radius evolution is changed like Pop II, but MT stability is not changed)



# However....

After GW150914, there are 1 bad news and 1 objection for Pop III BBH scenario

## 1.Bad news

$\lesssim 1/3$  decreasing expected Pop III SFR

Because of constraints by Planck  $\tau_e$

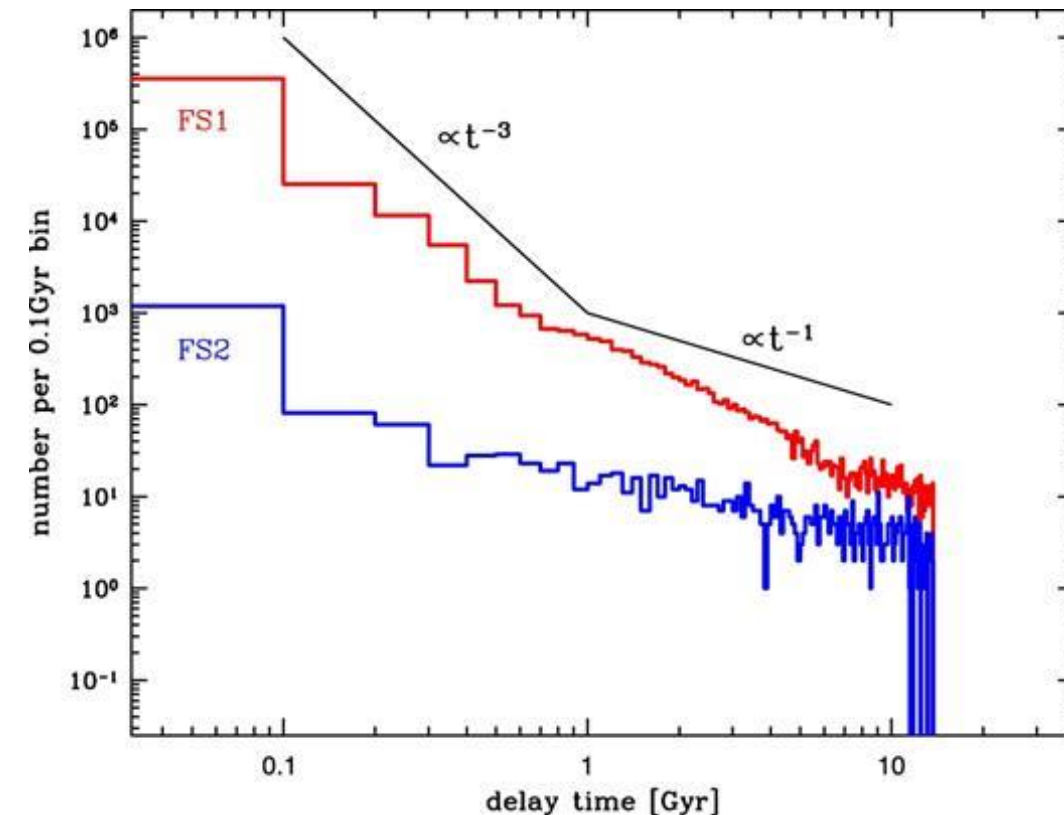
(Visbal et al.2015, Hartwig et al.2016, Inayoshi et al.2016)

## 2.Objection

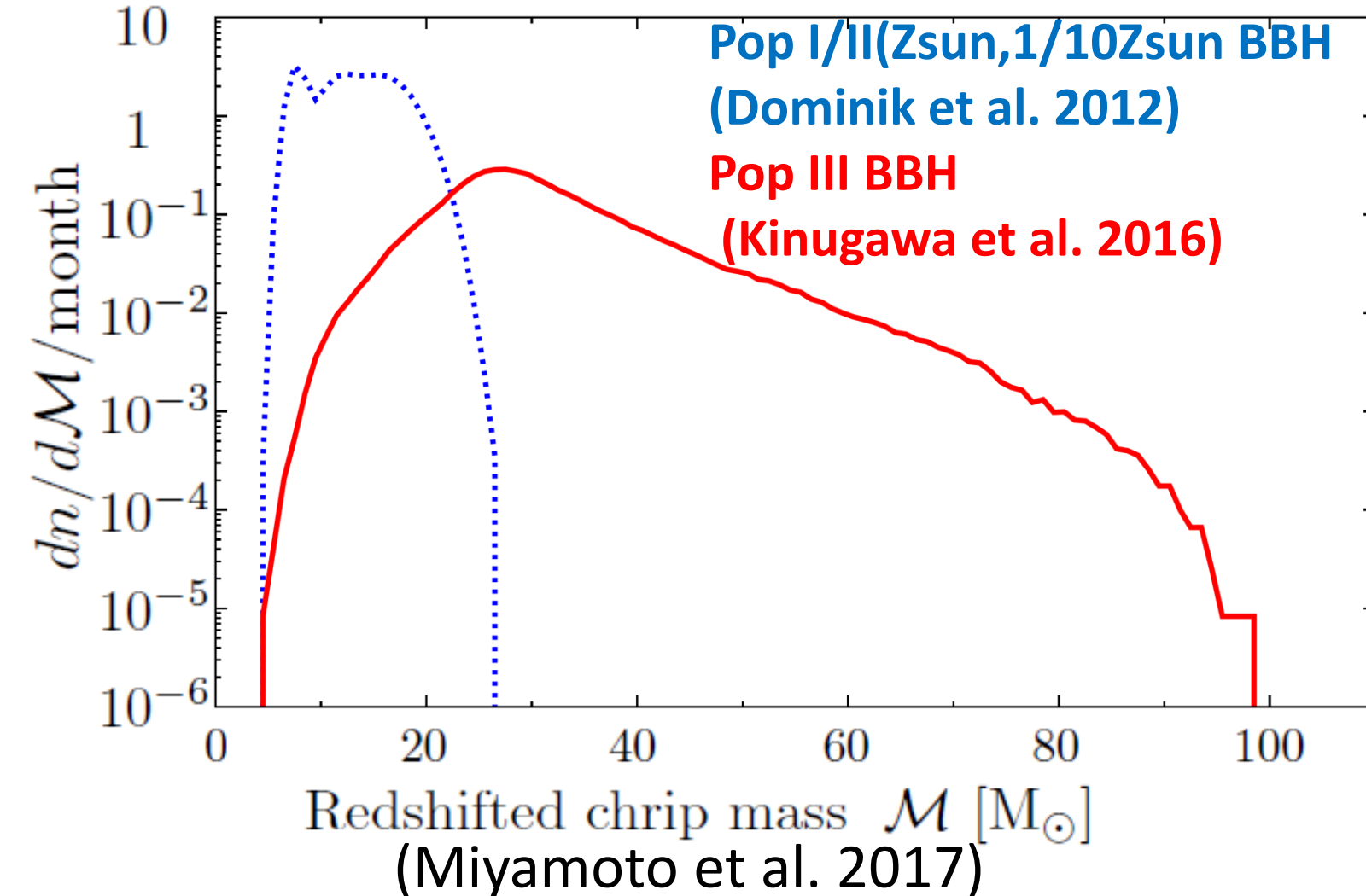
Cryzs Belczynski tried to calculate Pop III BBH merger rate.

In his calculation, almost all Pop III BBHs *merge at the early universe*

Belczynski et al. 2017



# Mass distributions of observable BBHs (KAGRA)



$$\mathcal{M} = (1+z) \frac{(M_1 M_2)^{3/5}}{(M_1 + M_2)^{1/5}}$$

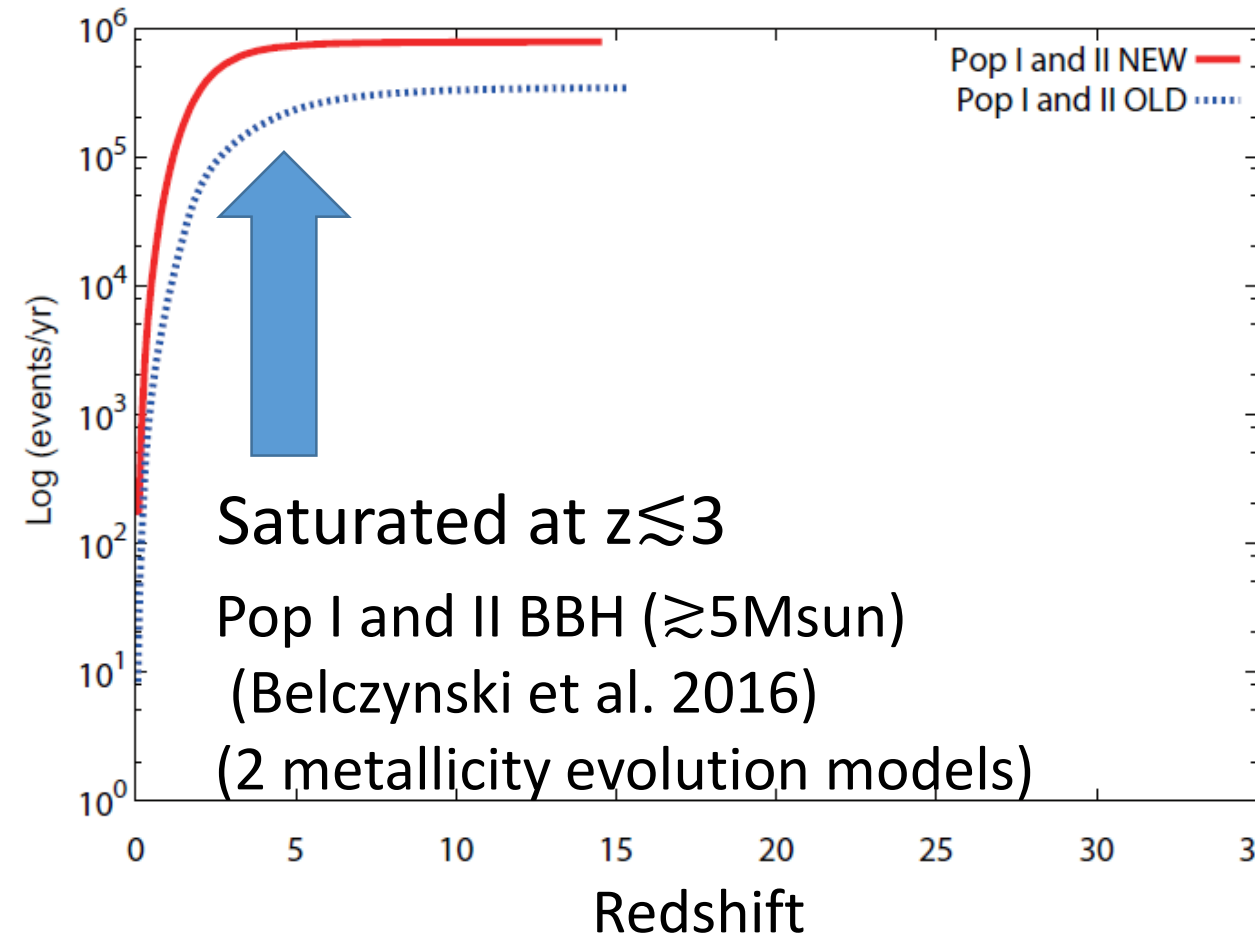
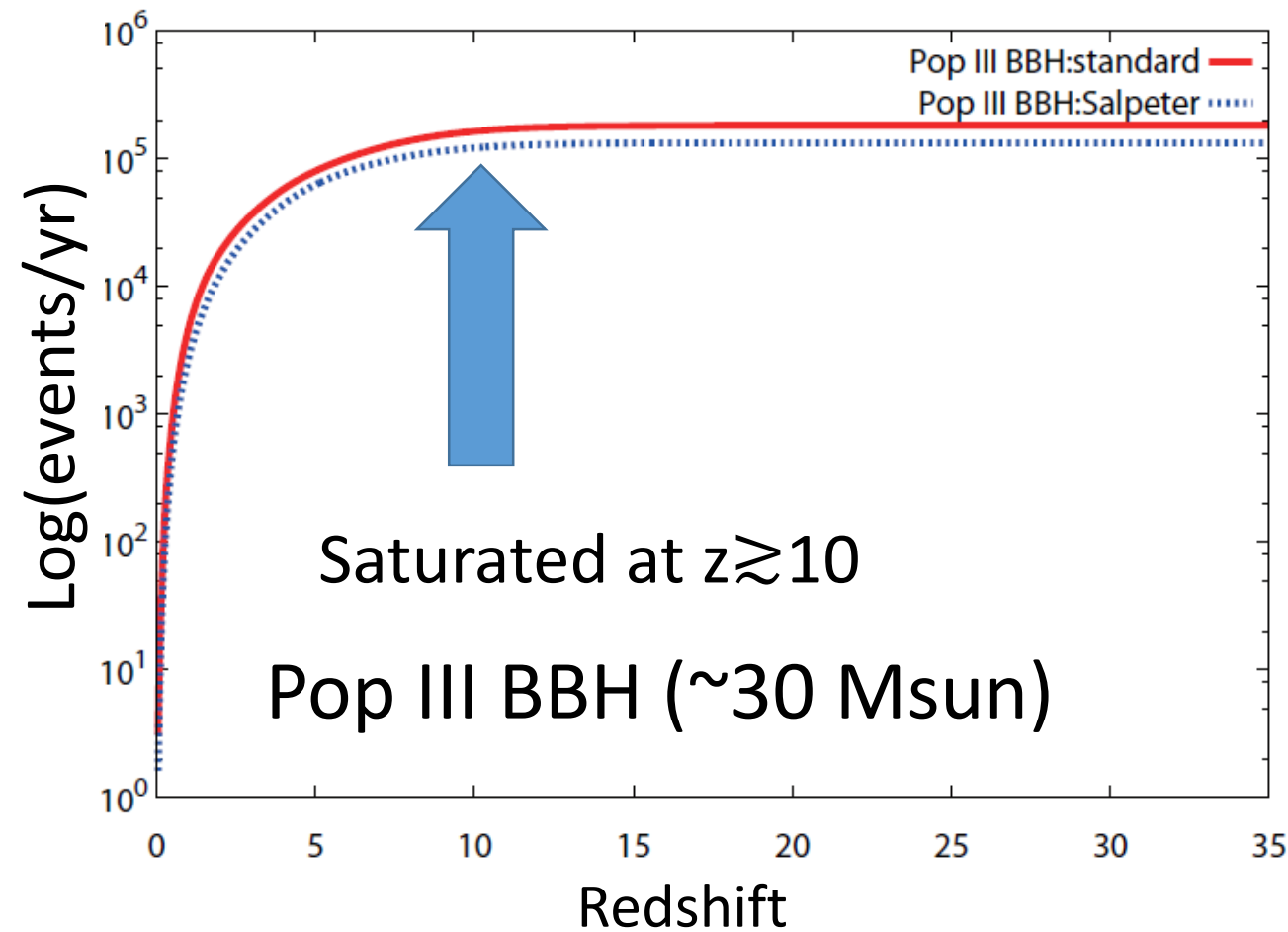
- The mass distribution might distinguish Pop III from Pop I, Pop II

→ The evidence of Pop III

Even if Mass dist. cannot distinguish

→ redshift dependence

# Cumulative BBH merger rate



# Merger time dependence of Pop III BBH spin

## Preliminary results

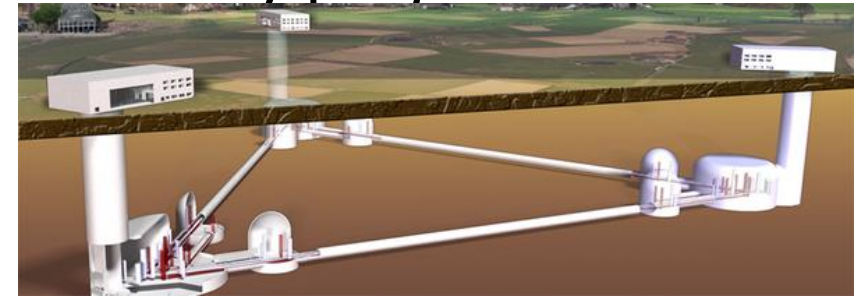
	$a_1/M_1 < 0.1$ $a_2/M_2 < 0.1$	$a_1/M_1 < 0.1$ $a_2/M_2 > 0.9$	$a_1/M_1 > 0.9$ $a_2/M_2 < 0.1$	$a_1/M_1 > 0.9$ $a_2/M_2 > 0.9$
Merger time <1Gyr	25%	36%	0%	23%
Merger time >10Gyr	70%	0.3%	4%	0%

- If the origin of massive BBHs is Pop III,  
high spin BBHs are easier to be detected at high redshift

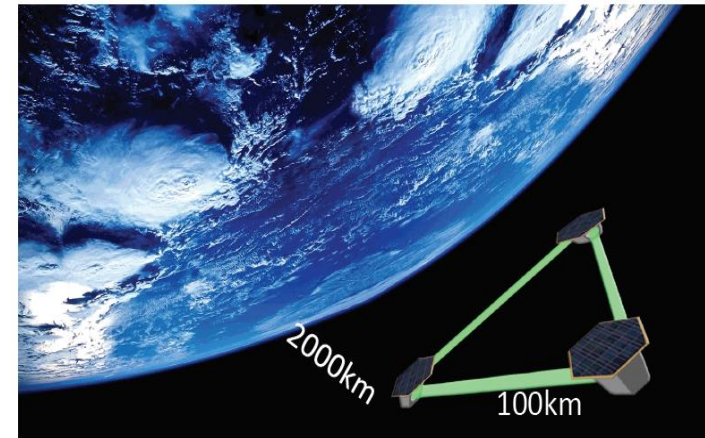
# Future plan of GW observer : ET, CE, B-DECIGO and DECIGO

- Einstein telescope (ET): the next generation GW observatory of Europe
- Cosmic explorer (CE) : the next generation GW observatory of US.
- DECIGO: Japanese space gravitational wave observatory project
- B-DECIGO: test version of DECIGO

ET, CE, B-DECIGO :  $z \sim 10$  (30  $M_{\text{sun}}$  BH-BH)

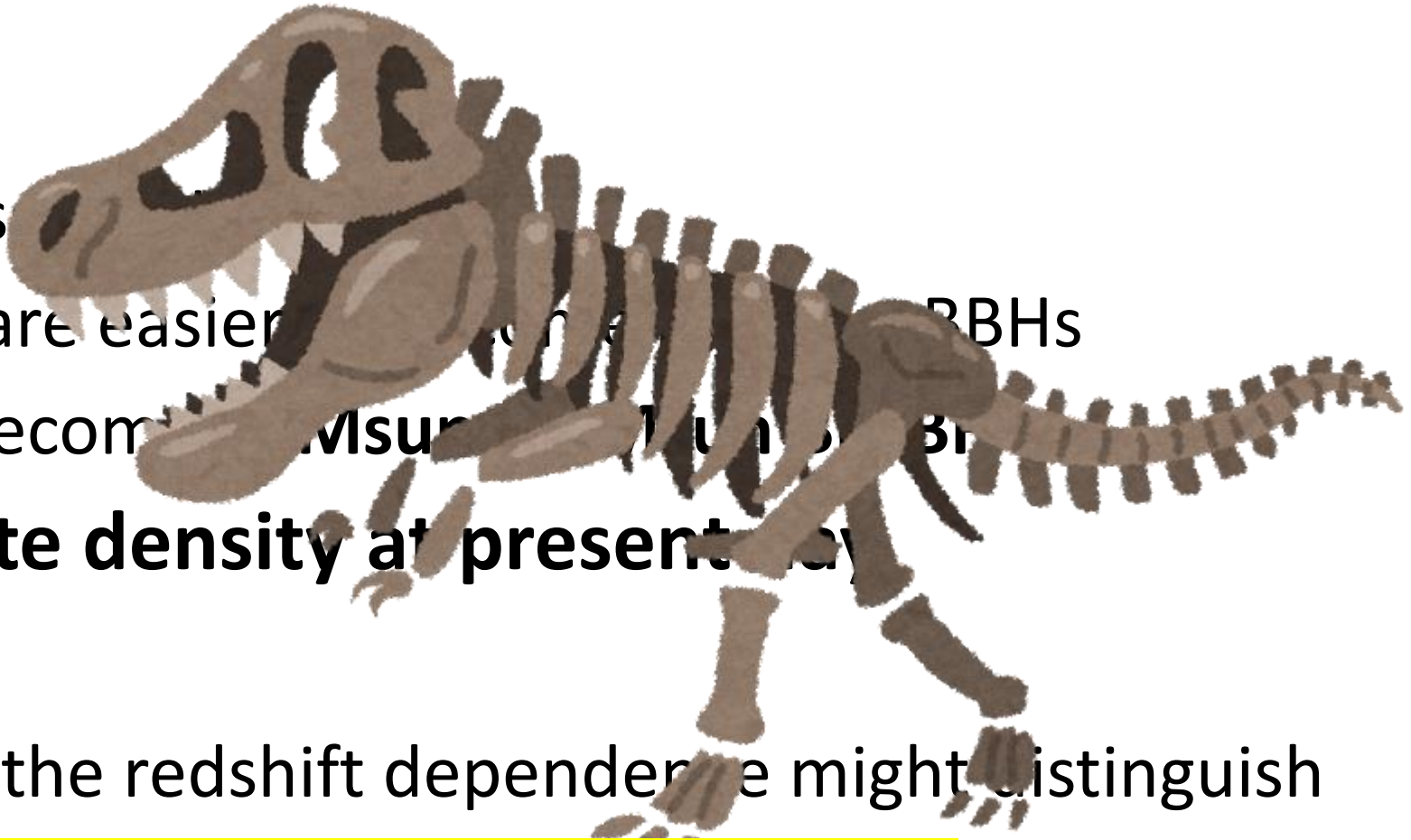


DECIGO can see Pop III BH-BHs  
when Pop III stars were born ( $z \sim 20$ )!  
(Nakamura, Ando, Kinugawa et al. 2016)





# PopIII BH-BH



- 7 out of 10 has
- Low masses are easier to find BBHs
- Pop III become visible
- **Pop III rate density at present day**
- **R** **]**
- The mass distribution or the redshift dependence might distinguish Pop III from
- **Massive BBHs = the fossil of Pop III ?**
- DECIGO can see Pop III BH-BH mergers when they were born

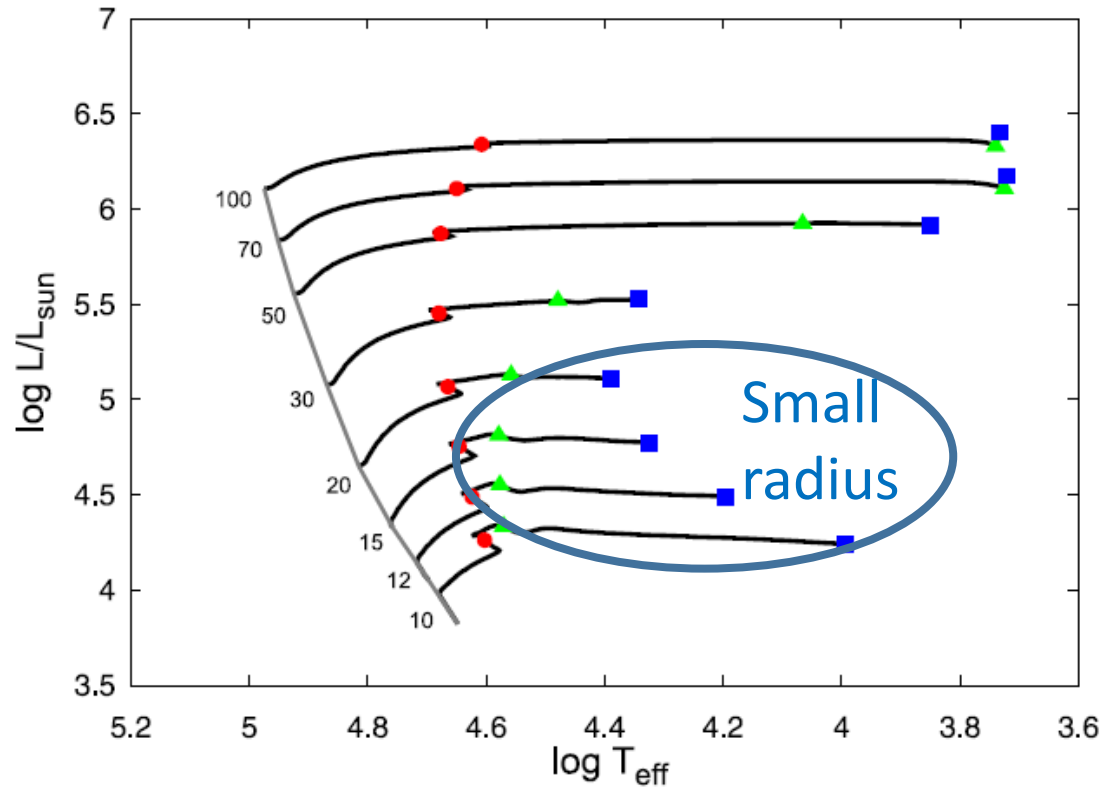
# Other Pop III compact binaries cases

- Pop III NSNS

Almost all binary NS (maybe) disrupt

- Pop III NSBH

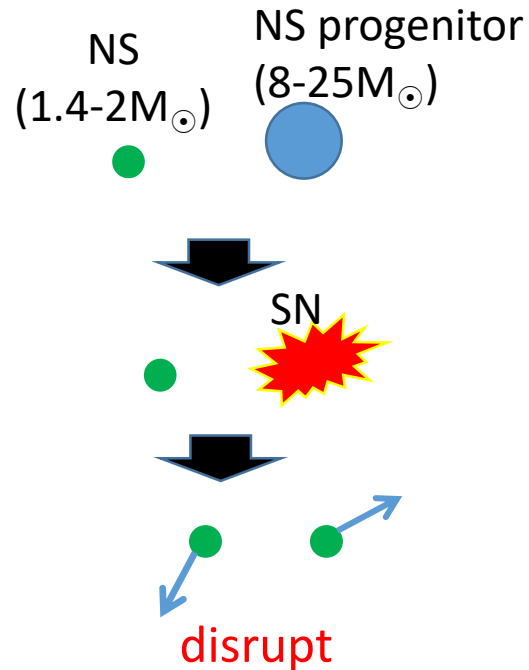
# Pop III NS progenitor evolution



- blue giant
  - Mass transfer is stable
  - mass loss is not so effective before supernova

# Pop III NS-NS disrupt

For example, we consider NS and NS progenitor binary.



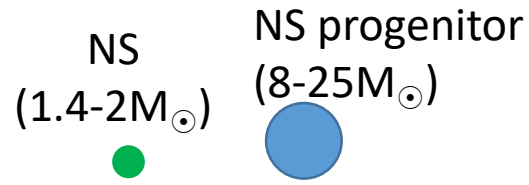
In the case of Pop III NS progenitor, wind mass loss and the mass loss due to binary interaction is not effective. When NS progenitor becomes supernova, NS progenitor suddenly loses mass and becomes NS.

Then, due to instant mass loss the binding energy of binary decreases and binary NS disrupts.

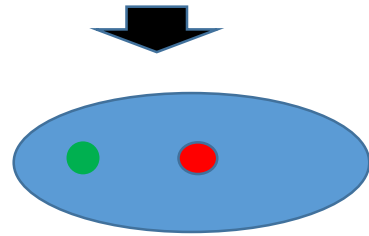
➡ Binary NS **cannot survive!**

# Pop I and II NS-NS

For example, we consider NS and NS progenitor binary.



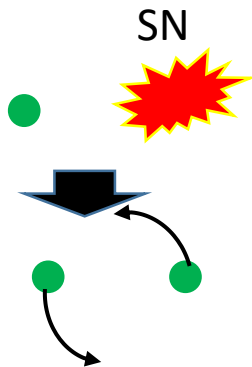
In the case of Pop I and II NS progenitor, wind mass loss and the mass loss due to binary interaction is effective.



NS progenitor can lose mass before SN.



Binary NS **can survive!**



# Other Pop III compact binaries cases

- Pop III NSNS

Almost all binary NS (maybe) disrupt

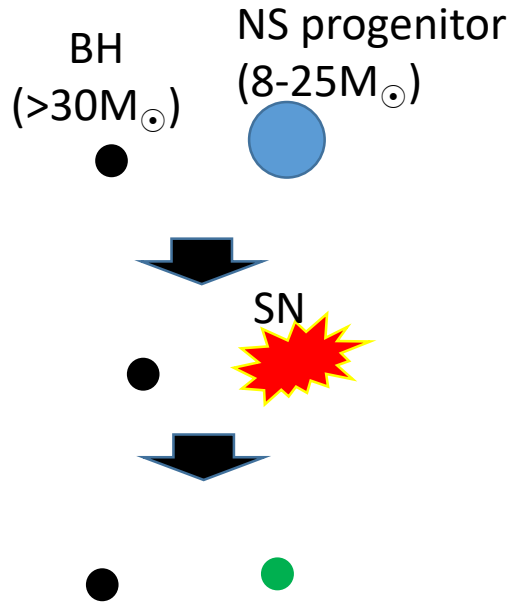
- Pop III NSBH

NSBH do not disrupt



# Pop III NS-BH do not disrupt

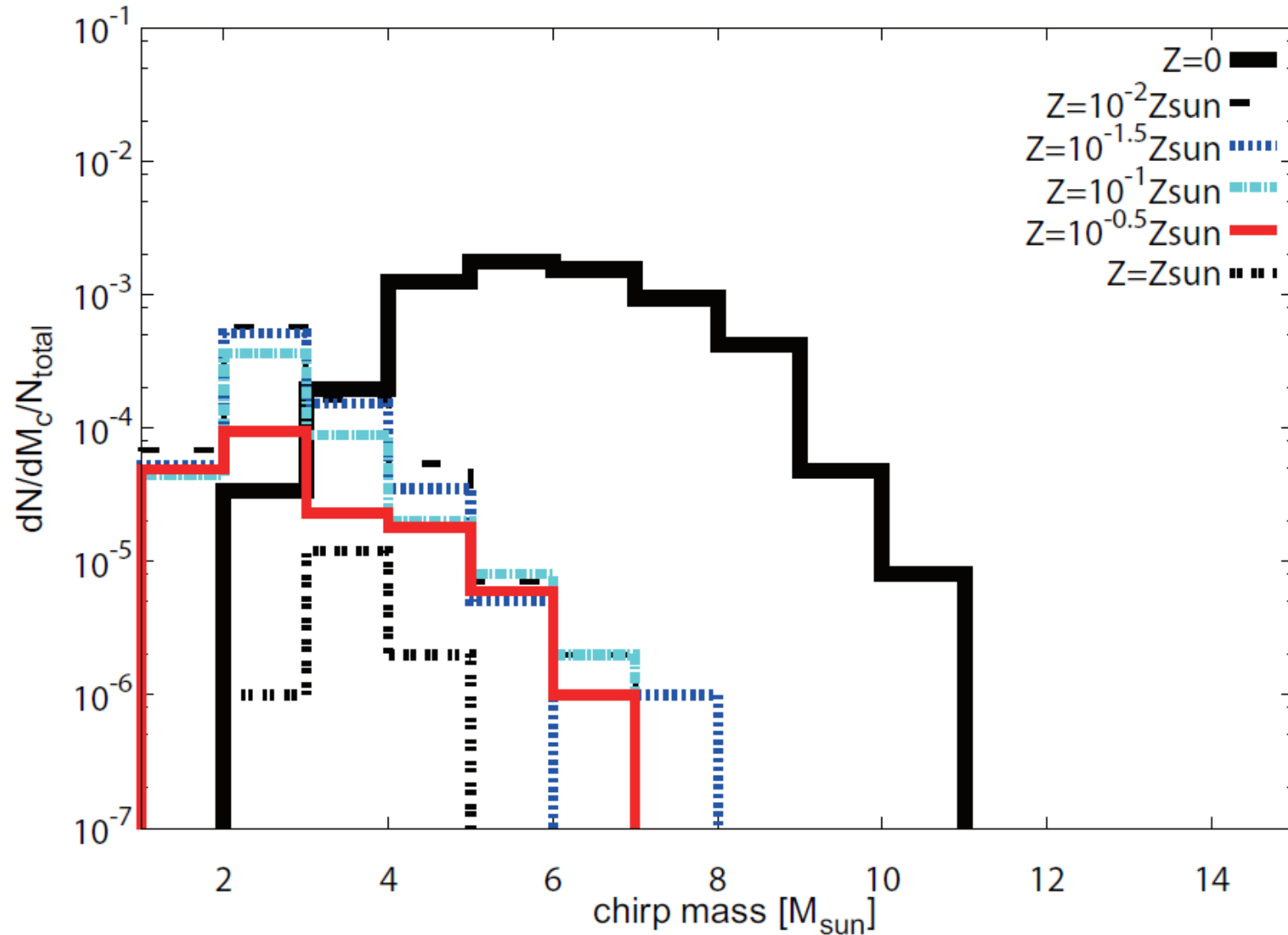
For example, we consider BH and NS progenitor binary.



In the case of Pop III NS progenitor, wind mass loss and the mass loss due to binary interaction is not effective. When NS progenitor becomes supernova, NS progenitor suddenly loses mass and becomes NS. But, due to massive BH, NS do not disrupts.

➡ NS BH **can survive!**

# Merging NSBH chirp mass distribution



# NSBH detection rate

	Merger rate [/yr/Gpc <sup>3</sup> ]	aLIGO (design sensitivity) detection rate [/yr]
Pop I+II	28.8 (Belczynski et al. 2016)	~10
Pop III	1.25	5.24(*)

\*For simplicity, as the assumption of the chirp mass of Pop III NSBH,  
we fixed  $M_c = 6M_{\odot}$  (Kinugawa et al.2016)

## Pop III GW summary

- **Merger rate of Pop III BBH at  $z \sim 0$  (GW150914 like massive BBH)**

$$R \sim 10 \text{ [yr}^{-1} \text{ Gpc}^{-3} \text{]}$$

- **Typical chirp mass**

$$M \sim 30 M_{\odot}$$

We might detect (detected?) Pop III BBHs by GW

- **Detection rate of Pop III NSBH for aLIGO designed sensitivity**

$$R \sim 5 \text{ [yr}^{-1} \text{]}$$

- **Typical chirp mass**

$$M \sim 6 M_{\odot} \text{ (1.4 Msun NS + 50 Msun BH)}$$