

# 相対論的衝撃波の Particle-in-Cellシミュレーション

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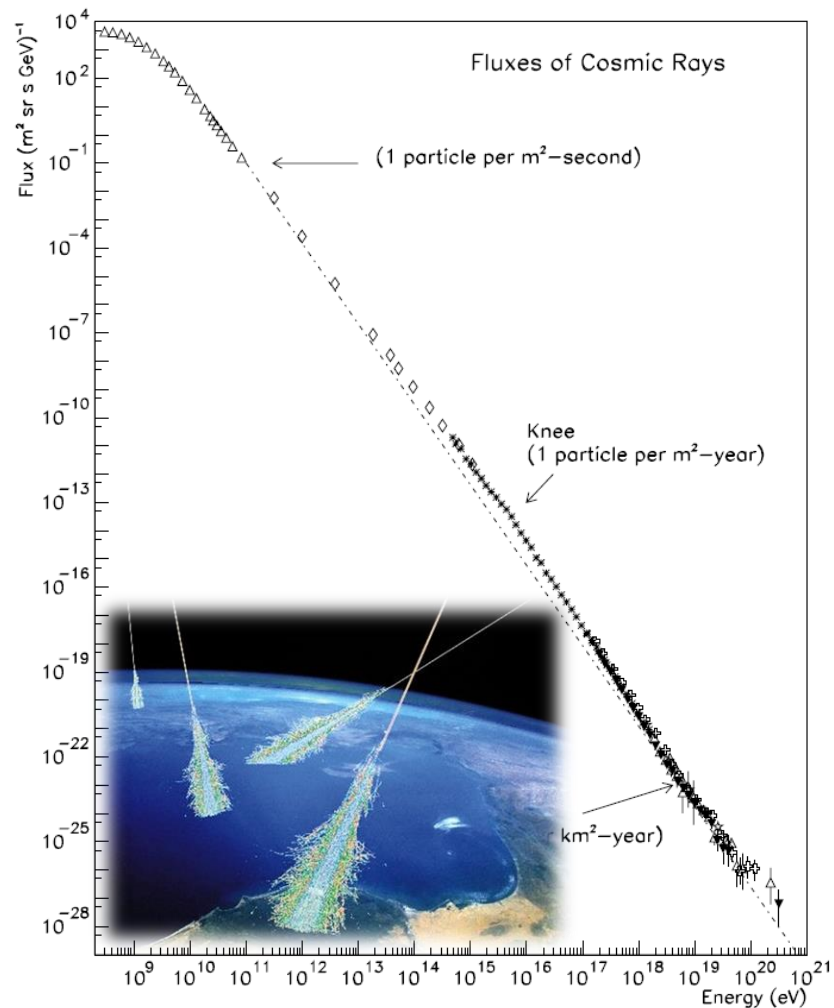
# 内容

1. 数値チェレンコフ不安定の抑制  
Ikeya & Matsumoto, PASJ, 2015
2. 相対論的衝撃波における大振幅先駆波励起と電子加速  
Iwamoto et al., ApJ, 2017; 2018
3. イオンワイベル不安定による磁場生成／飽和過程  
Takamoto, Matsumoto, Kato, ApJL, 2018

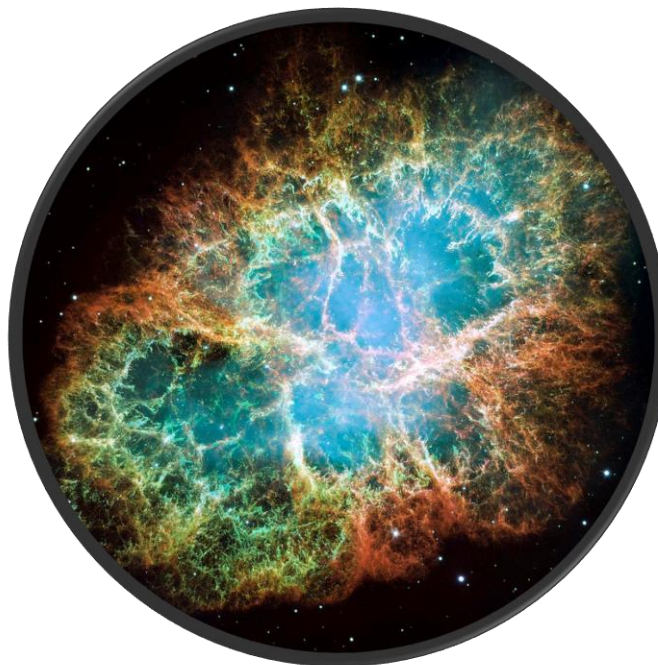
# 宇宙線の起源

地球に届く宇宙線のフラックス

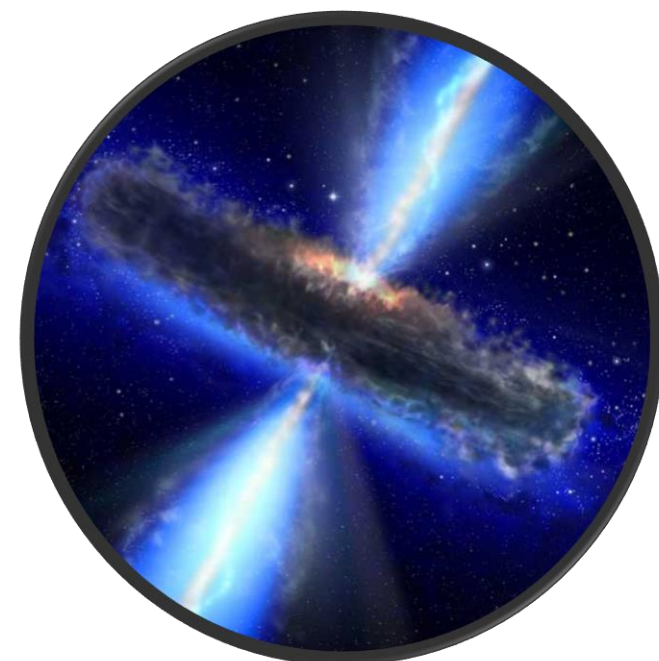
宇宙の爆発的現象に伴う**衝撃波**での荷電粒子（陽子・電子）の加速メカニズムを探求



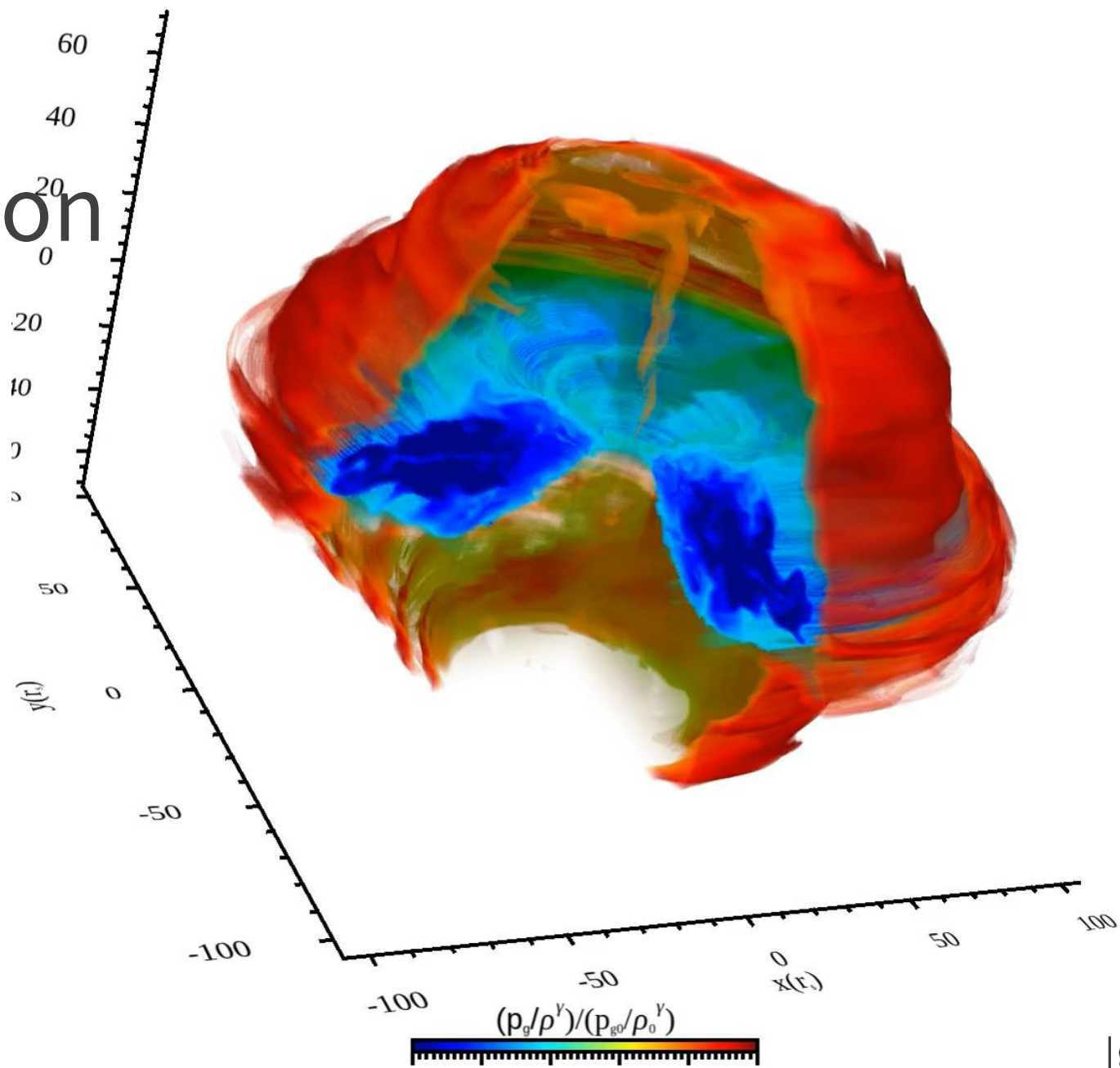
超新星爆発  
(銀河宇宙線)



宇宙ジェット  
(超高エネルギー宇宙線)



CANS+R  
(MHD + radiation  
transfer)



Igarashi+, in prep.

# Particle-in-Cell simulation

Vlasov eq. as particle motions

$$\frac{d\mathbf{x}_p}{dt} = \frac{\mathbf{u}_p}{\gamma_p}$$

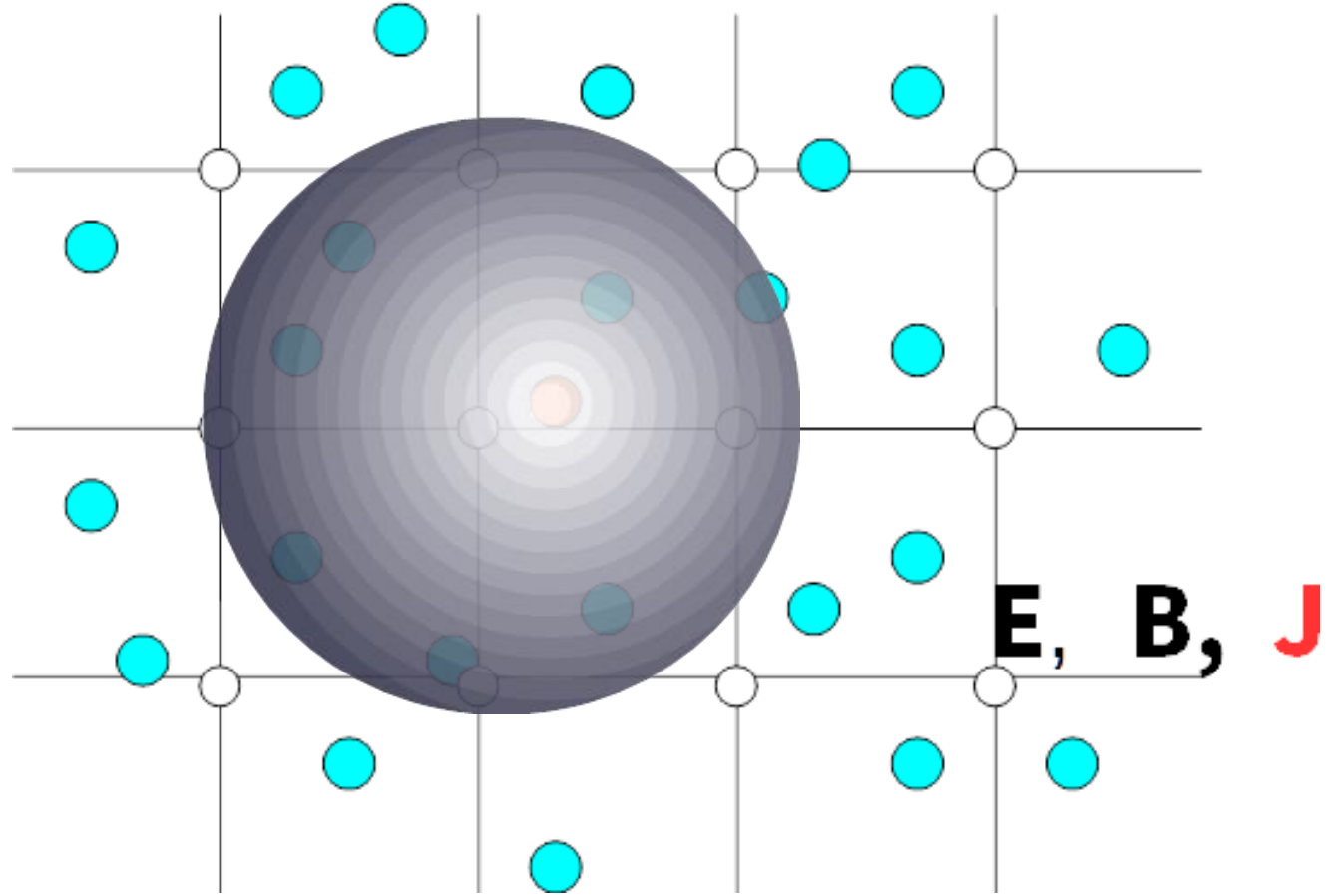
$$\frac{d\mathbf{u}_p}{dt} = \frac{q}{m} \left( \mathbf{E} + \frac{\mathbf{u}_p}{c\gamma_p} \times \mathbf{B} \right)$$

$$\mathbf{J} = \sum_p q_p \frac{\mathbf{u}_p}{\gamma_p}$$

Maxwell eqs. on grid points

$$\frac{\partial \mathbf{B}}{\partial t} = -c \nabla \times \mathbf{E}$$

$$\frac{\partial \mathbf{E}}{\partial t} = c \nabla \times \mathbf{B} - 4\pi \mathbf{J}$$



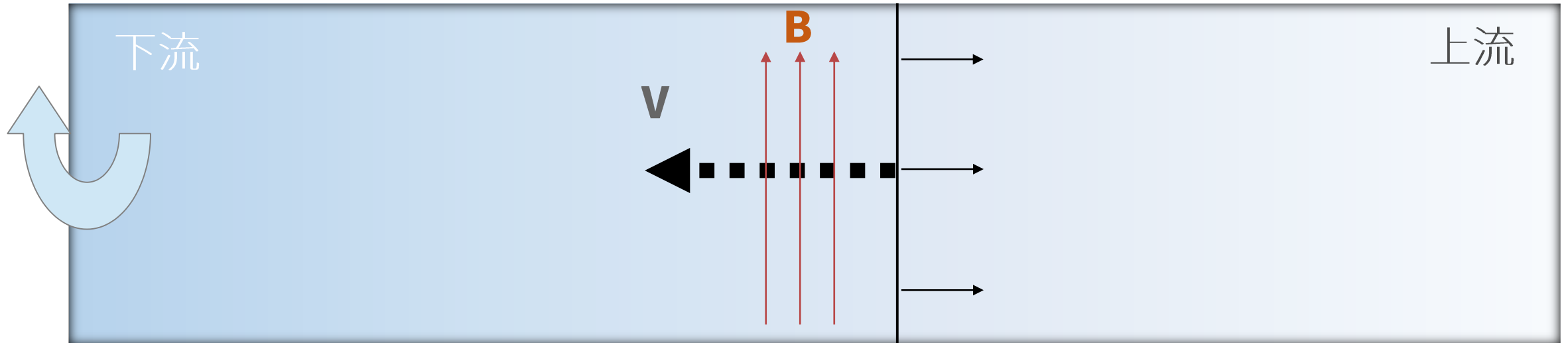
# スパコンで衝撃波を作る

- 自家製 2D/3D particle-in-cell コード
- SIMD最適化
- MPI+OpenMP ハイブリッド並列
- $10^5$  cores までスケール



**y**

衝撃波下流静止系 (波面は右に伝搬)



**x**

# PICの世界動向 (TRISTAN-PIC 対 その他)

$$\sigma = M_A^{-2}$$

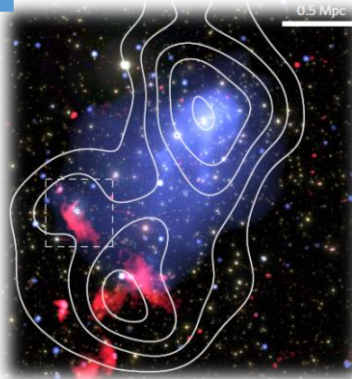
弱い衝撃波



強い衝撃波

1

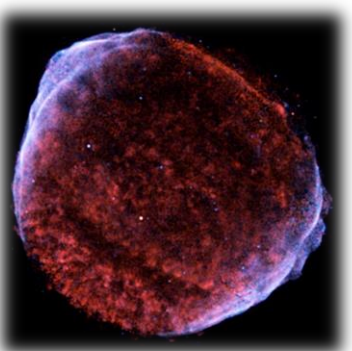
銀河団衝撃波



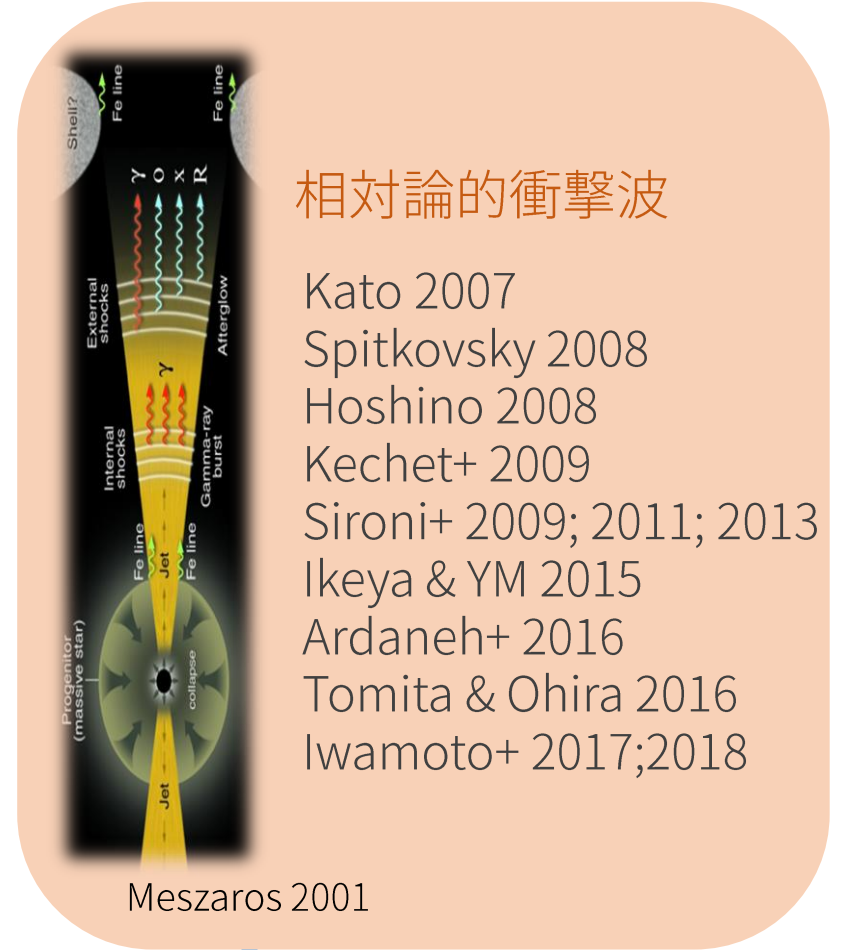
Weeren+ 2017

Matsukiyo+ 2011  
 Riquelme & Spitkovsky 2011  
 Guo+ 2014a,b  
 Matsukiyo & YM 2015  
 (太陽系衝撃波を入れると  
 もっとたくさん)

SNR衝撃波



Kato+ 2010; 2015  
 YM+ 2012;2013;2015;2017  
 Niemiec+ 2012; 2013  
 Park+ 2015  
 Wieland+ 2016  
 Bohdan+2017



相対論的衝撃波

Kato 2007  
 Spitkovsky 2008  
 Hoshino 2008  
 Kechet+ 2009  
 Sironi+ 2009; 2011; 2013  
 Ikeya & YM 2015  
 Ardaneh+ 2016  
 Tomita & Ohira 2016  
 Iwamoto+ 2017;2018

Meszáros 2001

0

非相対論的



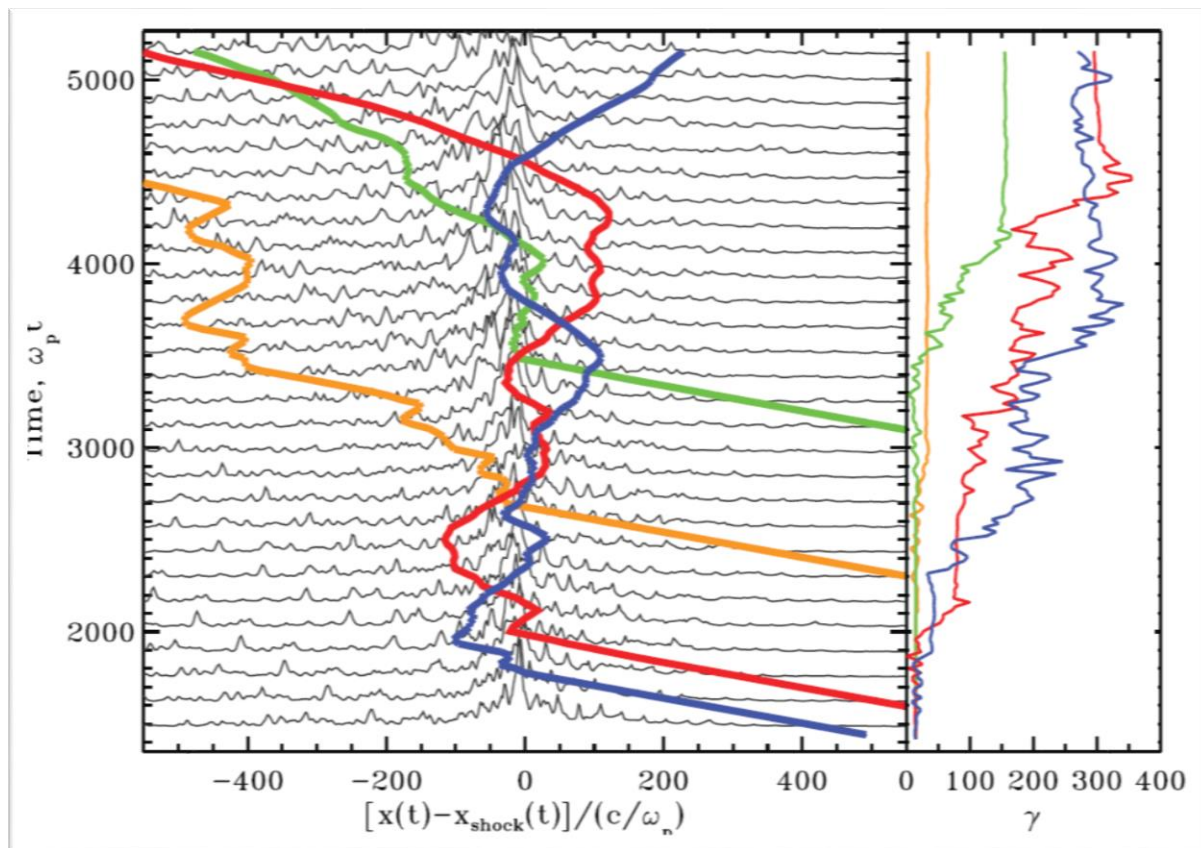
相対論的

1

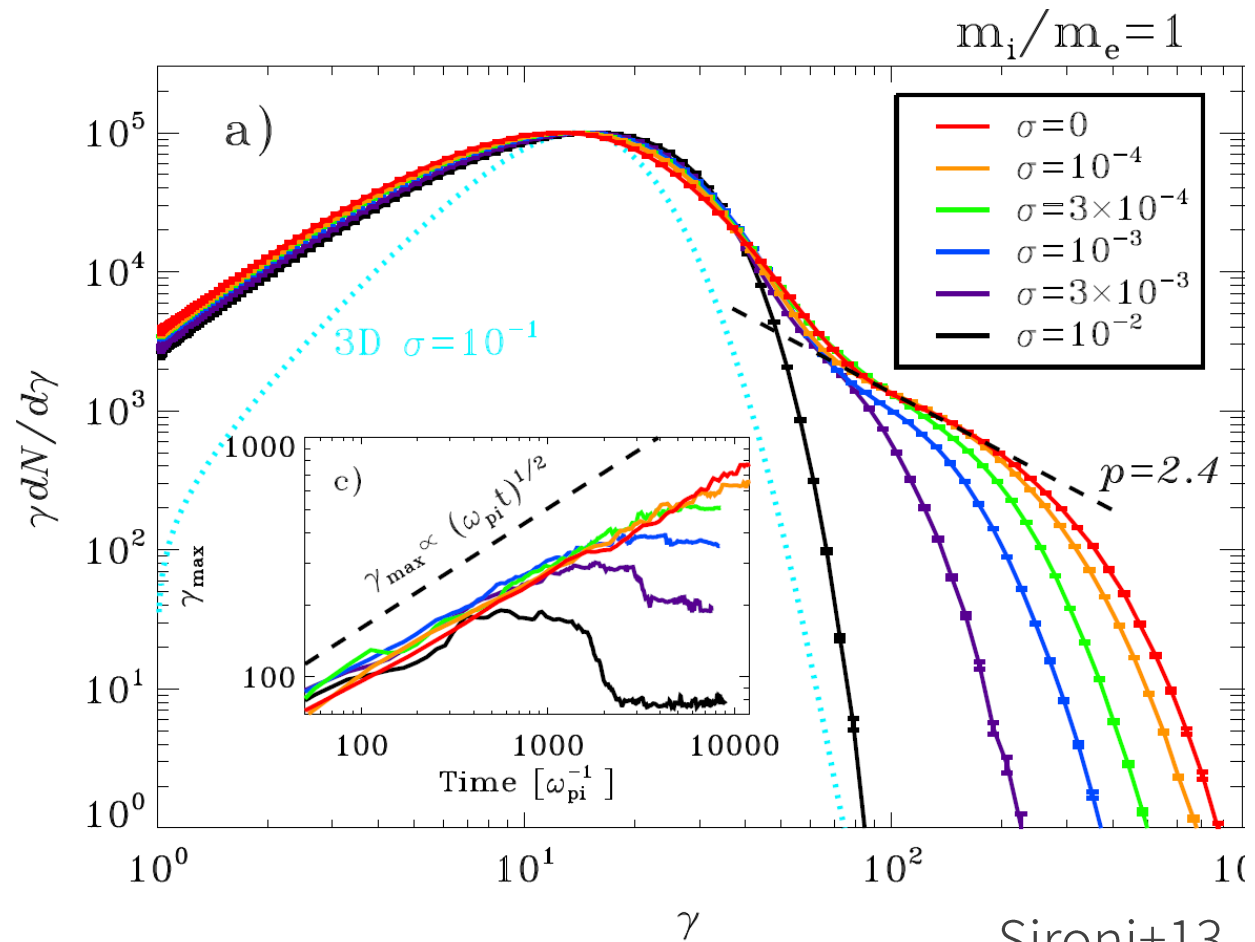
$V_{shock}/c$

# 相対論的衝撃波のPICシミュレーションと粒子加速

Weibel-mediated relativistic shock & Fermi acceleration



Kato07; Spitkovsky08



Sironi+13



数値チェレンコフ不安定の抑制

# 数値チェレンコフ不安定 (Godfrey74 JCP)

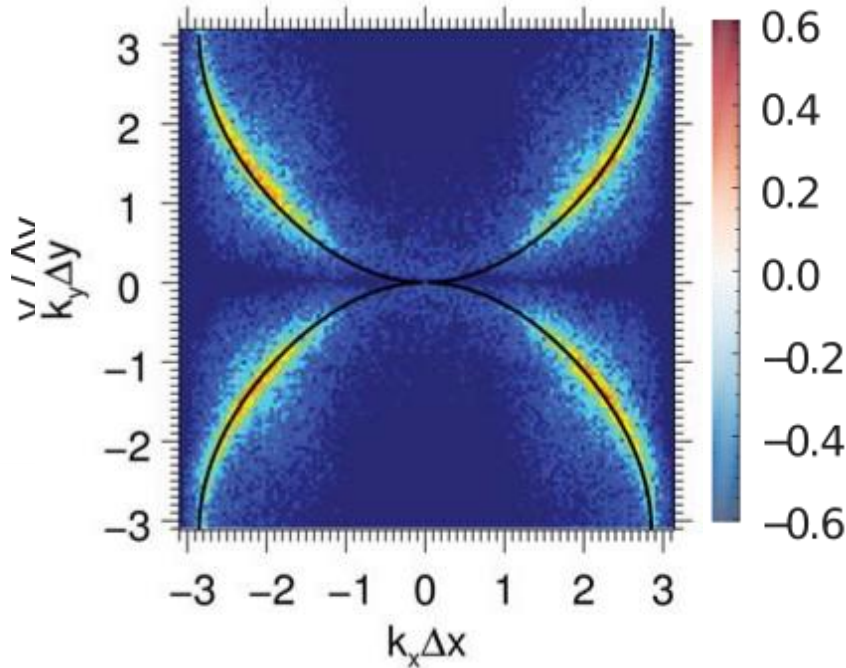
離散化した空間での光モードの分散関係

$$\sin^2\left(\frac{\omega\Delta t}{2}\right) = \frac{c^2\Delta x^2}{\Delta t^2}\sin^2\left(\frac{k_x\Delta x}{2}\right) + \frac{c^2\Delta y^2}{\Delta t^2}\sin^2\left(\frac{k_y\Delta y}{2}\right)$$

相対論的速度のビームとの共鳴

$$\omega = kV \quad (V \sim c)$$

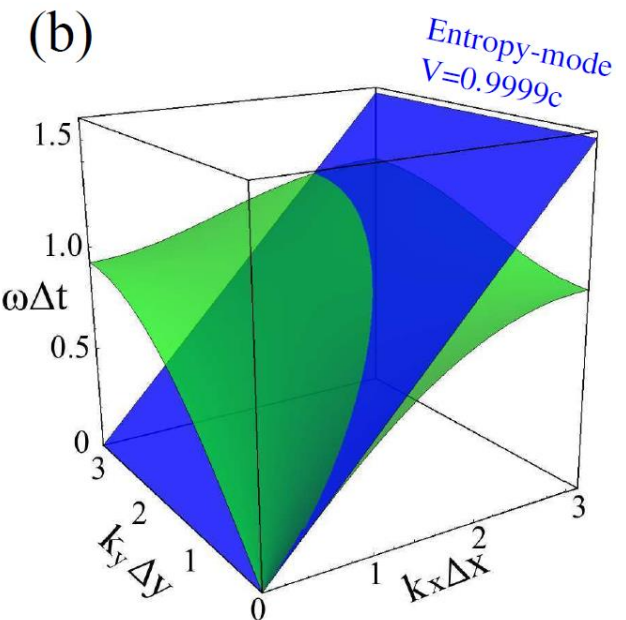
なので離散化された世界では物理的に正しい  
(エネルギー保存はされる)



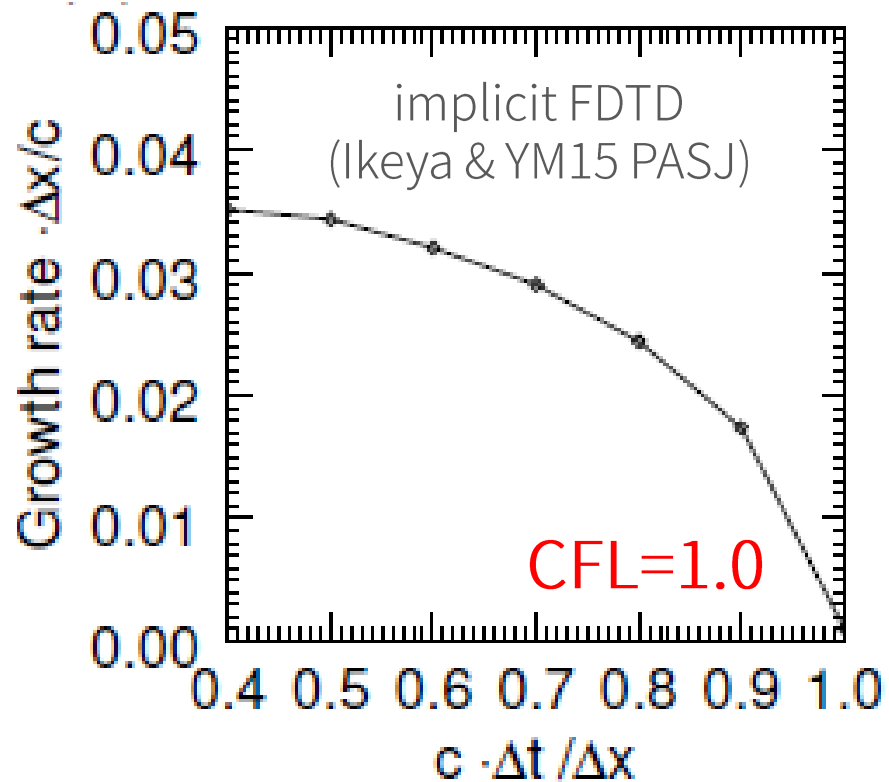
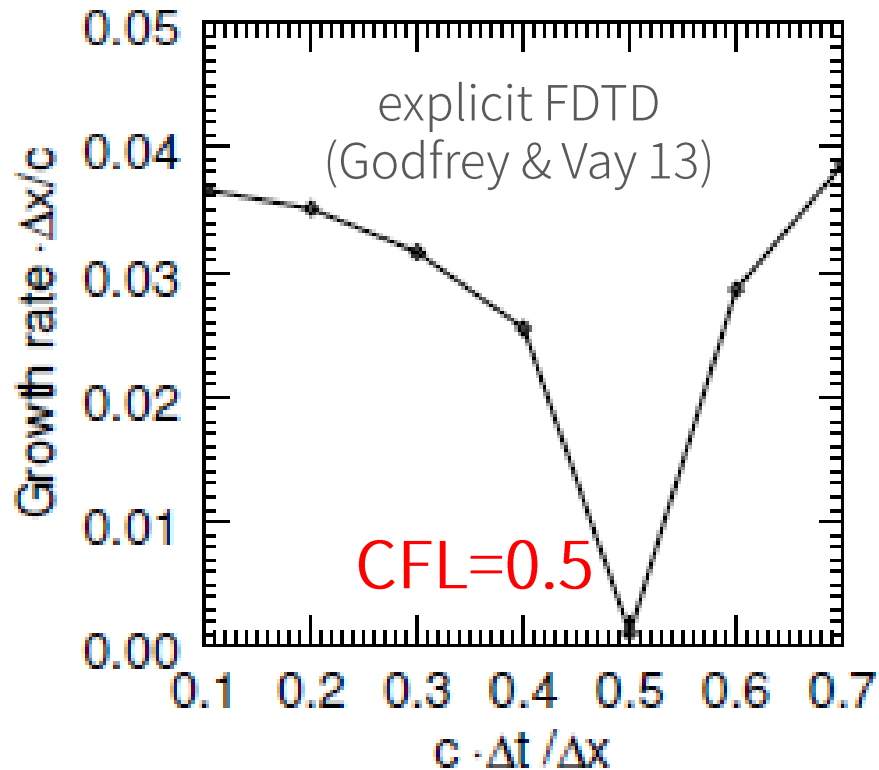
- ✓ 周期境界系
- ✓ X方向に相対論的流れ  
( $\gamma = 100$ )

## 対処方法

- スペクトル法で解く (Kato07; Vay+ 13)
  - ✓ フーリエ空間でマクスウェル方程式を解く
  - ✓ 波数空間でNCIを除去
  - ✓ 大規模並列計算で効率が出ない
- FDTD (差分法)
  - ✓ J, E, B に対してフィルタリング (Greenwood+ 04; Sironi & Spitkovsky 09)
  - ✓ 物理的な電磁波も同時に落としてしまう

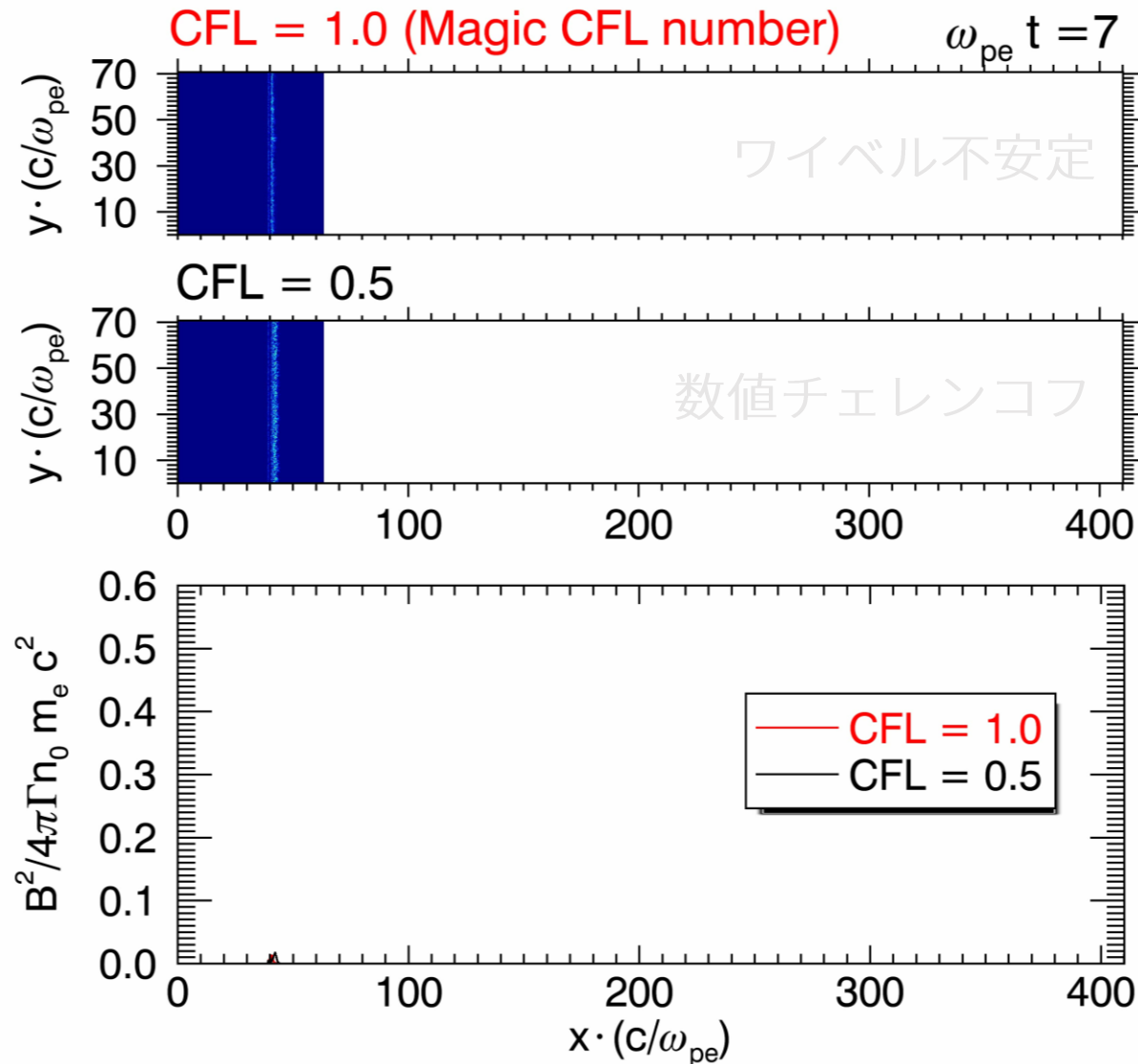


# マジカルCFL数の発見 (Vay+11 JCP)

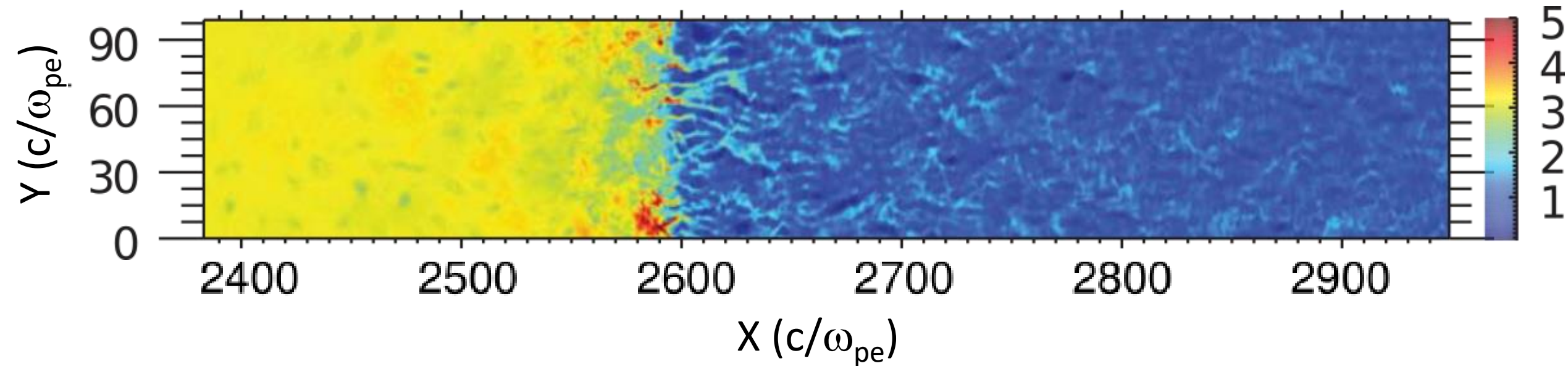


- NCIの成長率が特定のCFL数 ( $c\Delta t/\Delta x$ )で急激に減少することが発見 (Vay+ 11)
- PICアルゴリズムに強く依存
  - ✓ Esirkepovの電流計算法 (Esirkepov 00)
  - ✓ 電磁場に対して陽解法 (Godfrey & Vay 13) & 陰解法 (Ikeya & YM 15)
  - ✓ 粒子に対する電磁場の補間方法にも依存

# マジカルCFL数を用いた陽電子・電子 相対論衝撃波計算 (上流ローレンツ因子=100)



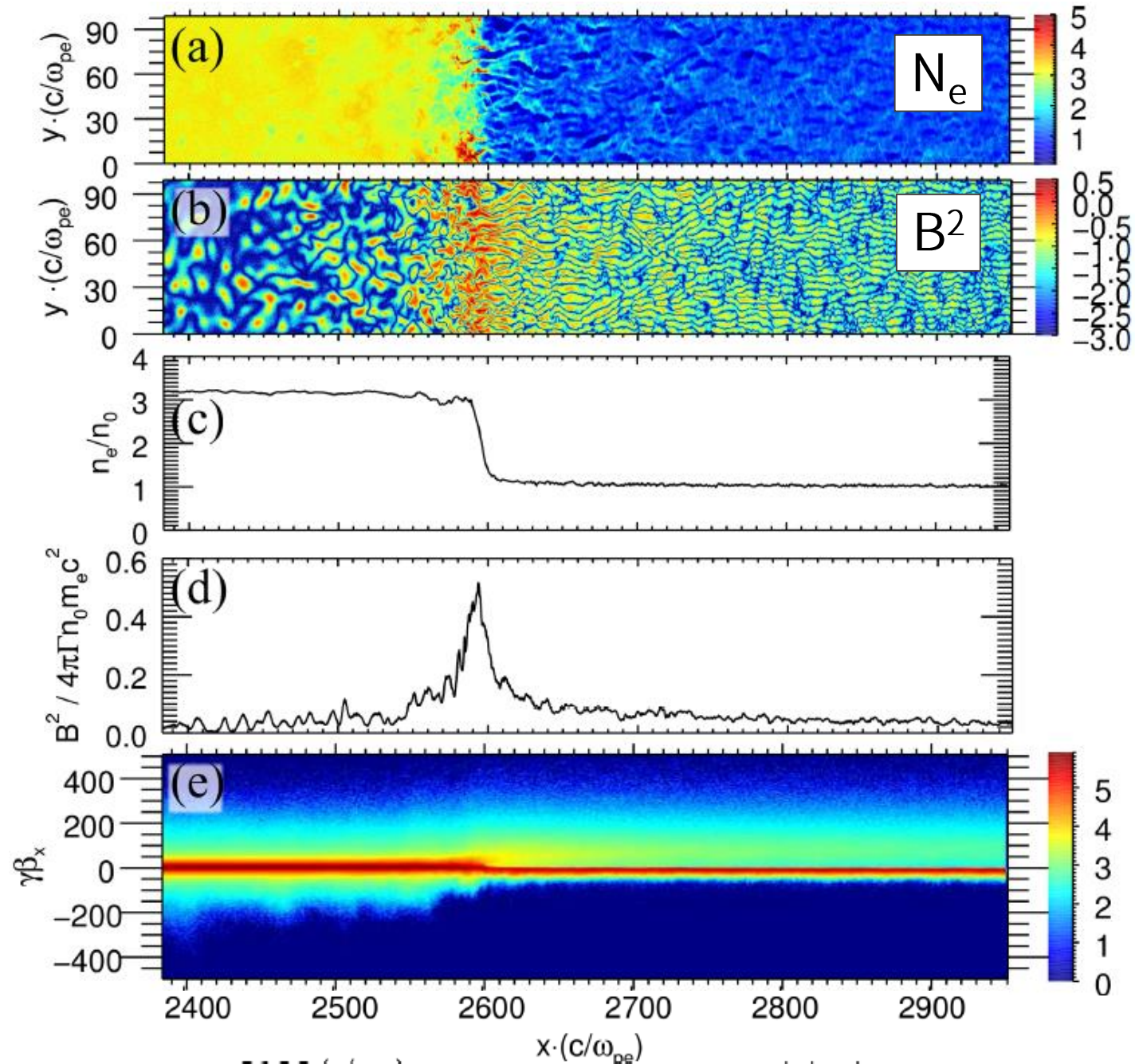
# Long-term 2D PIC simulations



- Un-magnetized Weibel-mediated shocks in pair ( $e^+$  &  $e^-$ ) plasmas
- Magic CFL number method vs. previous results (smoothing current density, *cf.* Sironi+13)
- Long-term simulations  $\omega_{pe} t \sim 6000$  to observe particle accelerations
- $\Gamma = 15$ ,  $N_0 = 8, 16, 25, 50$  /cell in the upstream region
- $N_x \times N_y = 80,000 \times 1,400$
- 648 cores & 3.5TB memory on Cray XC30 at CfCA, NAOJ

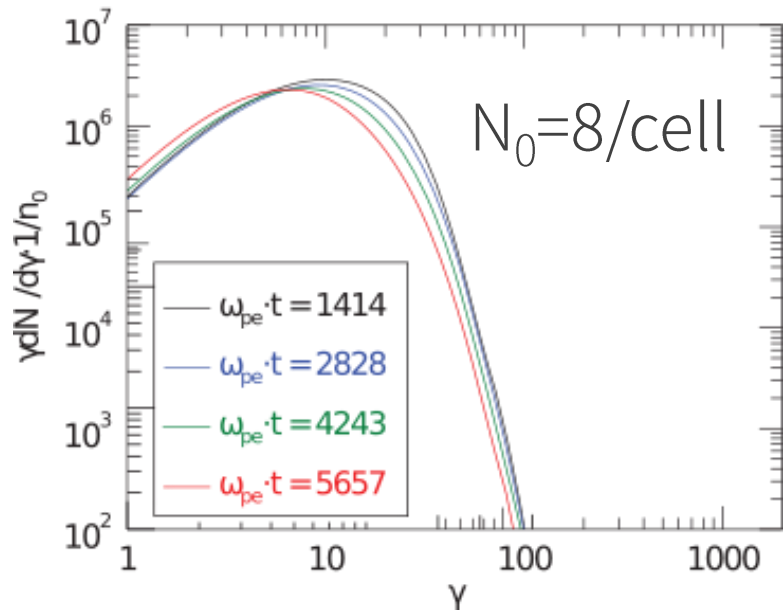
# 2D shock structure ( $N_0=50$ /cell)

- Basically followed Sironi+ 13
- Weibel filaments in far upstream regions
- ~10 times stronger magnetic field turbulence
- Present method can maintain magnetic fields in larger amplitudes

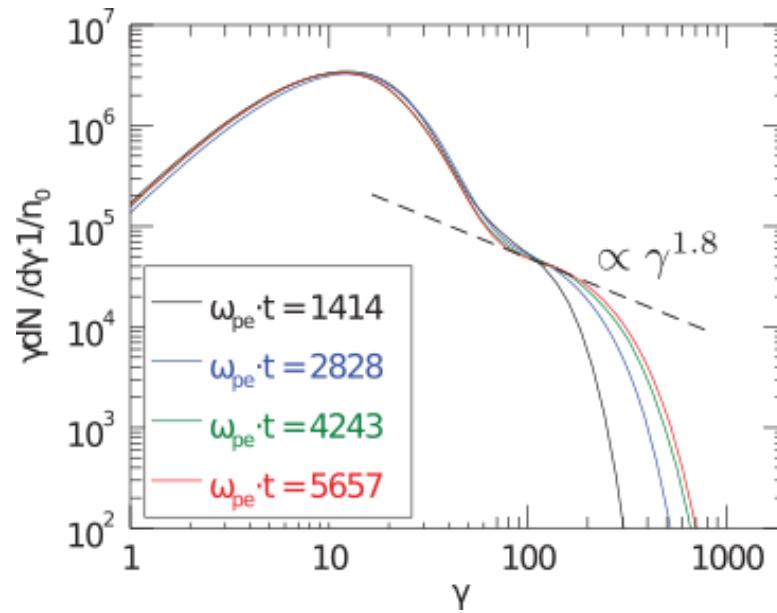


# Time evolutions of downstream energy spectrum

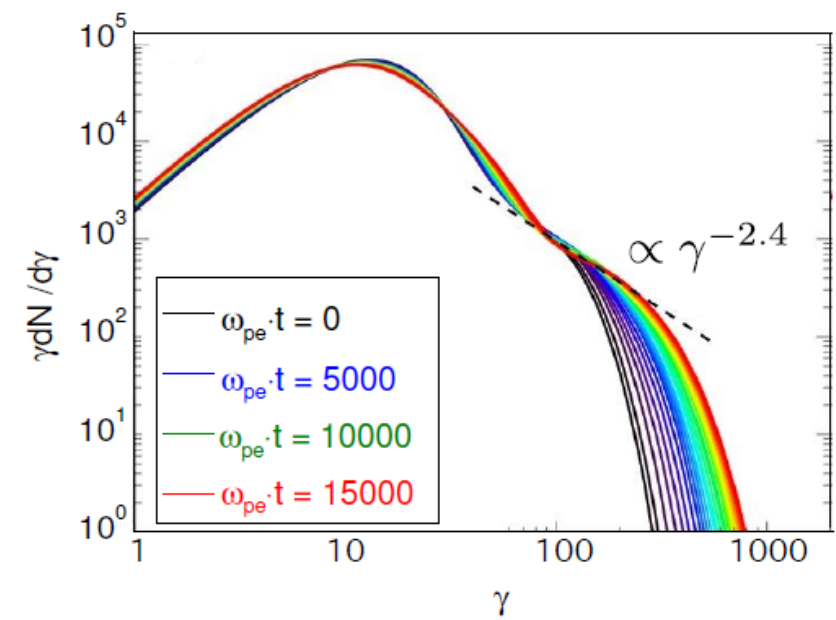
CFL=0.5 (Non-magic)



$N_0=50/\text{cell}$ , CFL=1.0



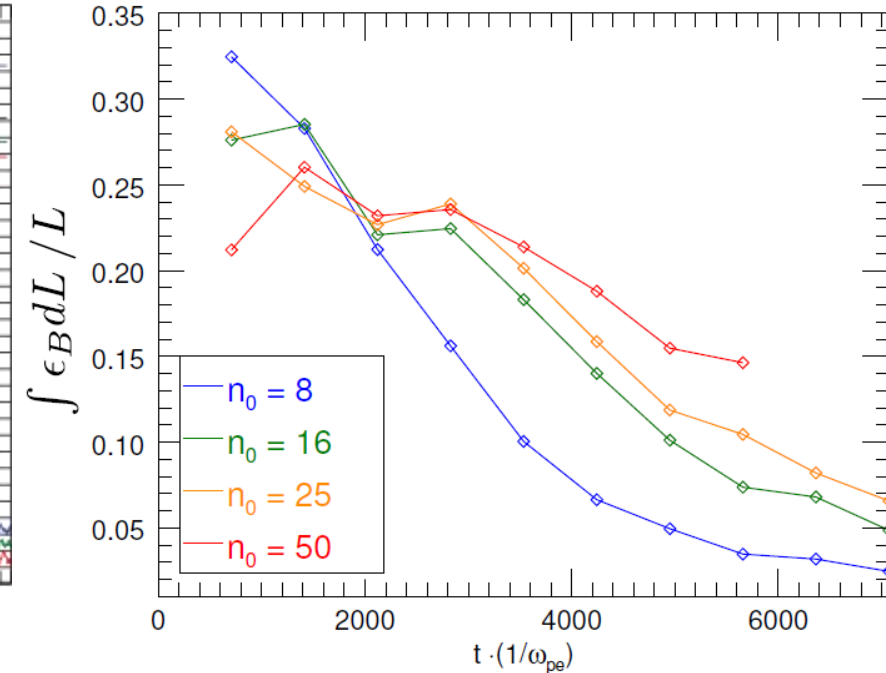
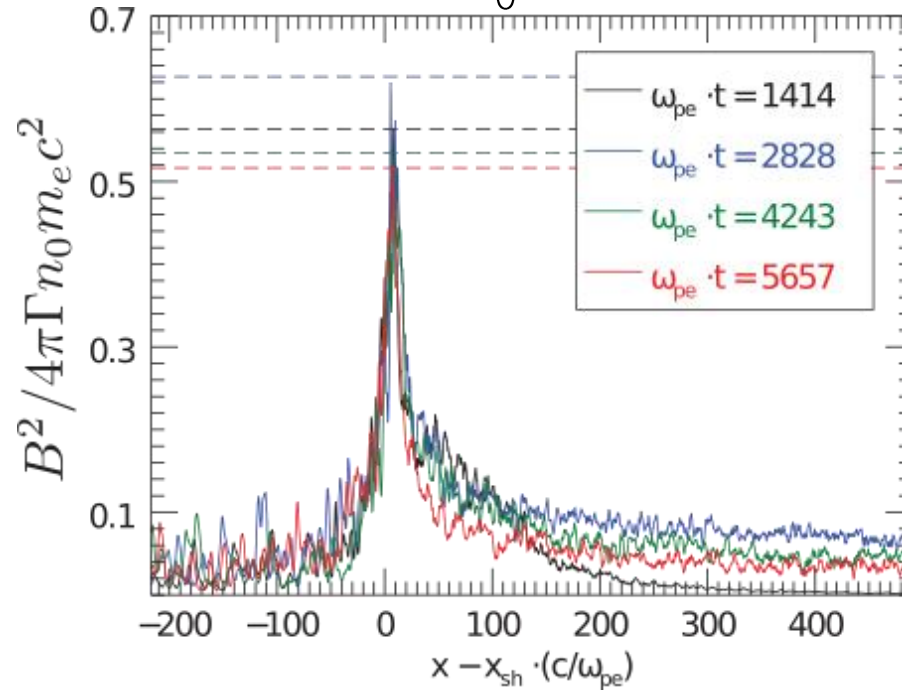
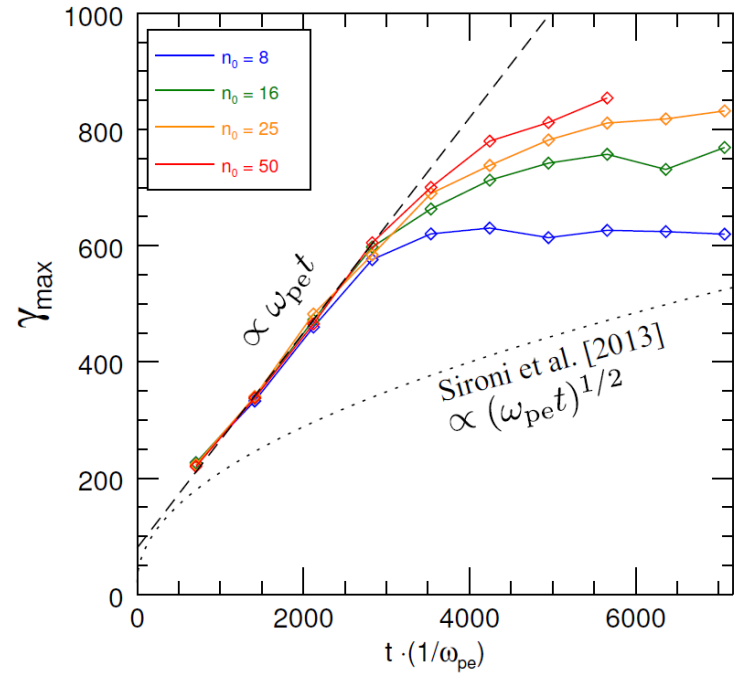
Sironi+13



- ❑ No efficient acceleration with NCI
- ❑ Efficient accelerations were obtained by using Magic CFL number method
- ❑ Spectral index -1.8 vs. -2.4

# $\gamma_{\max}$ time evolutions

$N_0=50$



- Rapid acceleration until  $\omega_{pe} t \sim 3000$  due to strong magnetic turbulence around the shock front ( $\gamma_{\max} \propto t$ )
- Followed by slow acceleration stage due to decay of the Weibel magnetic turbulence ( $\gamma_{\max} \propto t^{1/2}$  by Sironi+ 13)
- Results seemed to be converged with  $n_0=50$

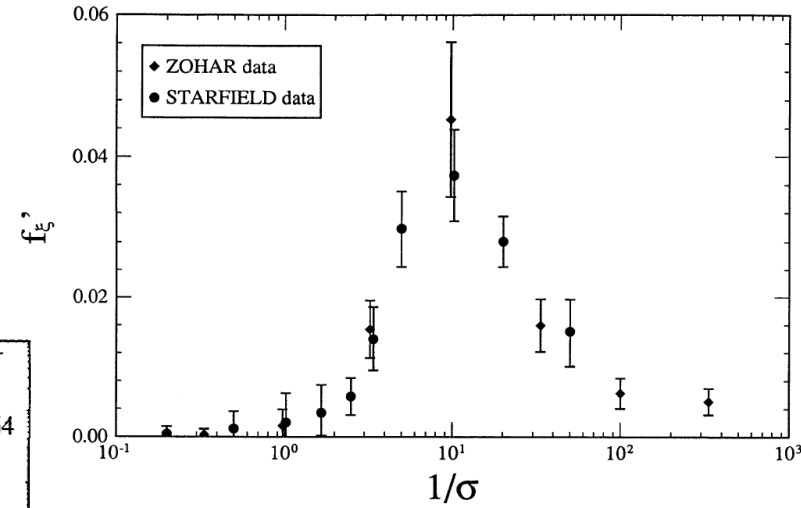
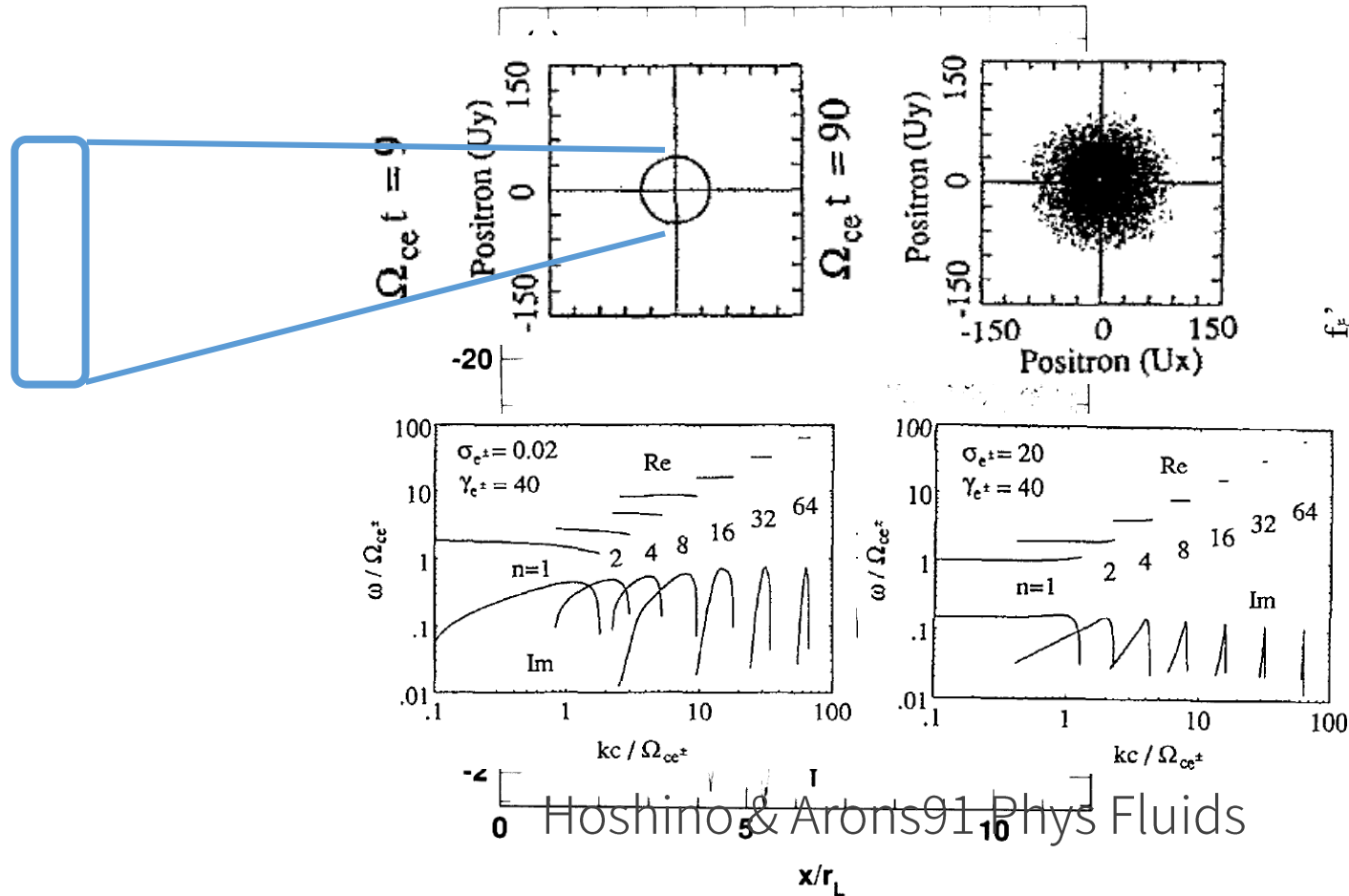


# 相対論的衝撃波からの大振幅 先駆波励起と電子加速

# 1990s~ 1D PICシミュレーション

最初のPIC計算  
synchrotron maser instability  
precursor wave

emission efficiencies for  
shock magnetizations

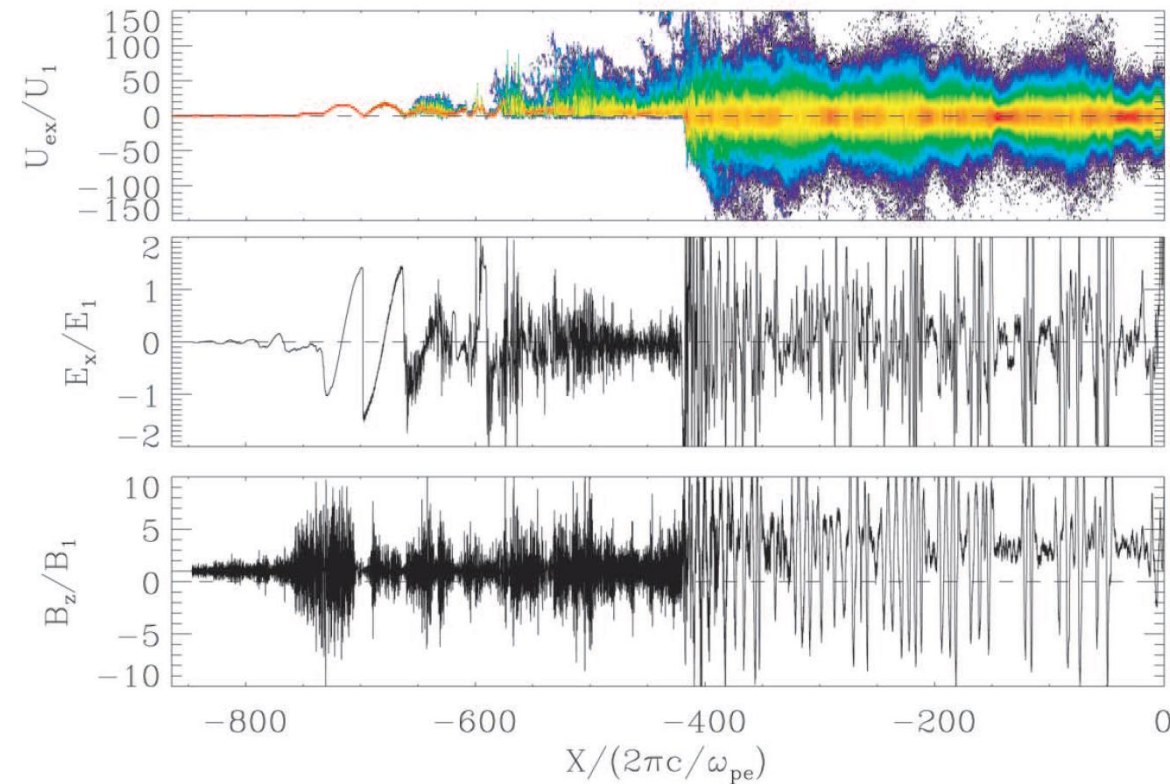


Gallant+91 ApJ;  
Hoshino+92 ApJ

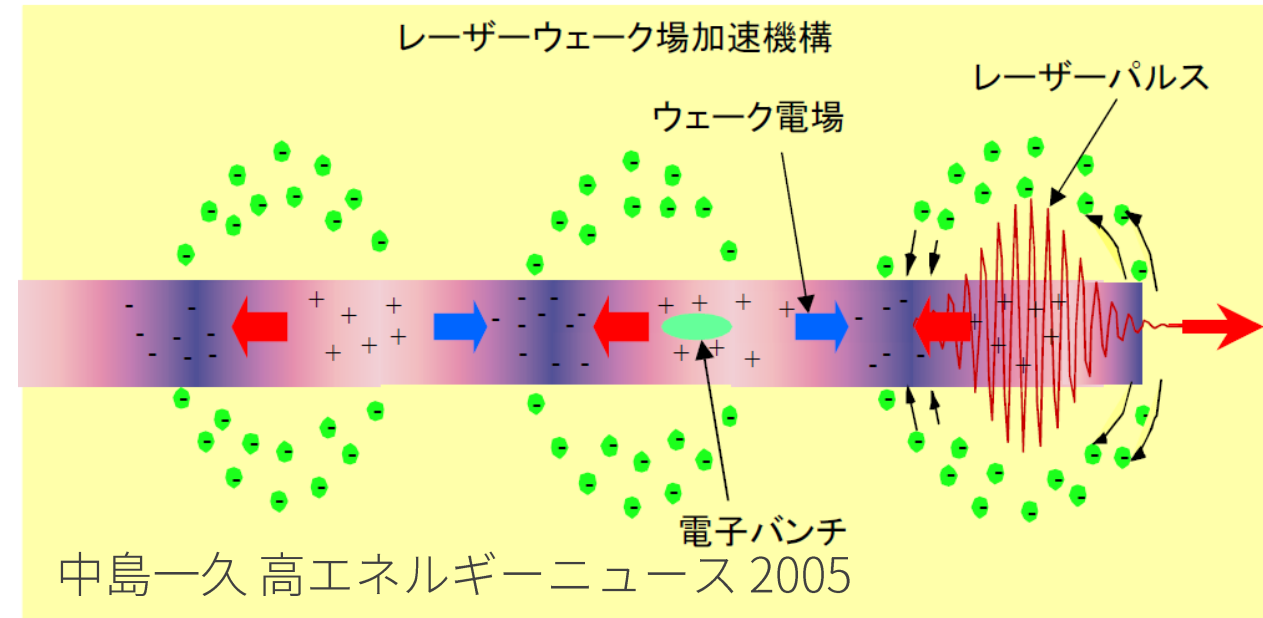
Hoshino & Arons 91 Phys Fluids  
Langdon, Arons, Max88 PRL

# 先駆波励起と航跡場加速 (1D PIC, Lyubarsky06; Hoshino08)

□ Tajima & Dawson 79 PRL



Hoshino08

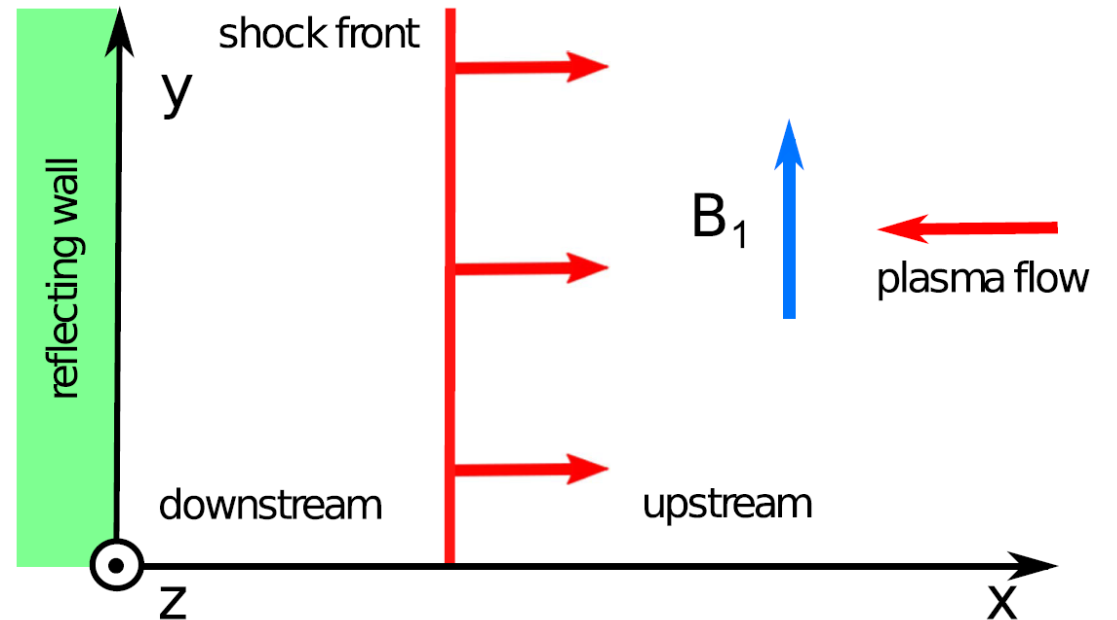
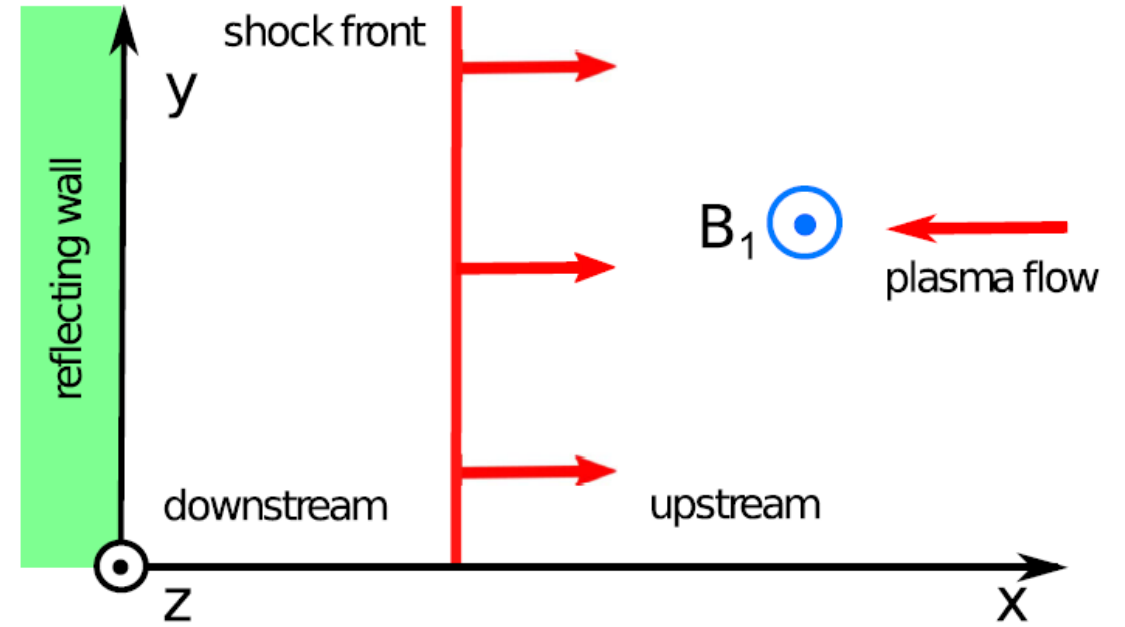


□ レーザー（光モードの電磁波）のポンデロモータイブ力  $(-\frac{e^2}{4m\omega^2} \nabla E^2(x))$  によって電子が掃きだされ、ウェークに強い静電場（ラングミュア波）が励起

□ ラングミュア波とのランダウ共鳴

# Iwamoto+17;18 ApJ

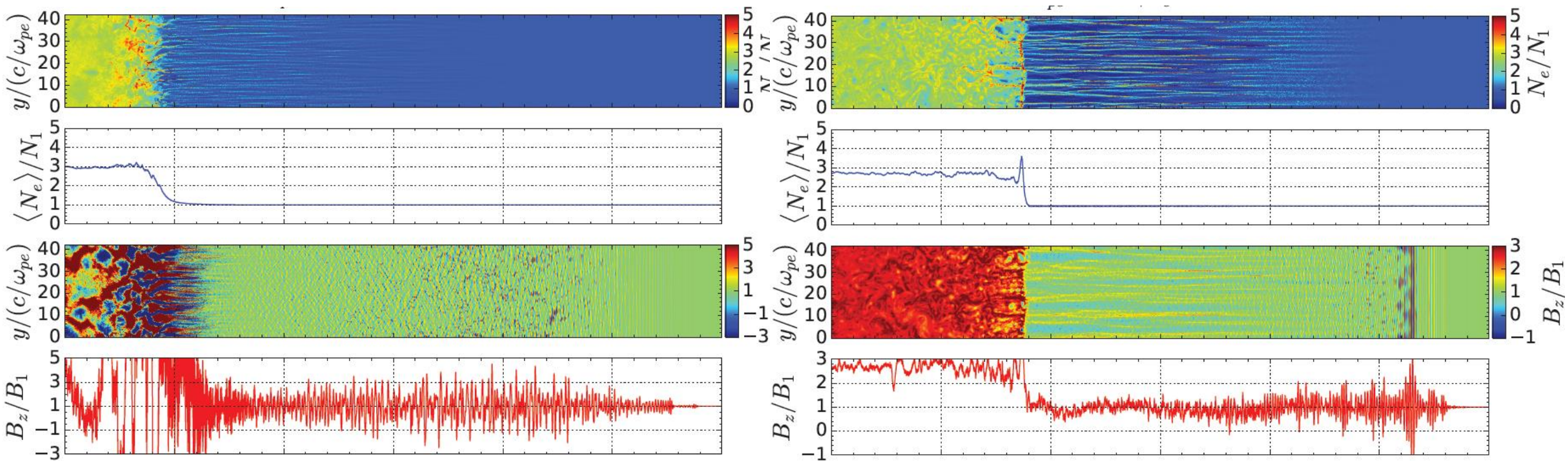
- 2次元PICシミュレーションでワイベル不安定との競合を調べた
- 上流バルク $\Gamma=40$
- 陽電子-電子系 ( $M/m=1$ )
- イオンはないので航跡場加速は起きない。Precursor waveの放射効率のみを議論
- マジカルCFL法
- $N=64$  /cell/species
- $dx/(c/\omega_{pe}) = 1/40$  (typically  $\sim 1/10$ )



# Persistence of precursor waves (out-of-plane $B_0$ )

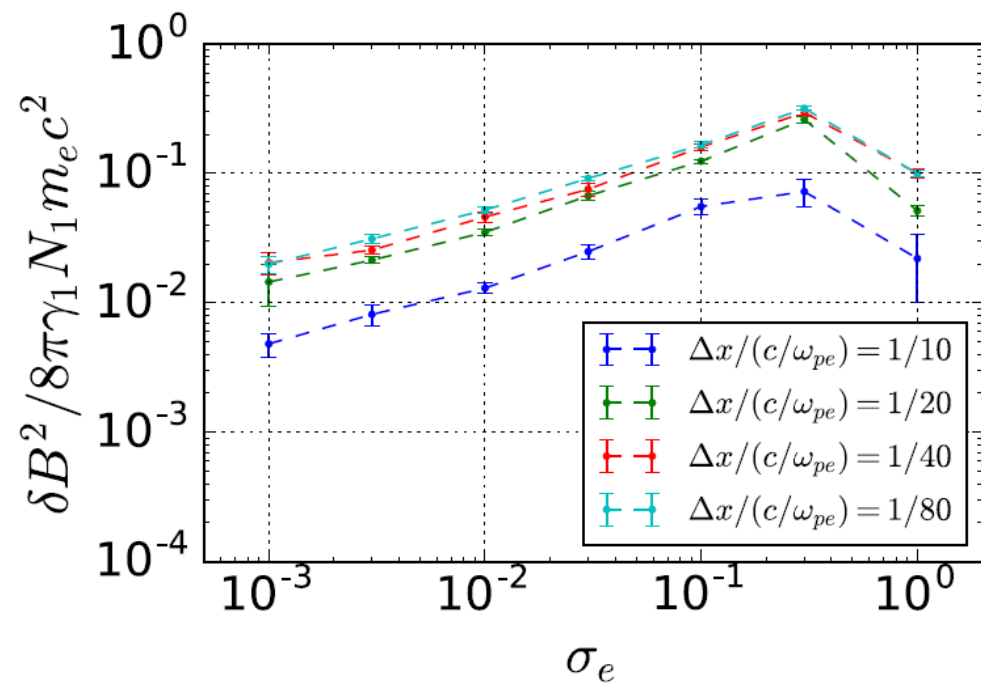
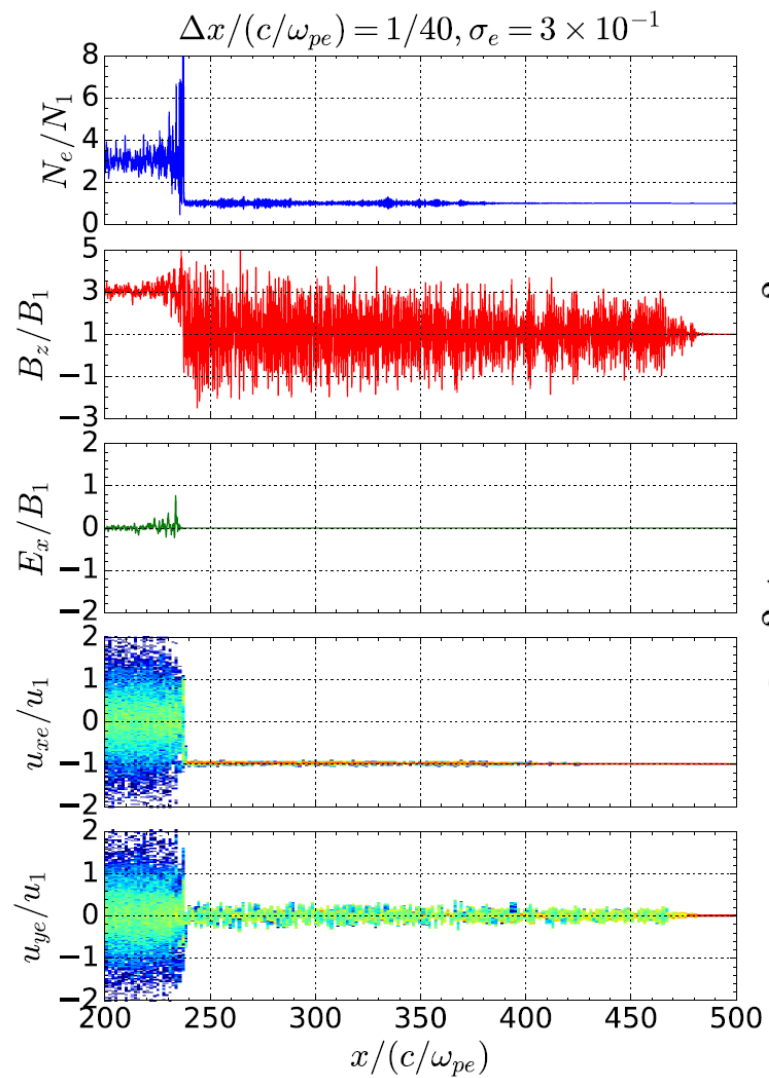
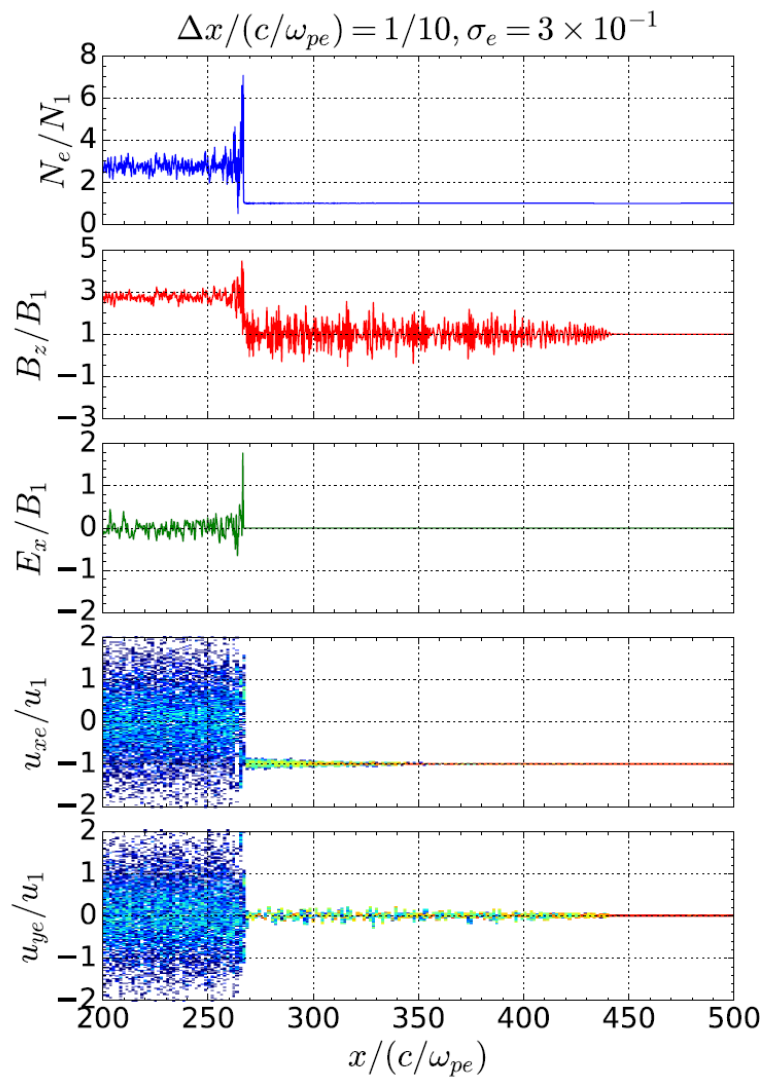
$$\sigma_e = 3 \times 10^{-3}$$

$$\sigma_e = 3 \times 10^{-1}$$

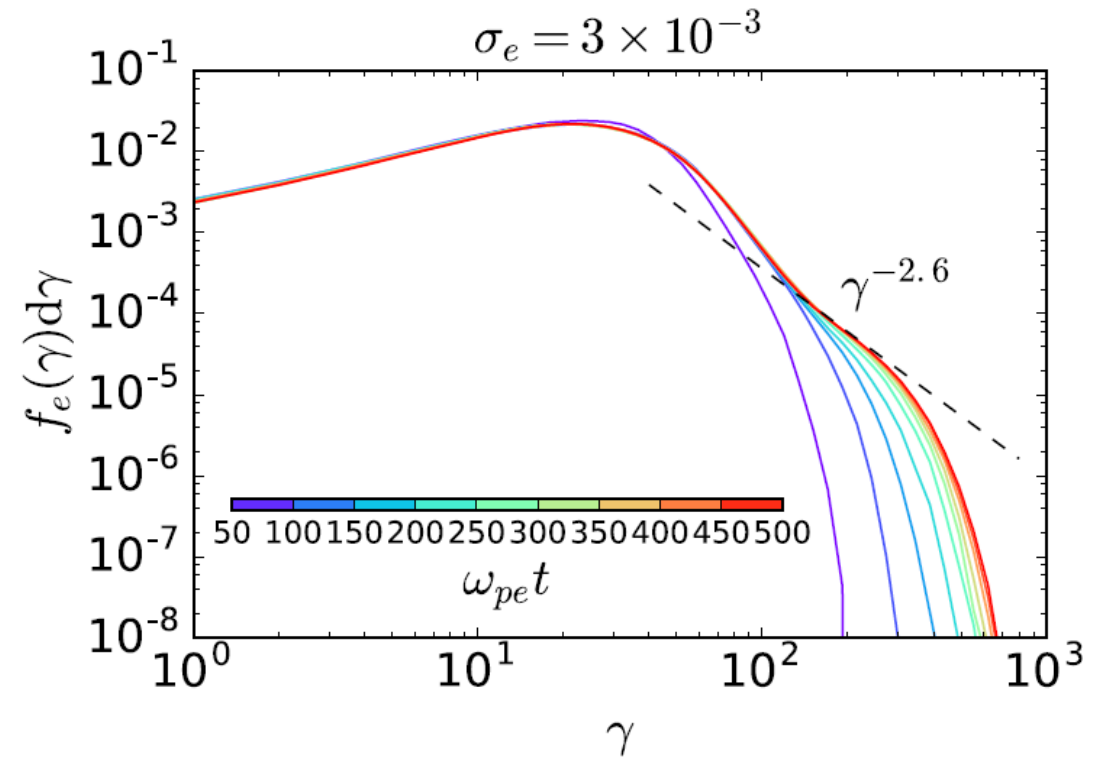
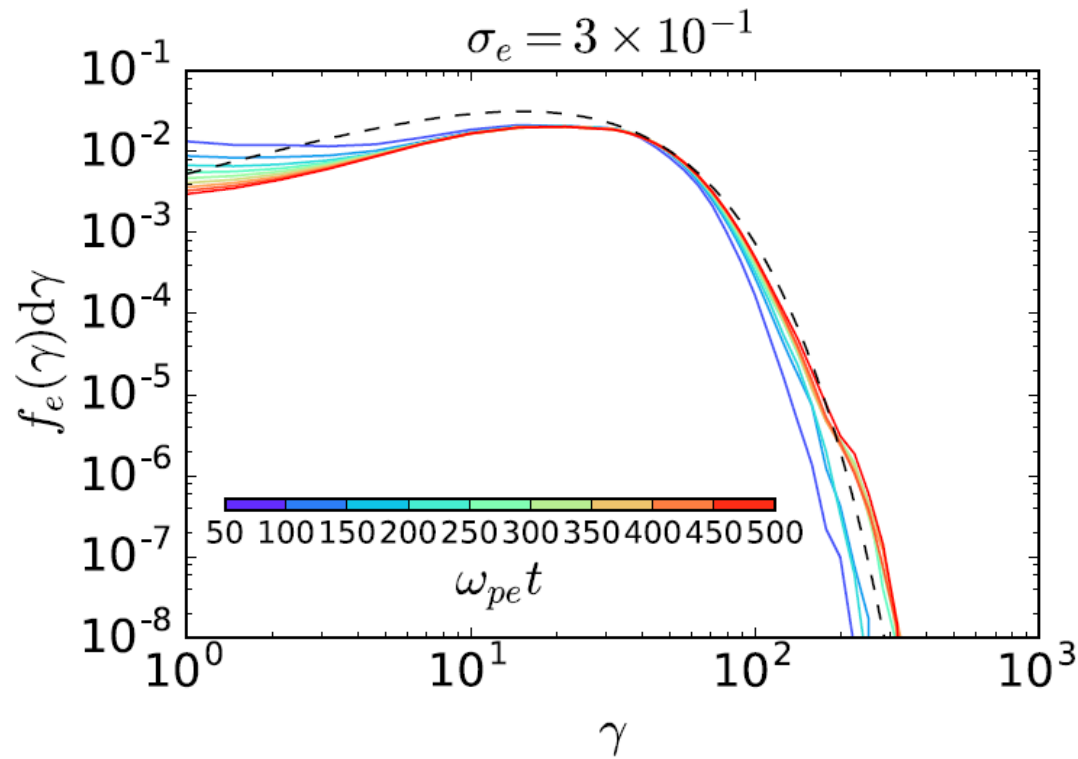


- ❑ 多次元計算では放射されないと言われていたが (Sironi+11;13) 低 $\sigma$ でも大振幅の電磁波放射 (ワイベル不安定と共存が可能) !
- ❑ 多次元性による効果 : 電磁波の自己集束による上流密度揺らぎ

# 解像度依存性

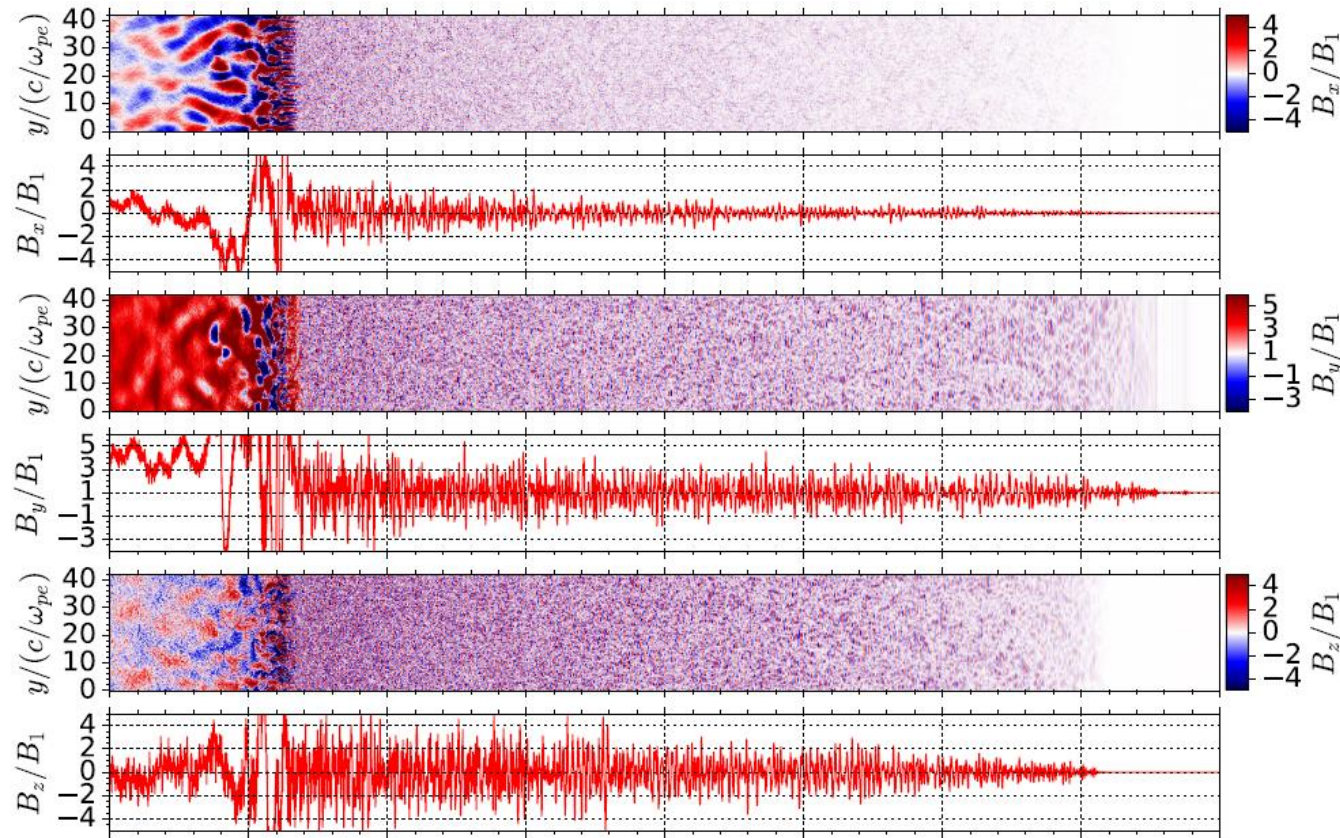


# 下流のエネルギーースペクトル



- 陽電子・電子系なので、航跡場加速は起きない
- 低 $\sigma_e$ での非熱的粒子加速はWeibel乱流によるフェルミ加速 (Sironi+ 13)
- Sironi+13より早い時間スケールでの加速 (マジカルCFL法の恩恵)

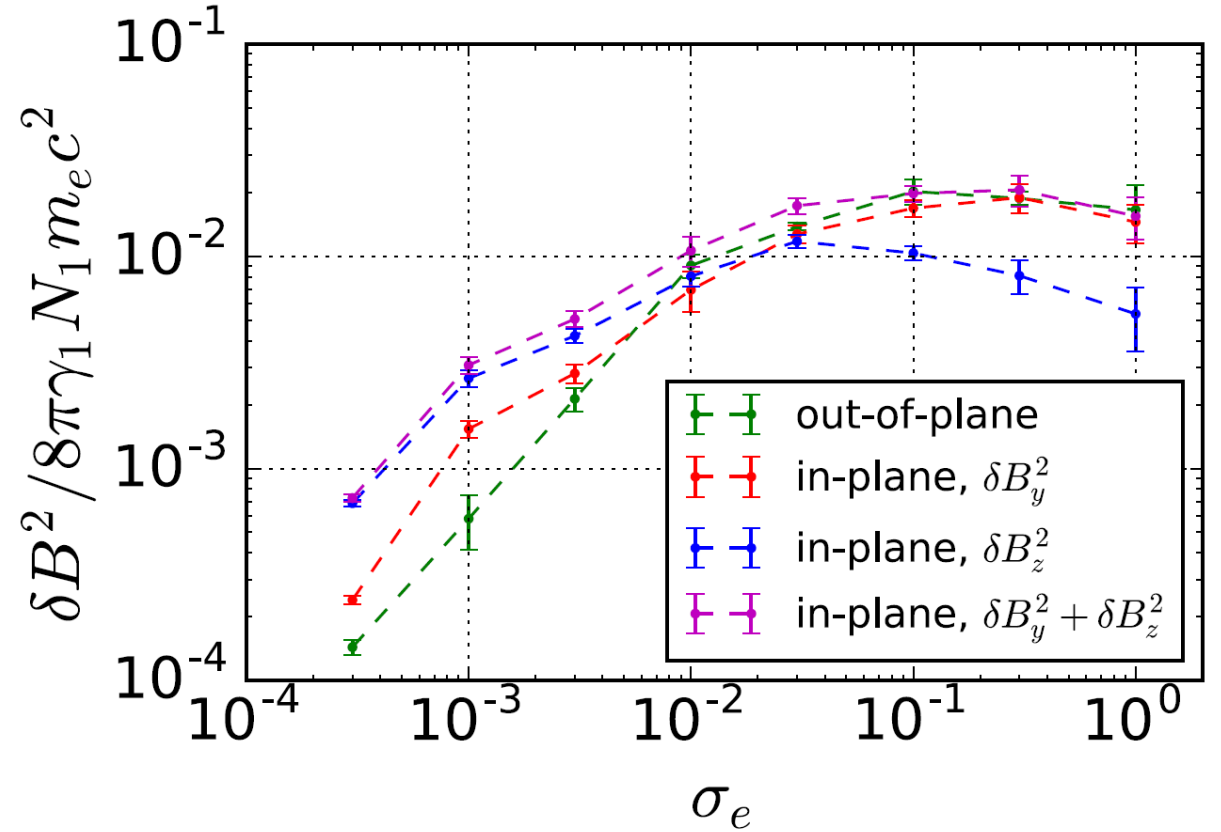
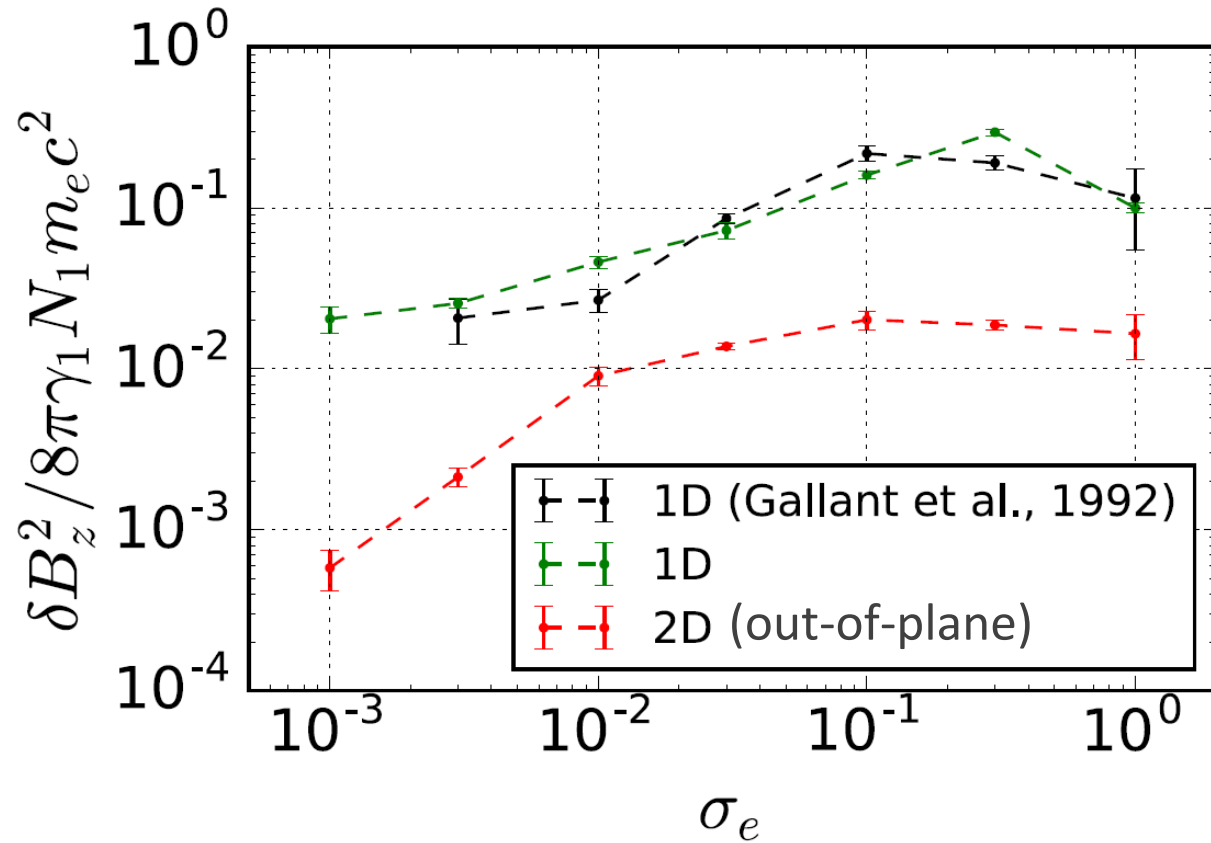
# Precursor wave emission enhanced by the Weibel instability (in-plane $B_0$ )



- いわゆるSMIで励起される電磁場 (X-mode) は背景磁場方向に変動成分 ( $\delta B_y$ )
- $\delta B_z$ 成分も遅れて強い励起 (電磁場O-mode)
- 衝撃波面近傍の強いワイベル磁場が $\delta B_z$ 成分を生成するため、粒子のジャイロ運動は紙面垂直方向に傾くため、( $\delta B_y, \delta B_z$ )を励起
- 上流領域 ( $B \sim B_y$ )でO-modeにモード変換

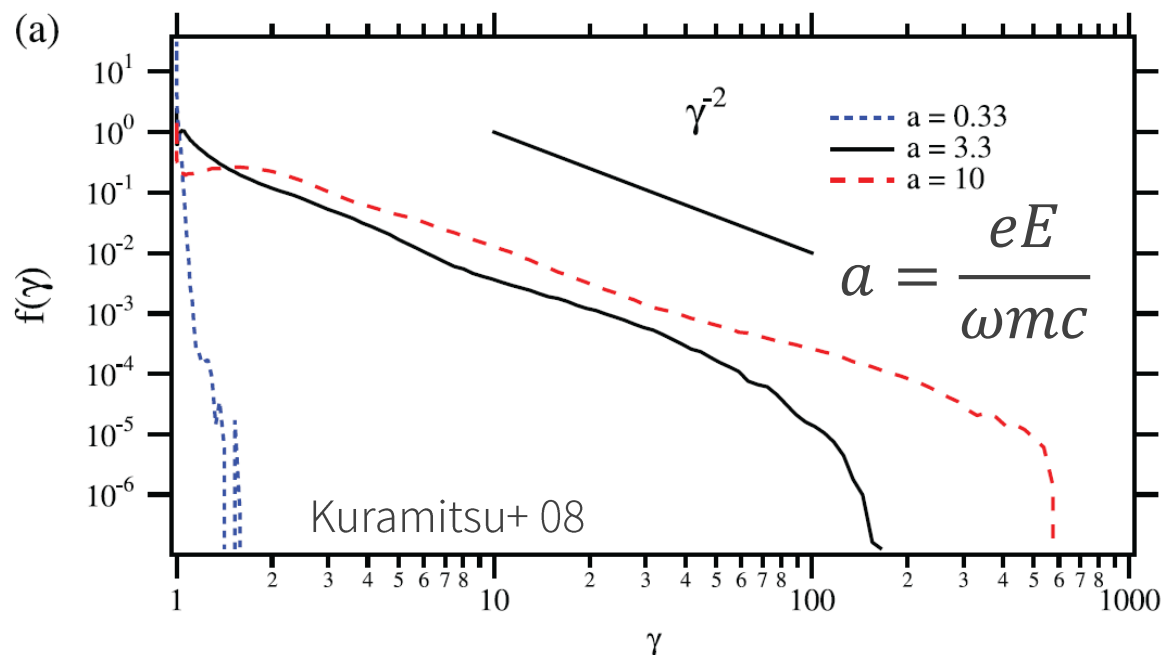


# 先駆波の振幅の $\sigma$ 依存性



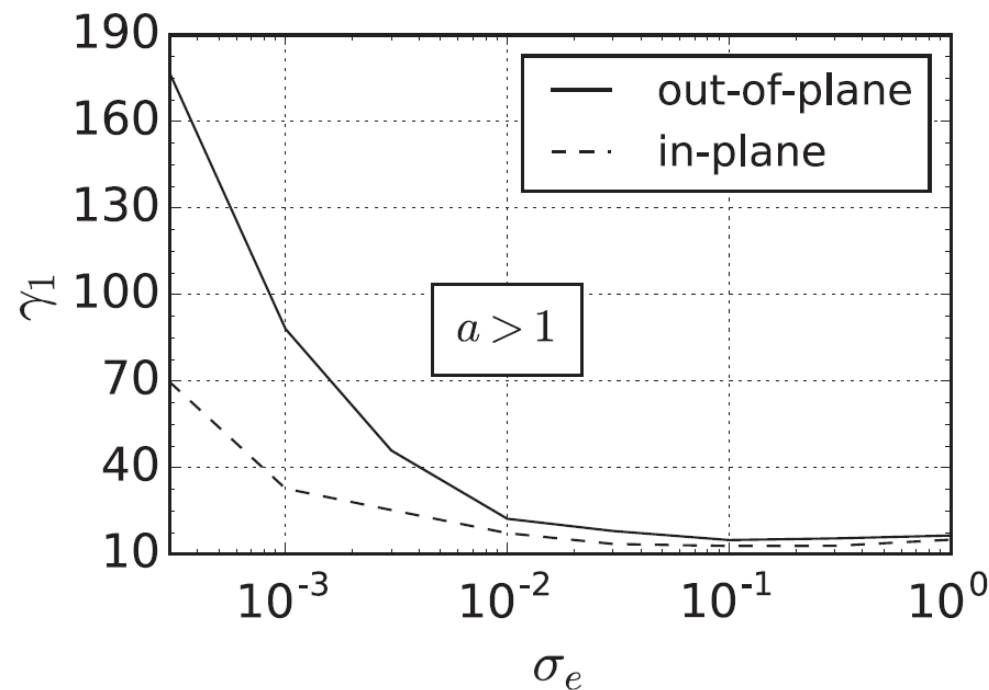
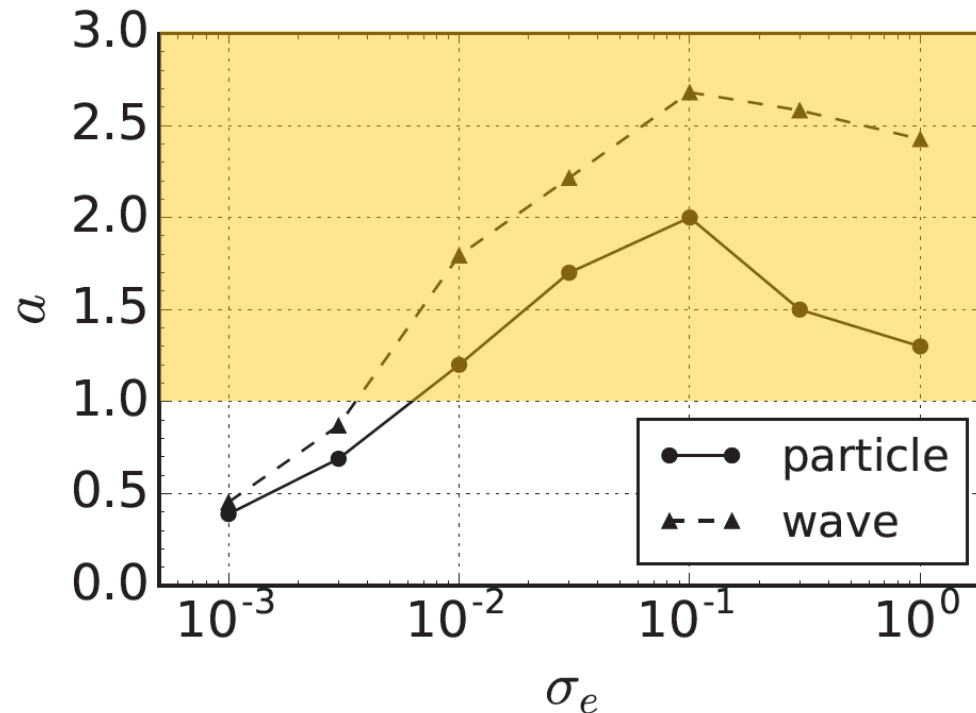
3D計算結果を待ちましょう！

# 航跡場加速への期待



Strength parameter  $a > 1$  で効率的な航跡場加速 (Kuramitsu+08)

幅広い  $\sigma$  で  $a > 1$  の大振幅電磁波励起が実現



# イオンワイベル不安定による 磁場生成と飽和過程

# Origin of the seed magnetic field (non-MHD process)

$$\frac{\partial B}{\partial t} = -c \nabla \times E = \nabla \times (v \times B)$$

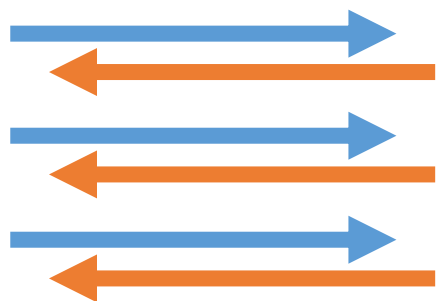
MHD dynamo  
for global magnetic fields

Seed magnetic field?

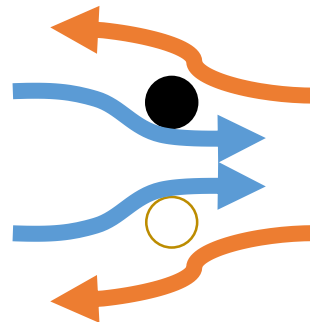


Weibel instability (Weibel, PRL, 1959; Fried, Phys. Fluid, 1959)

Beams / temperature anisotropy



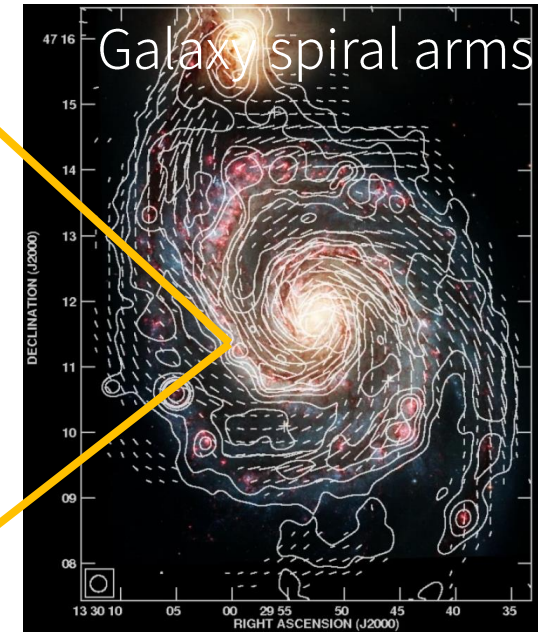
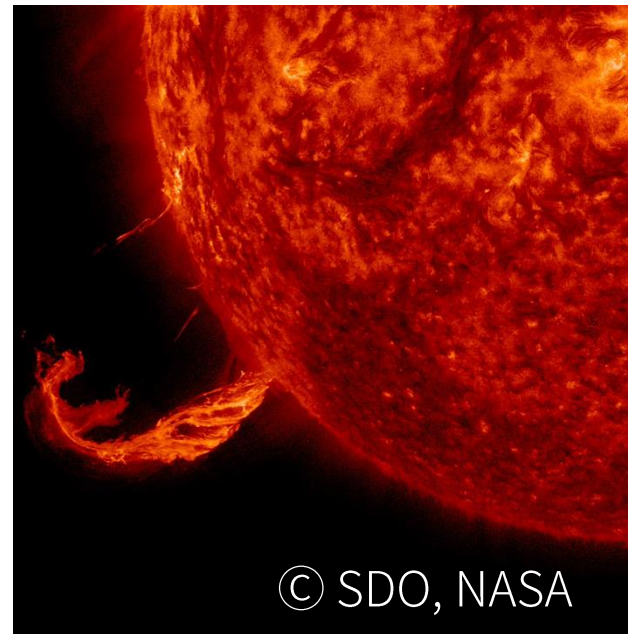
$+\delta B$



$$j \rightarrow \nabla \times E$$



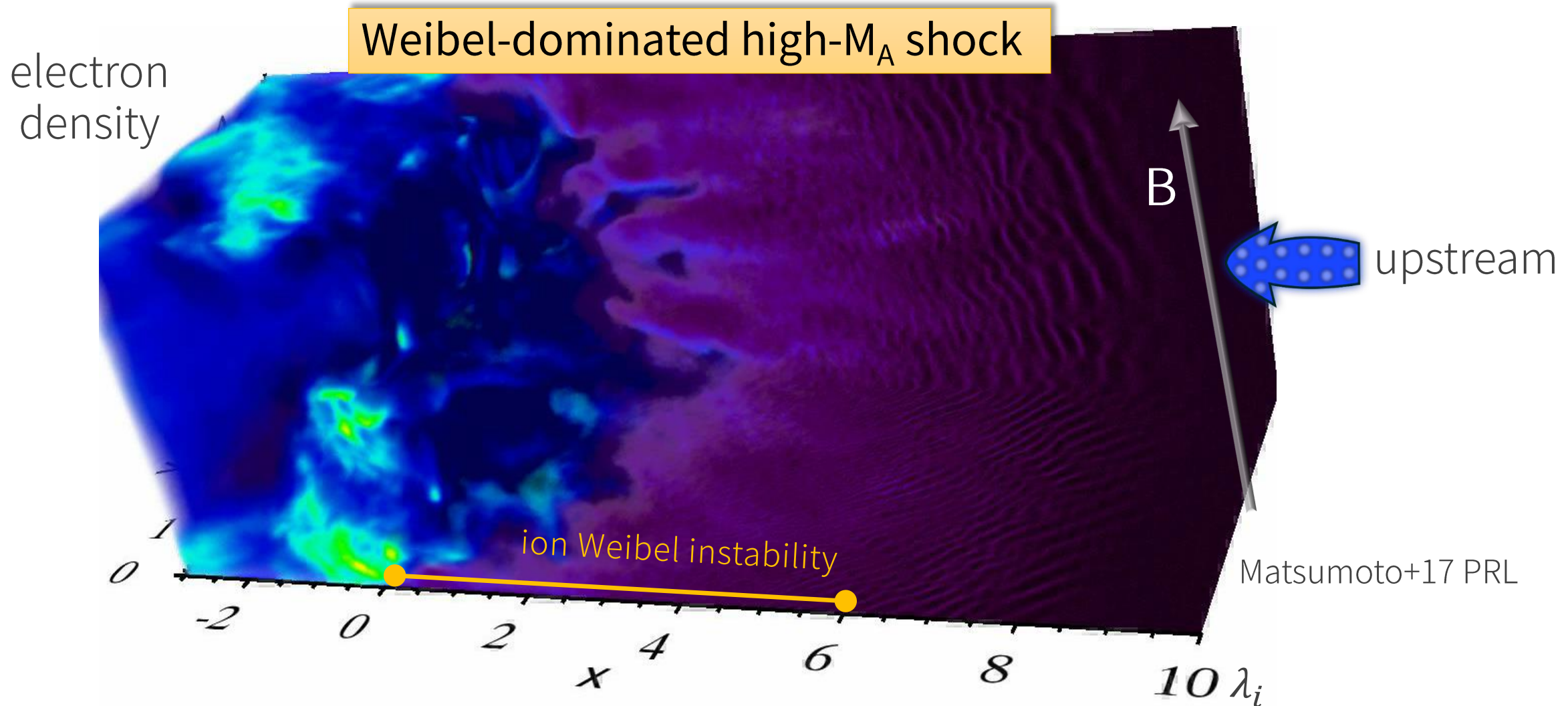
$c/\omega_{p,s}$



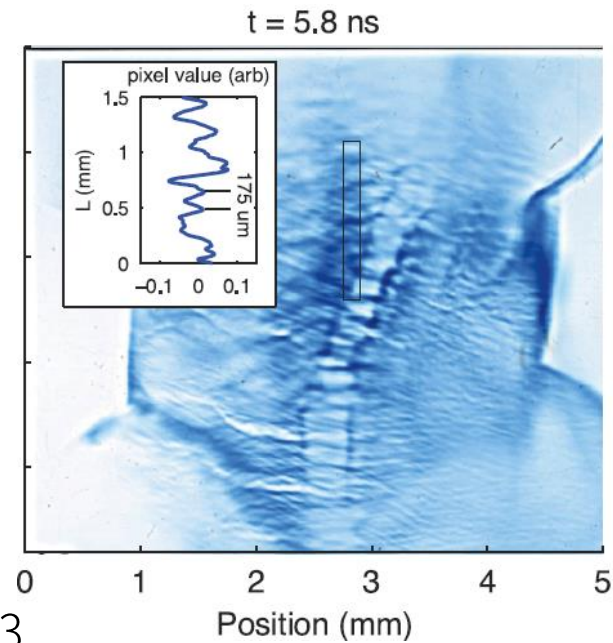
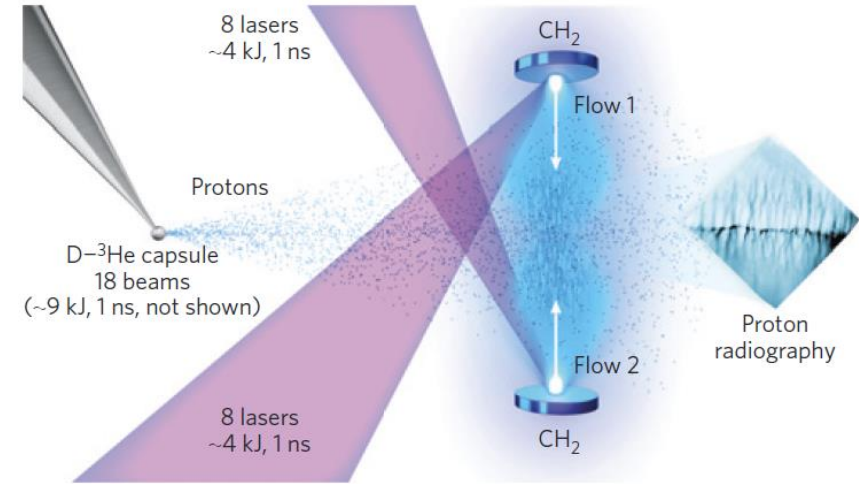
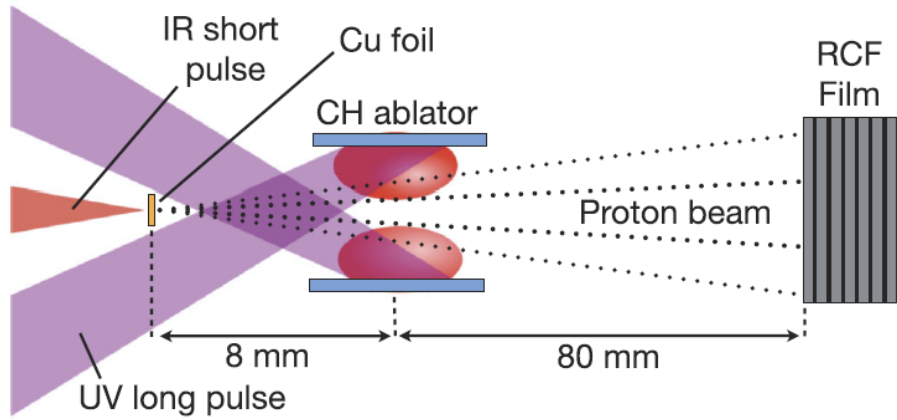
Fletcher+11

# Weibel instability in high-energy astrophysics

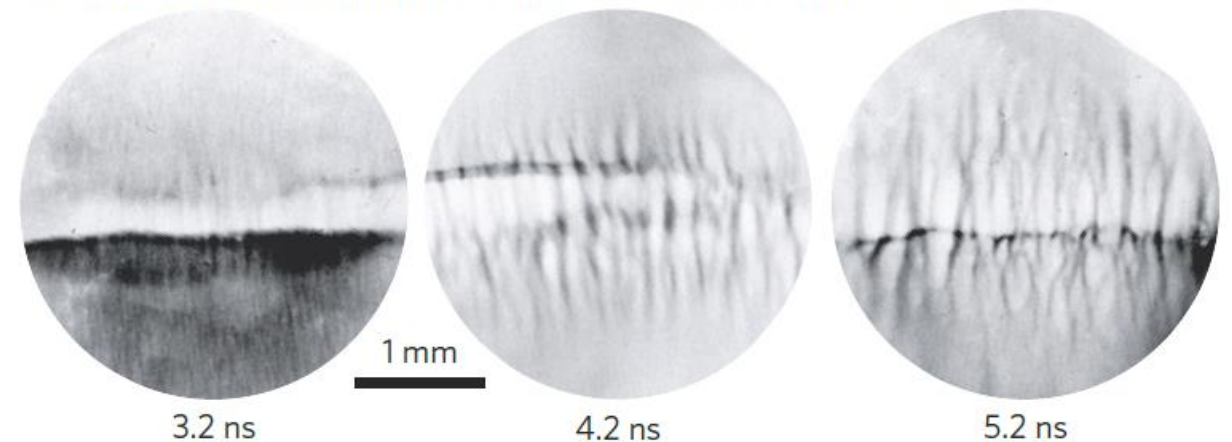
- Weibel-mediated shocks (Kato07; Kato & Takabe08)
- Particle accelerations at shocks (Spitkovsky08; Sironi+11,13; Matsumoto+15,17)



# Weibel instability in laboratory experiments



Experimental proton radiographs from 14.7 MeV (D-<sup>3</sup>He) protons



# Saturation mechanisms of the ion Weibel instability $\Leftrightarrow$ maximum magnetic energy

■ Linear growth :  $\omega_{pi}^{-1}, c/\omega_{pi}$

■ Inverse cascade

✓ coalescence of current filaments with same polarity

✓  $I = j\pi R(t)^2, j = \text{const.}$

✓  $R(t) \propto t^2 (M/m)^{-\frac{1}{2}}$  (Ruyer+15;16)

■ Upper limit of the current

✓ Alfvén current limit (Alfvén38)

✓ Bio-Savart law

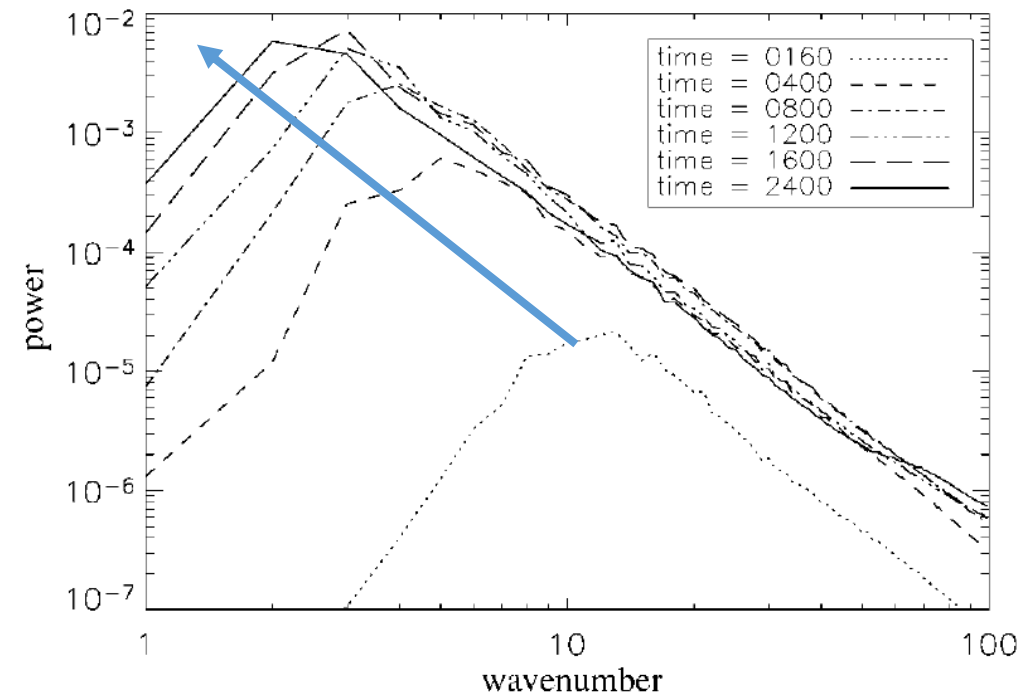
$$B = \frac{2Ir}{\xi R^2} \quad (r < R)$$

$$= \frac{2I}{cr} \quad (r > R)$$

✓ When  $r_{ge} \sim R(t)$  in self-generated magnetic field

$$I_A = \frac{mc^2 \gamma c}{q}$$

✓ Constant  $I_A$  with growing  $R$  results in monotonic decay of  $|B|$  after saturation



Frederiksen+04 ApJL

# Alfven current limit in the Weibel instability

■ 2D electron-positron plasmas (Kato05 PoP)

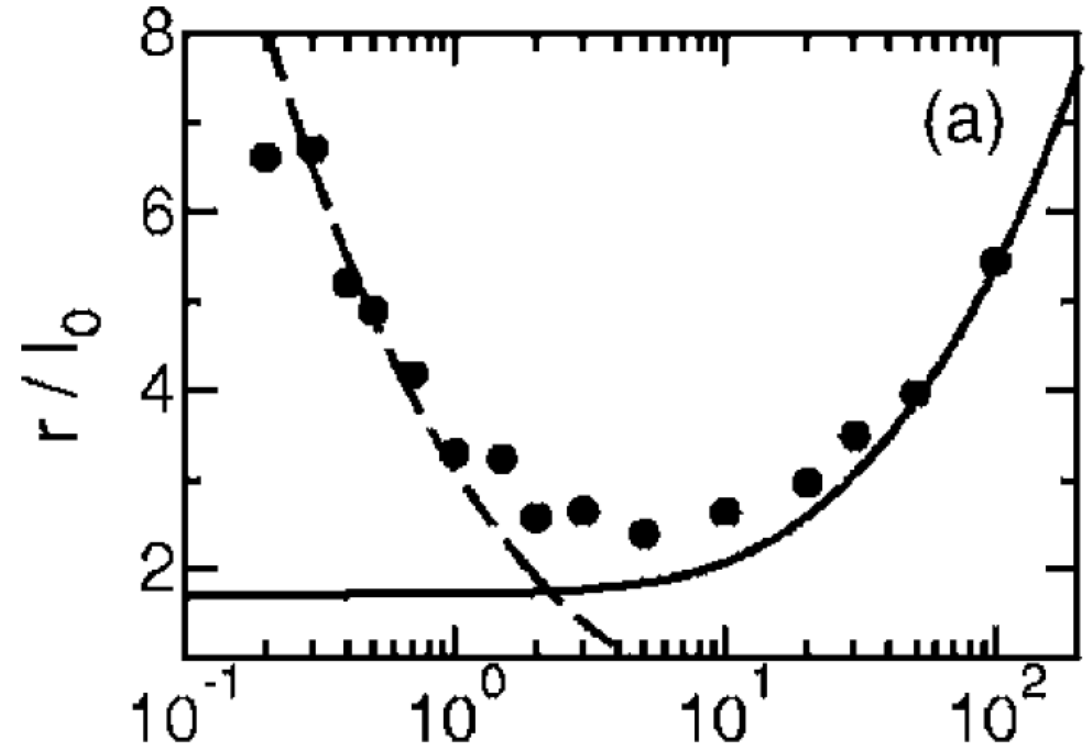
■ In the strong anisotropy,  $I_A = \pi R^2 q n \beta c$  gives the filament size at saturation as

$$r_p \sim 2\sqrt{\gamma} \frac{c}{\omega_{p,s}}$$

■ Magnetic field energy at saturation:

$$\epsilon_B = \frac{B_{max}^2}{8\pi\gamma\rho c^2} \text{ with } B_{max} = \frac{2I_A}{cr_p}$$

$$\Rightarrow \epsilon_B \sim \frac{\beta^2}{2} \sim \frac{1}{2}$$



weak anisotropy  
(non-relativistic)

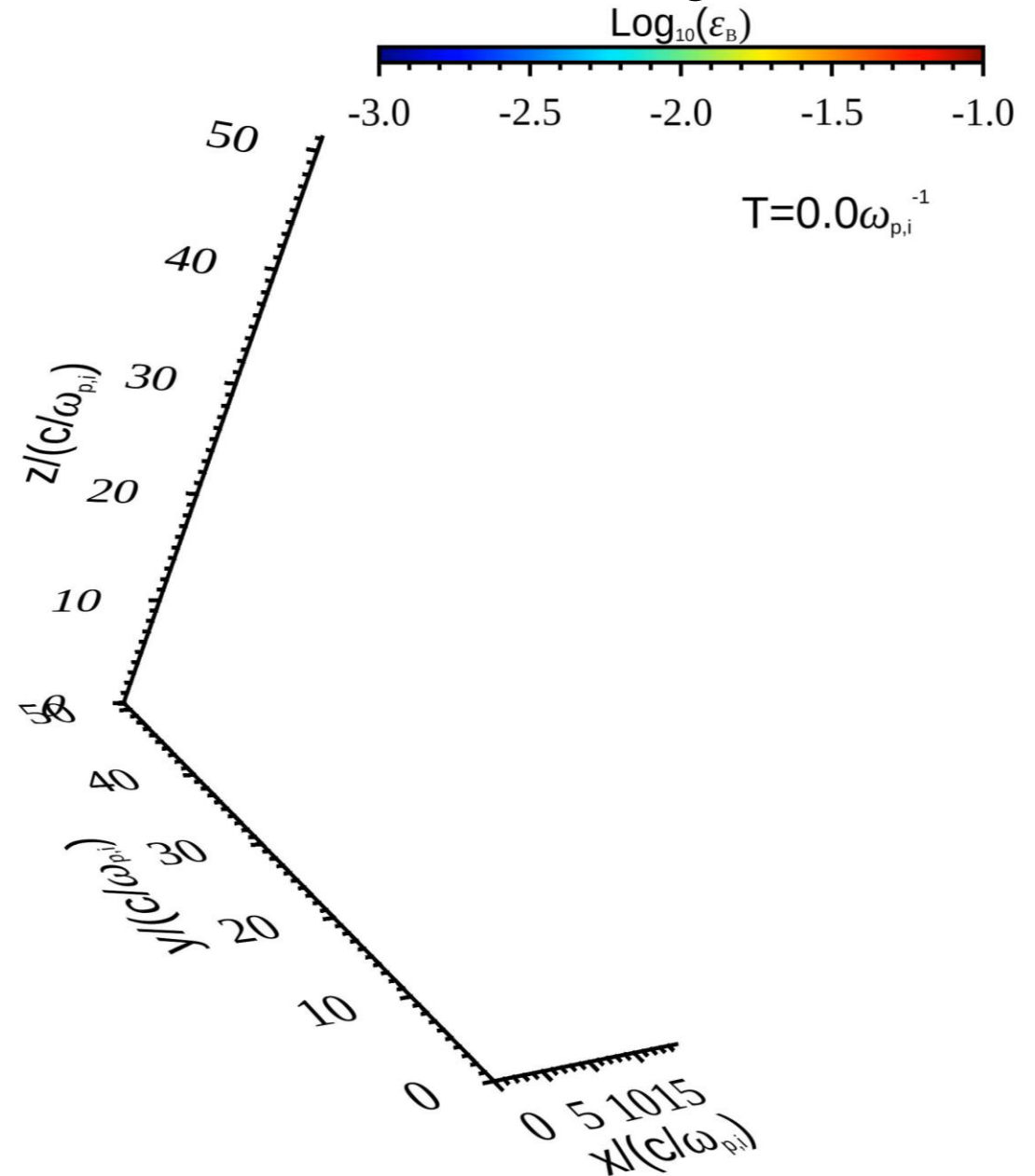
strong anisotropy  
(relativistic)

$\sigma_{\parallel} / \sigma_{\perp} - 1$   
Kato05 PoP

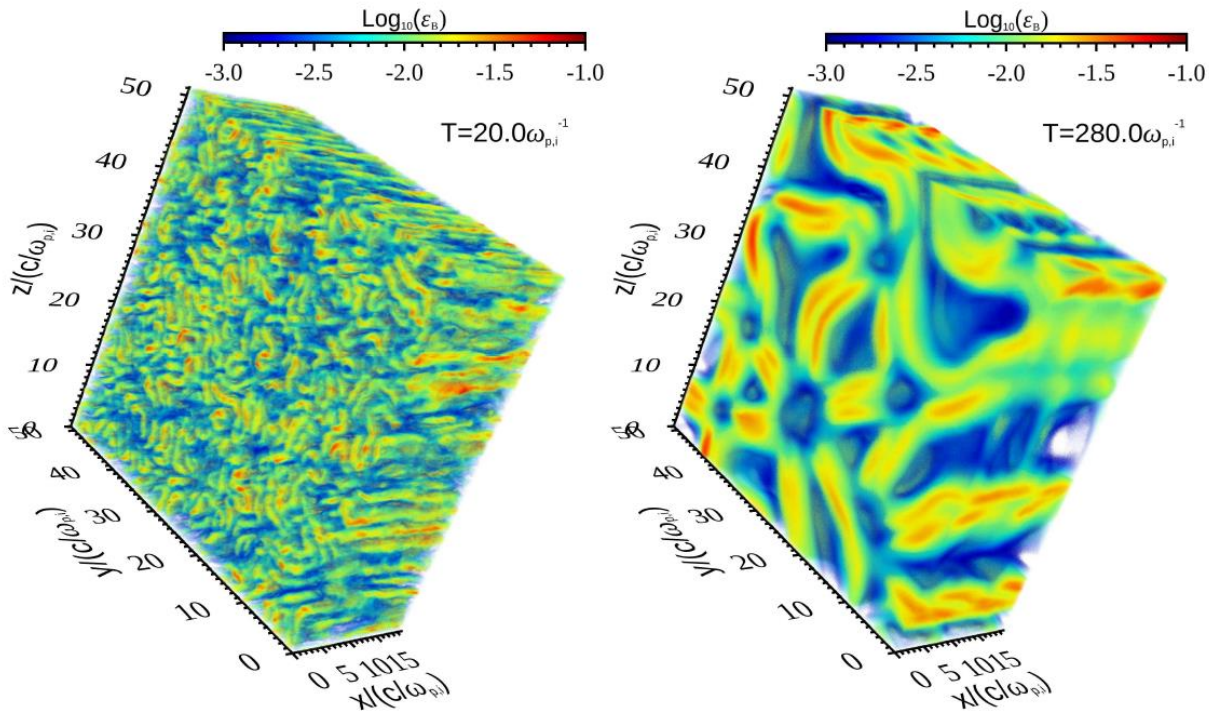


# Saturation of the Ion Weibel instability

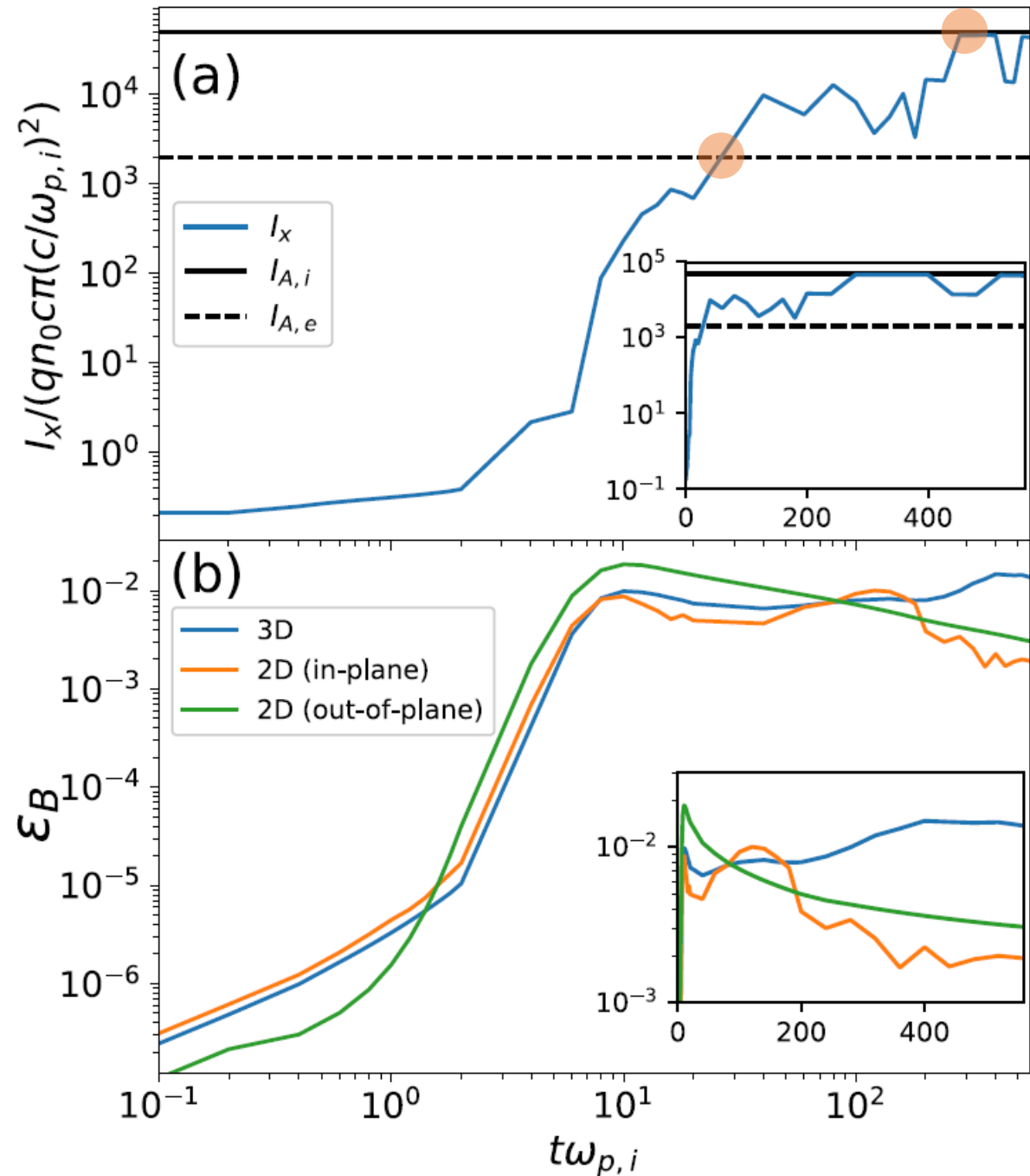
- $M/m=25$ , relativistic cold beams ( $\Gamma=5$ ) in a periodic box
- $(N_x, N_y, N_z) = (1000, 2560, 2560)$
- $(L_x, L_y, L_z) = (20, 51.2, 51.2) c/\omega_{pi}$
- 10 ptcls/cell/species
- 131,072 processor cores with 200TB memory on the K computer
- $\epsilon_B = B^2 / 8\pi\rho_0 c^2 (\gamma - 1) \Rightarrow$



# Time evolution



- First to present true saturation of the ion Weibel instability reaching Alfvén current limit
- 3D ion Weibel instability could sustain strong magnetic field ( $\epsilon_B \sim 1.5\%$ ) after saturation (contrary to the rapid decay in 2D)



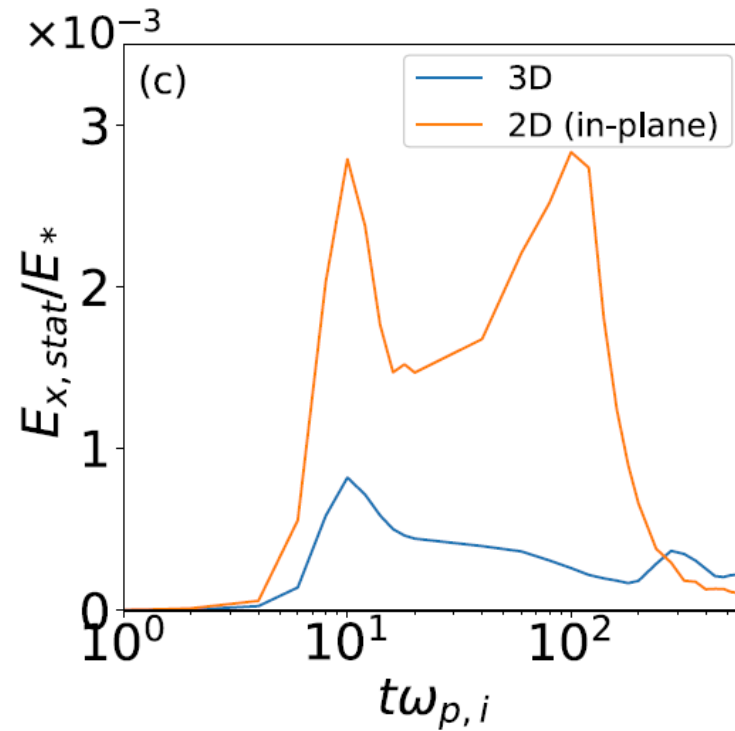
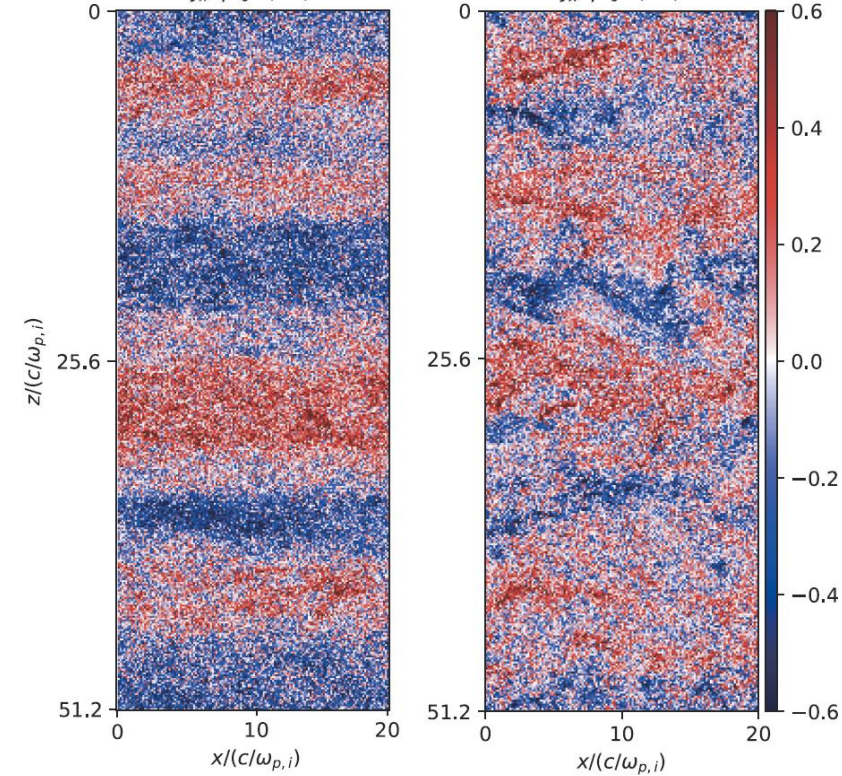
# 3D vs. 2D in-plane

3D

$J_x/qn_0c$  (3D)

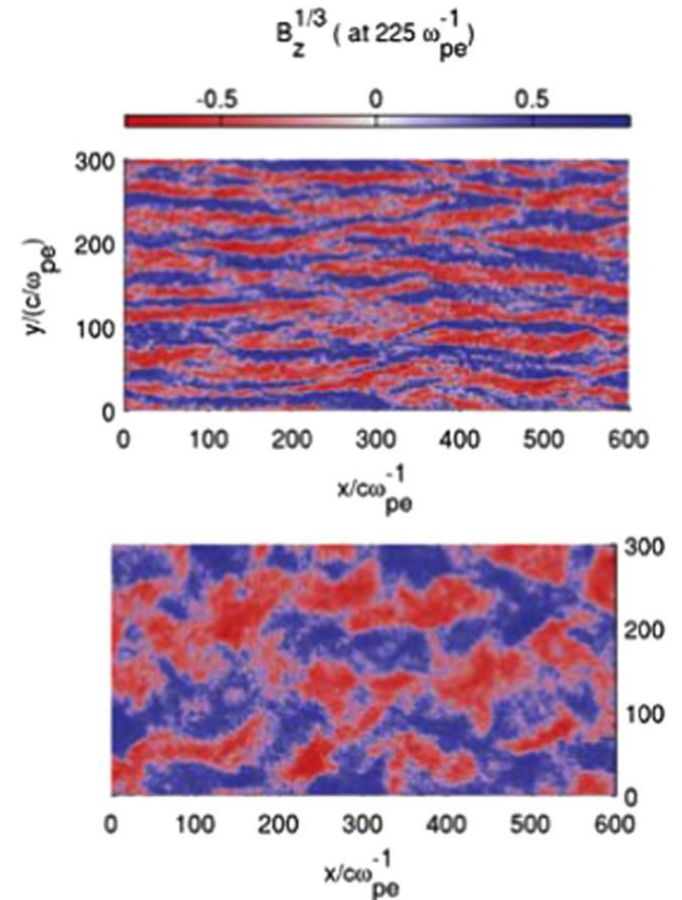
2D

$J_x/qn_0c$  (2D)



Takamoto+18ApJL

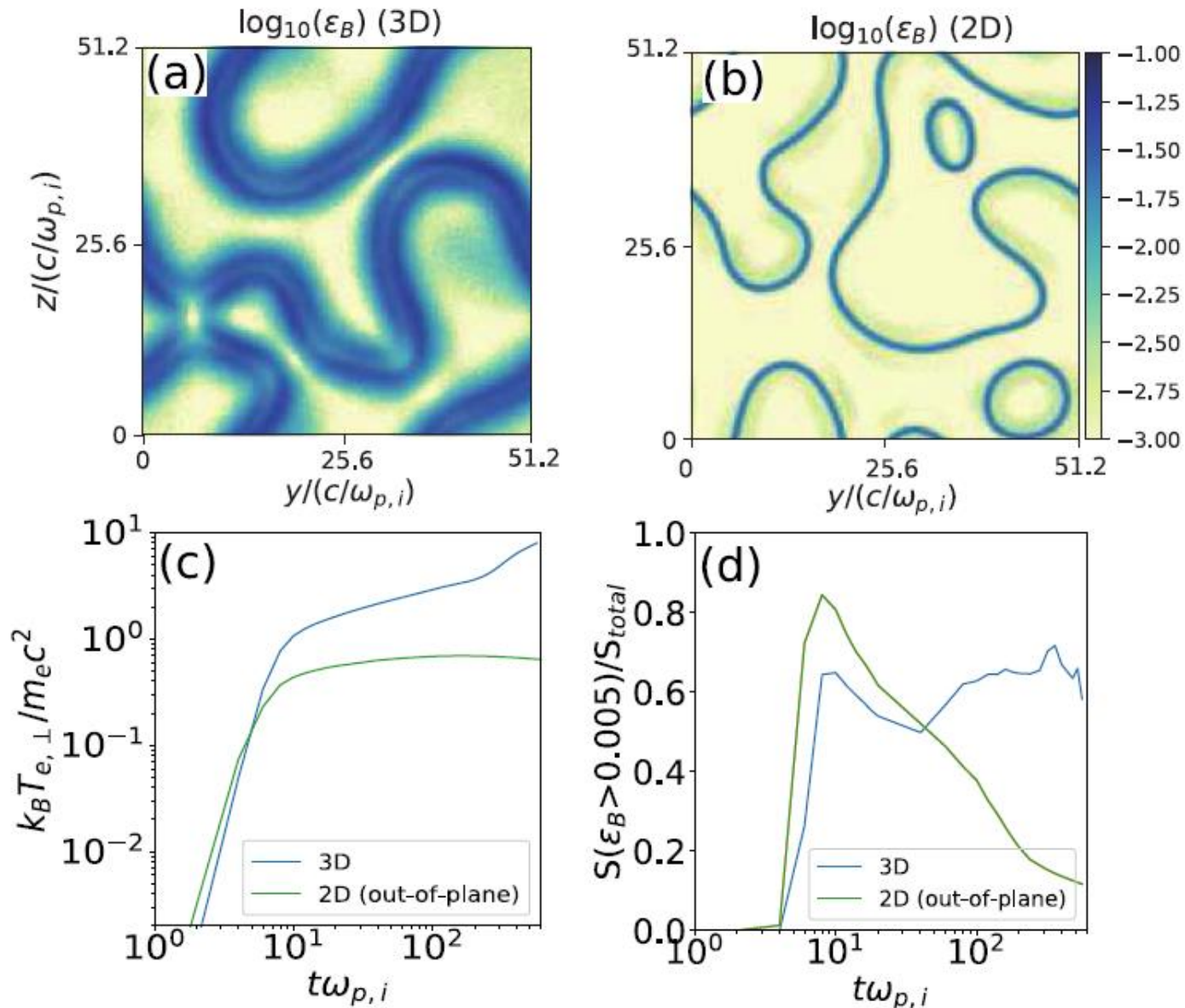
breakup of filaments  
in 2D in-plane case



Kumar+15ApJ

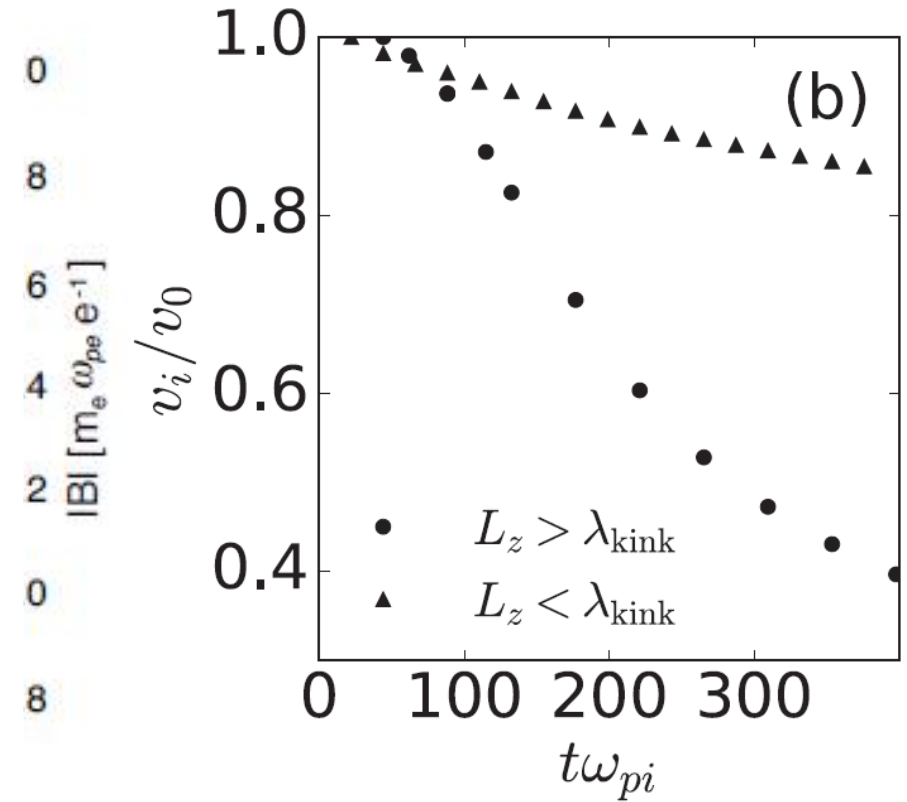
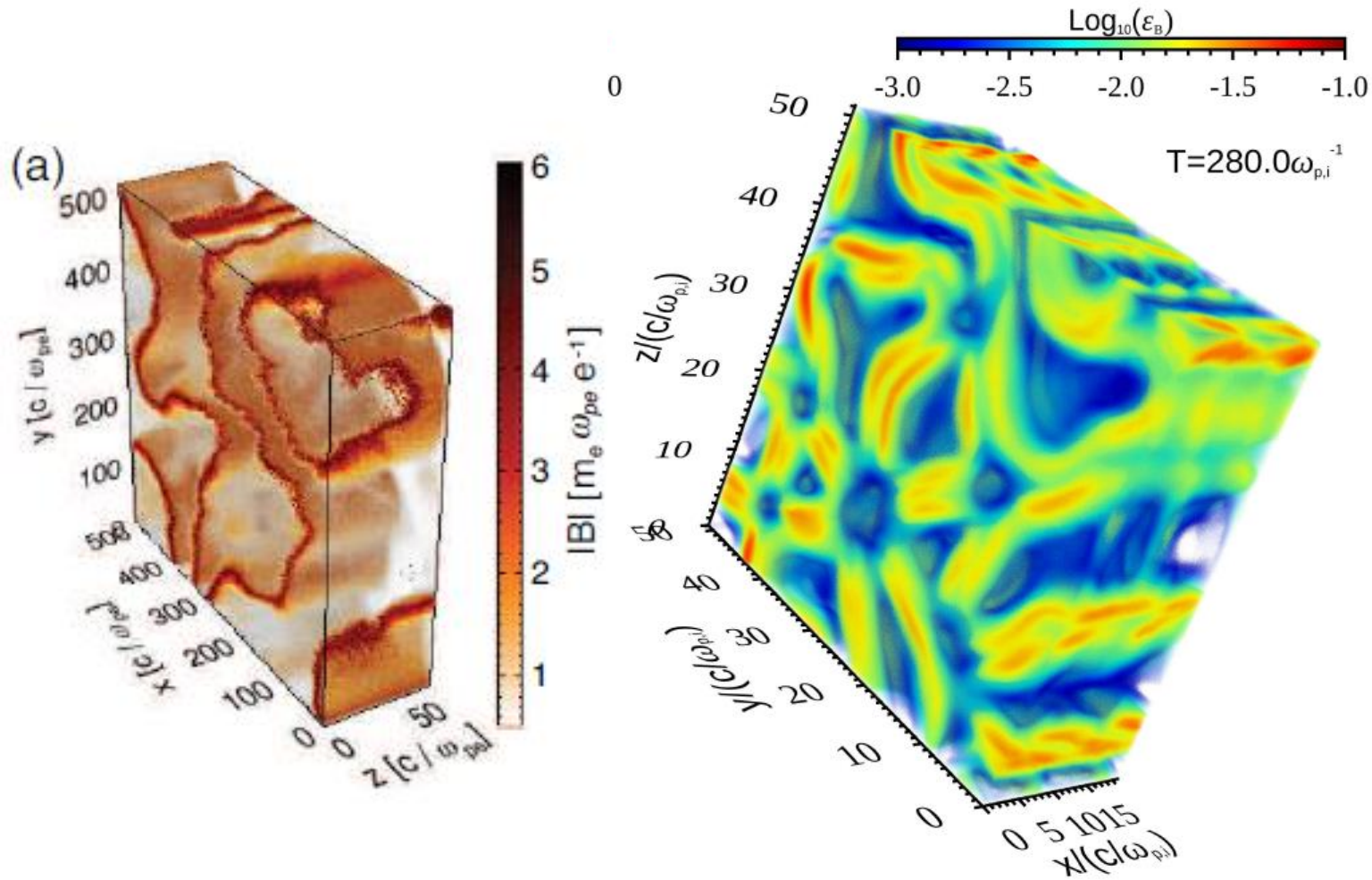
Dimensional limitation could cause two-stream electrostatic instabilities, resulting in breakup of current filaments in the 2D in-plane case (cf. Kumar+15)

# 3D vs. 2D out-of-plane



- Larger filling factor of large  $\epsilon_B$  regions in 3D
- Continuous electron heating during coalescence of filaments un-magnetized electrons with relativistic temperature ( $T_e \sim 10mc^2$ )
- Relativistically hot electrons contributed ion-scale magnetic field generation

# ion Weibel vs. Kink mode



# まとめ

- Ikeya & Matsumoto15で開発したマジカルCFL法で、実用上数値チェレンコフ不安定問題は解決し、相対論的衝撃波の多次元構造が急速に明らかになりつつある
- Precursor waveの相対論的衝撃波面からの放射は幅広い $\sigma$ で普遍的に存在している (Iwamoto+17; 18)
- イオン・電子系での航跡場加速 (Iwamoto+, in prep.)
- イオン・電子系での乱流リコネクション加速 ( $\sigma \ll 1$ , cf. Matsumoto+ 15 Science)
- 超相対論的衝撃波 ( $\Gamma > M/m$ ) では、イオン加速も
- 3次元計算もしたい…けど「京」は来年度中に運用停止

# Post-K supercomputer

- Next generation Japanese flagship supercomputer
- Present-K operation will terminate in FY2019
- Operation will start in 2021 (Approved to proceed to manufacturing phase from CPU design phase)
- No accelerator!
- CPU: Arm, 48cores, 512bitSIMD, 2.7TFlops
- Memory: HBM2 high band-width memory, but <1GB/core (32GiB/cpu)
- B/F=0.37 : high as GPUs
- Total Flops < 1Exa Flops, but x100 application speed up
- expected # of CPUs:  $1\text{EFlops}/2.7\text{TFlops} \sim 3 \times 10^5$  ( $\sim 10^7$ cores)

