

Impact of Lyman Alpha Pressure on Metal-Poor Dwarf Galaxies

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in collaboration with

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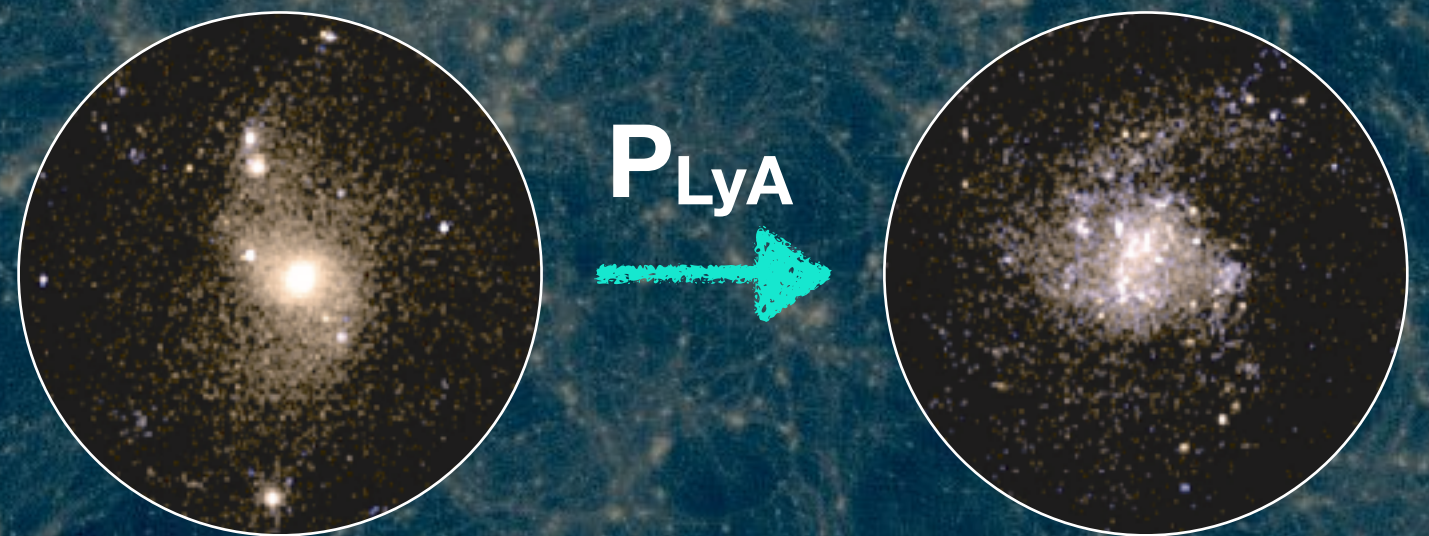
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Thibault **Garel** (Lyon)

Leo **Michel-Dansac** (Lyon)

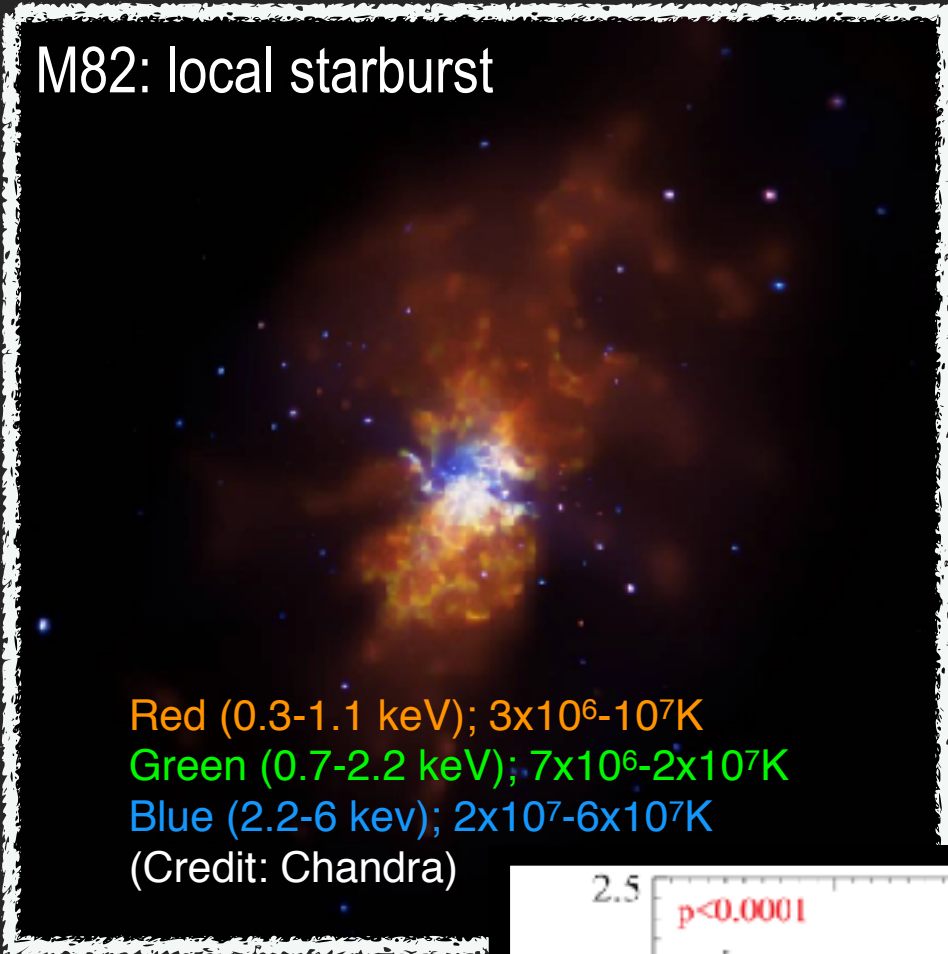
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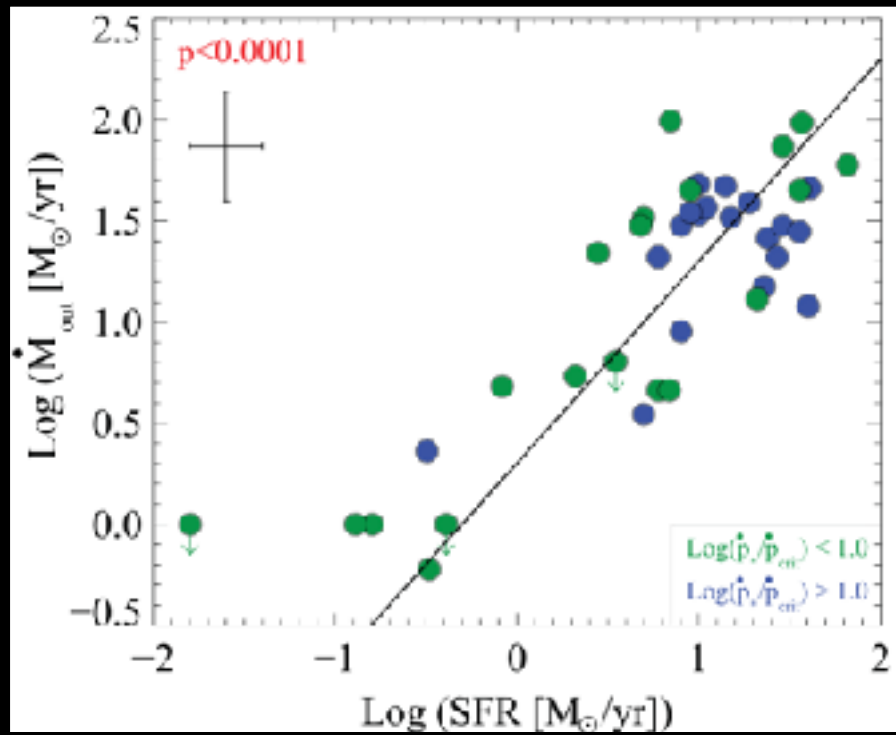


Common feature of SFRs - Outflows

M82: local starburst

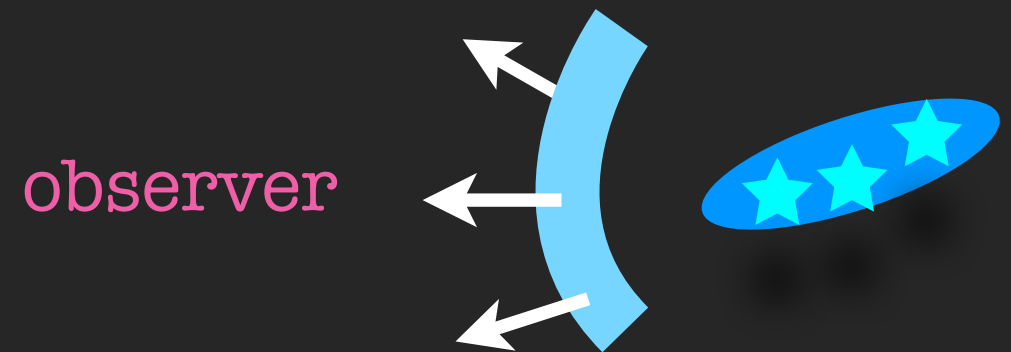
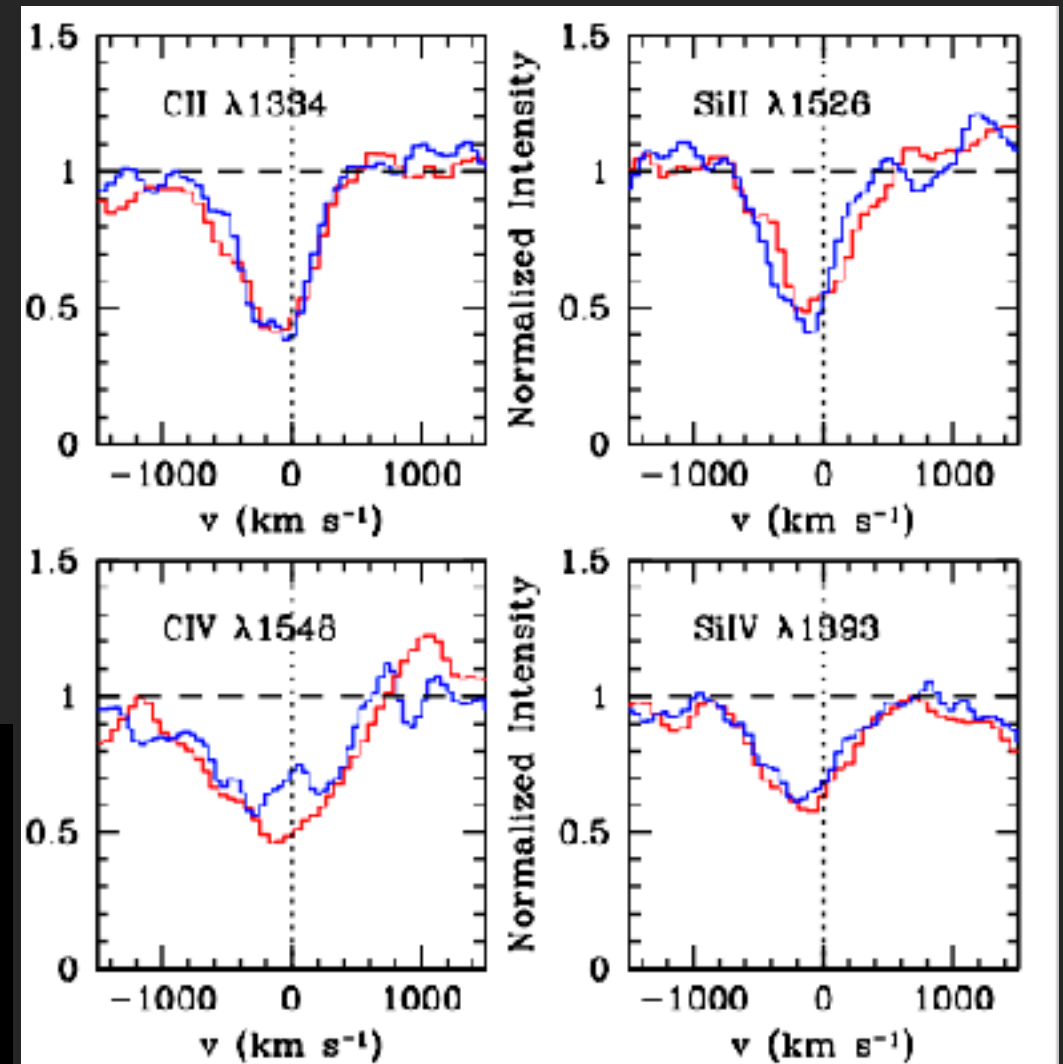


Red (0.3-1.1 keV); $3 \times 10^6 - 10^7 K$
 Green (0.7-2.2 keV); $7 \times 10^6 - 2 \times 10^7 K$
 Blue (2.2-6 keV); $2 \times 10^7 - 6 \times 10^7 K$
 (Credit: Chandra)



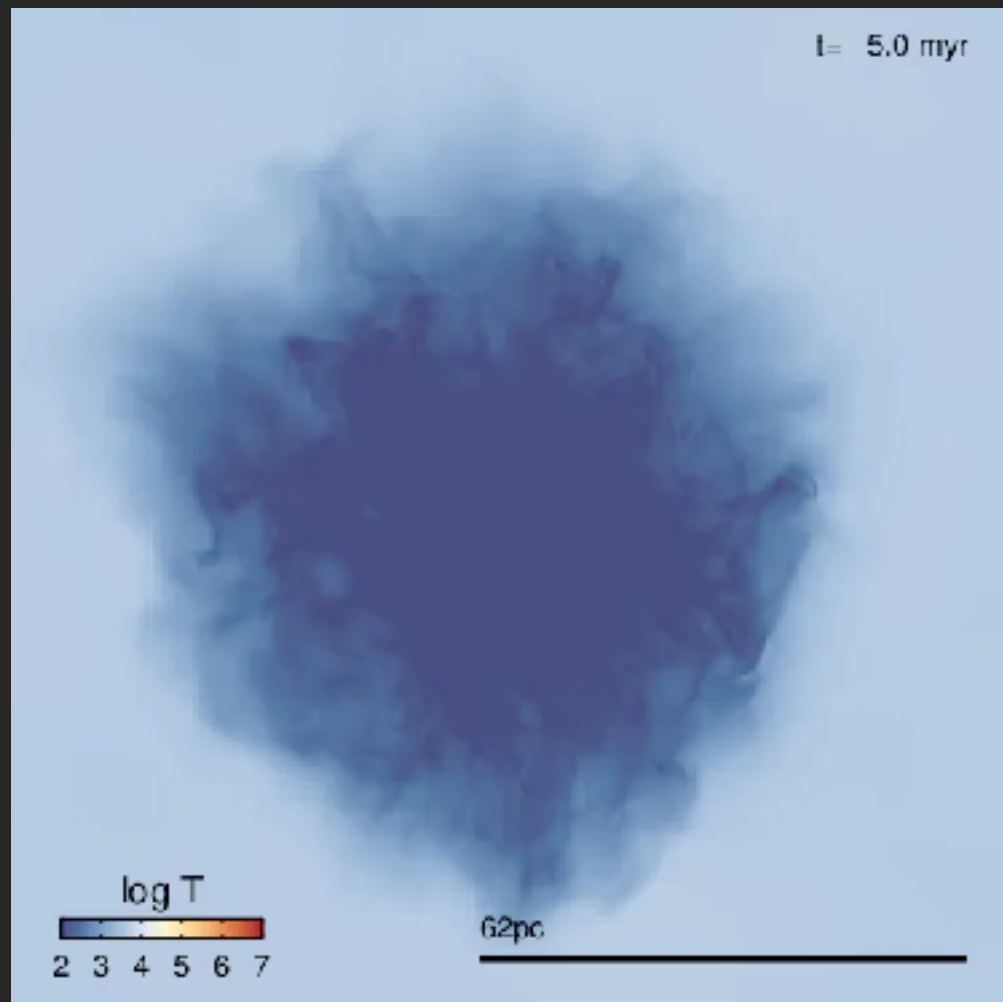
Heckman+(15; z~0 local starburst)

Ubiquitous **Outflows** in LBGs at z~2 (Steidel+10)



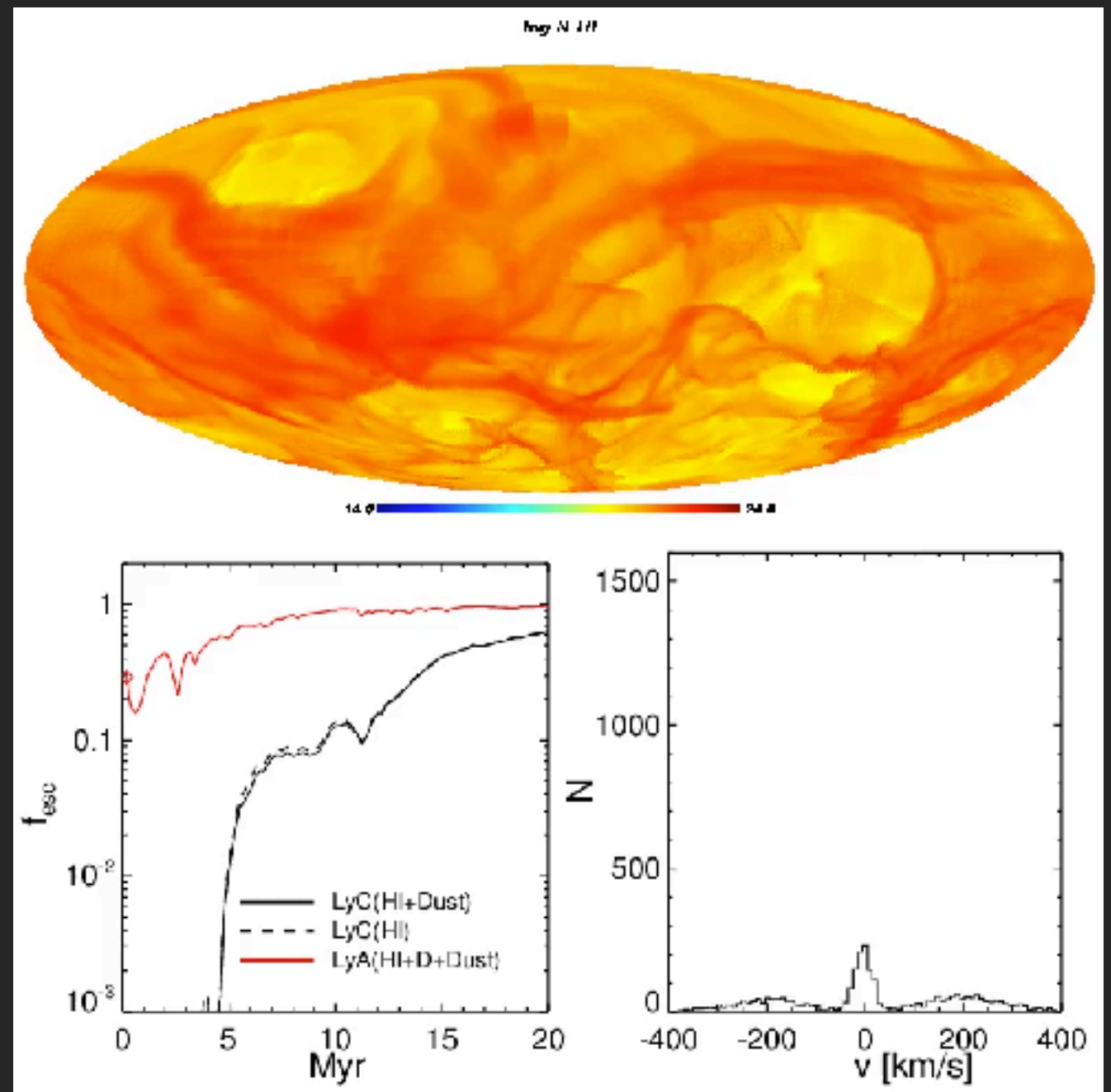
Lyman alpha profile on cloud scales

c.f. Neufeld (90;91); Zheng+(02); Verhamme et al. (06,12);
Dijkstra et al. (06); Gronke & Dijkstra (15)



Radiation-Hydrodynamic simulations of a
turbulent cloud with
 $M_{\text{cloud}} = 10^6 M_{\text{sun}}$, $\text{SFE} = 1\%$

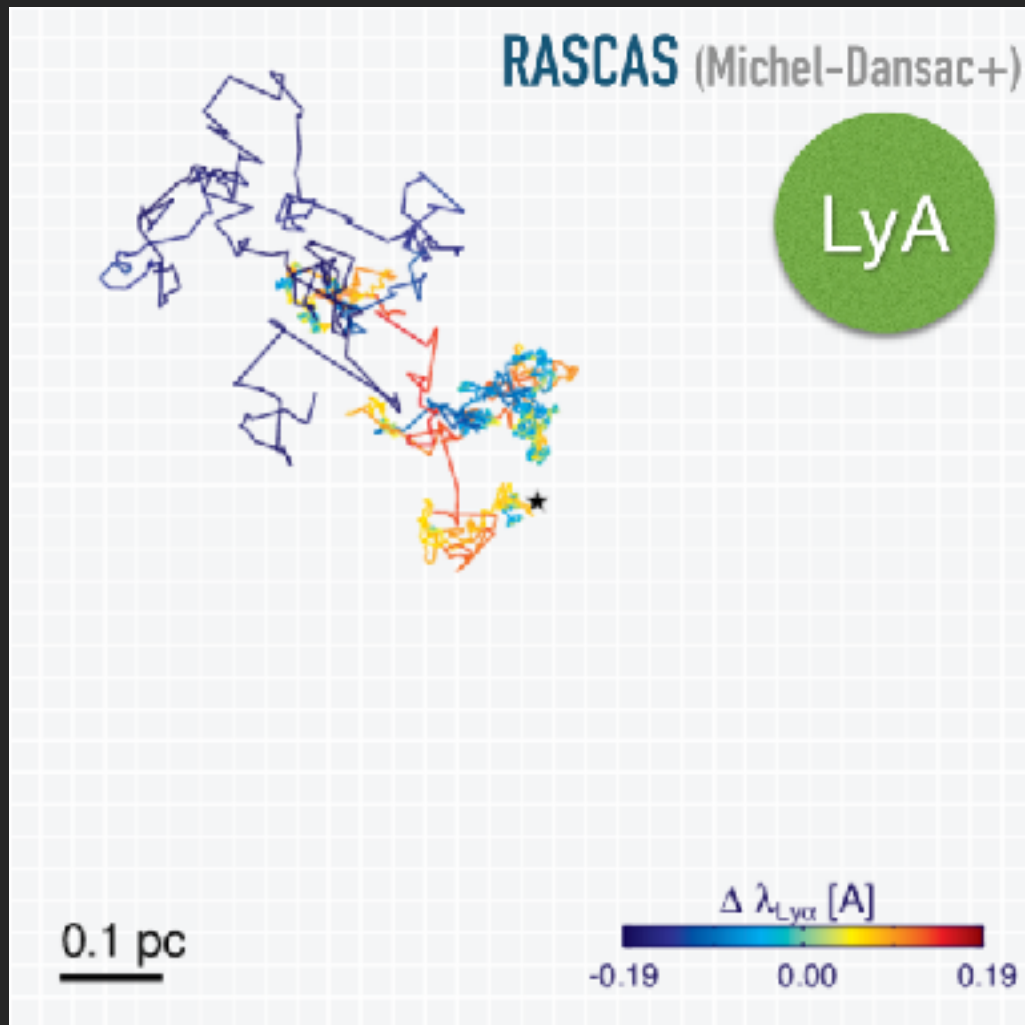
POST-PROCESSING WITH **RASCAS** (Jeremy Blaizot's talk)



recombination+collisional radiation

Kimm, Blaizot et al. (in prep)

Scattering of Ly α photons transfers momentum to the surroundings



WHY DO WE CARE?

Optical depth to Ly α is huge!

$$\tau_{Ly\alpha} = \left(\frac{N_{HI}}{2 \times 10^{13} \text{ cm}^{-2}} \right) \left(\frac{T}{10^4 \text{ K}} \right)^{-1/2}$$

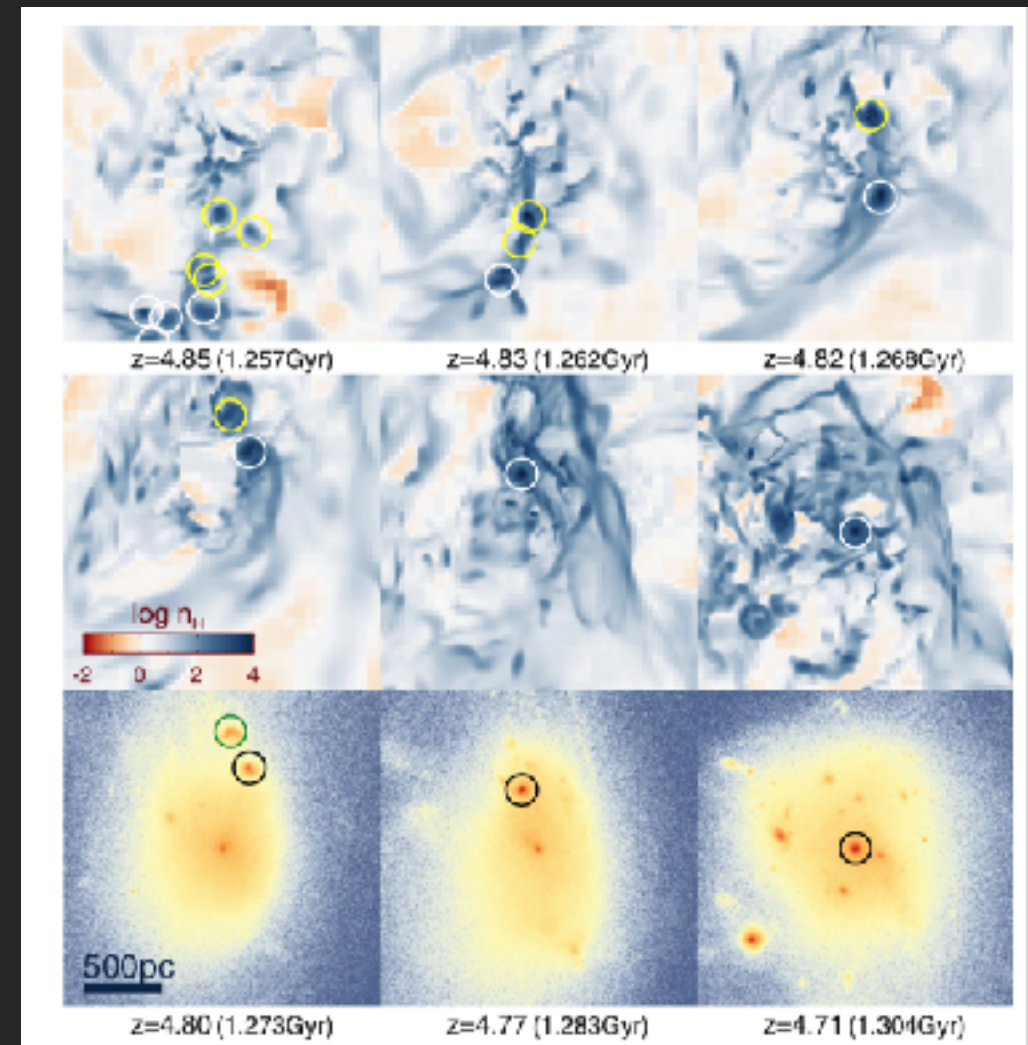
Solution to the Over-cooling problem at high z ?

Momentum transfer

$$\Delta \vec{p} = \frac{h\nu}{c} (\nu_{in} \hat{n}_{in} - \nu_{out} \hat{n}_{out})$$

Multiplication factor

$$F_{Ly\alpha} = M_F \frac{L_{Ly\alpha}}{c}$$



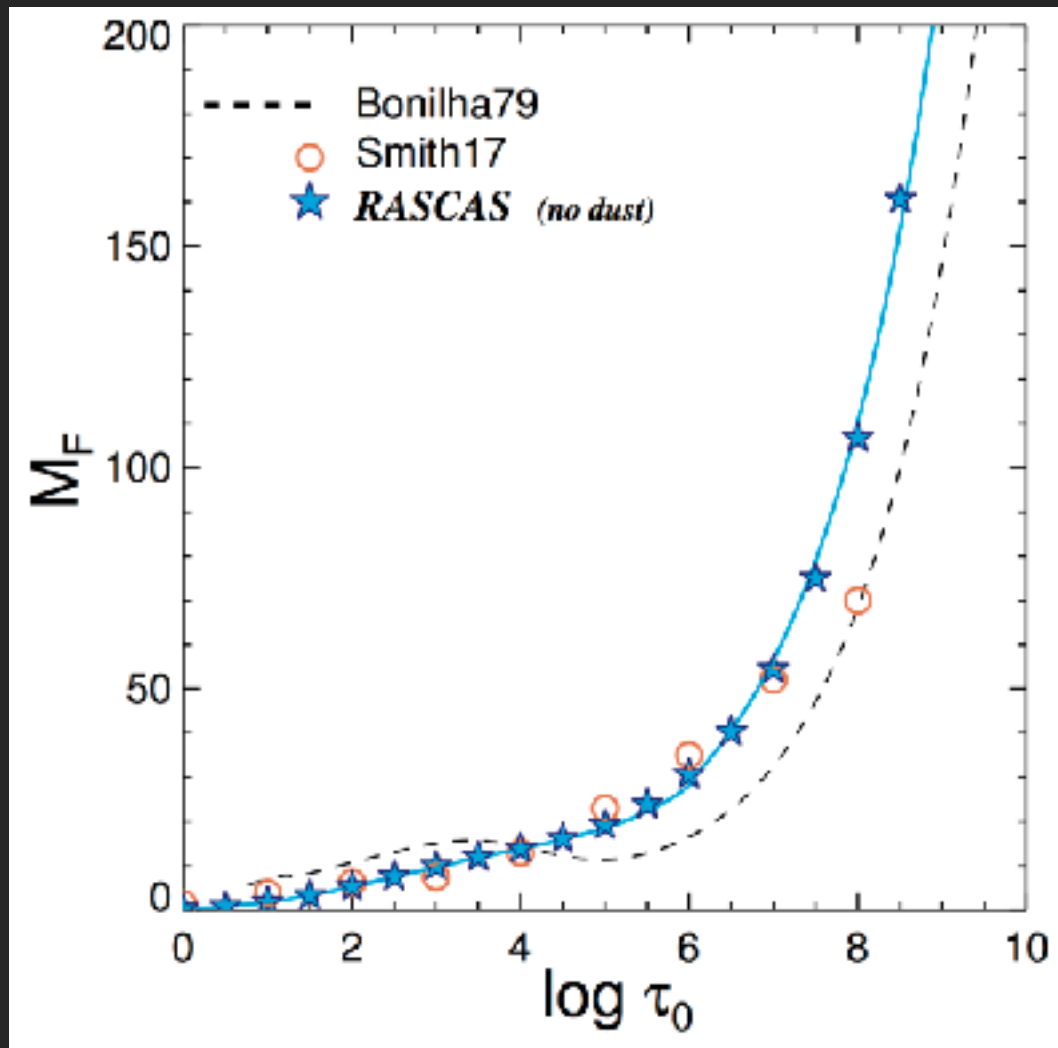
Kimm, Cen, Devriendt+(15)

See also Dijkstra & Loeb (08), Smith et al. (16)

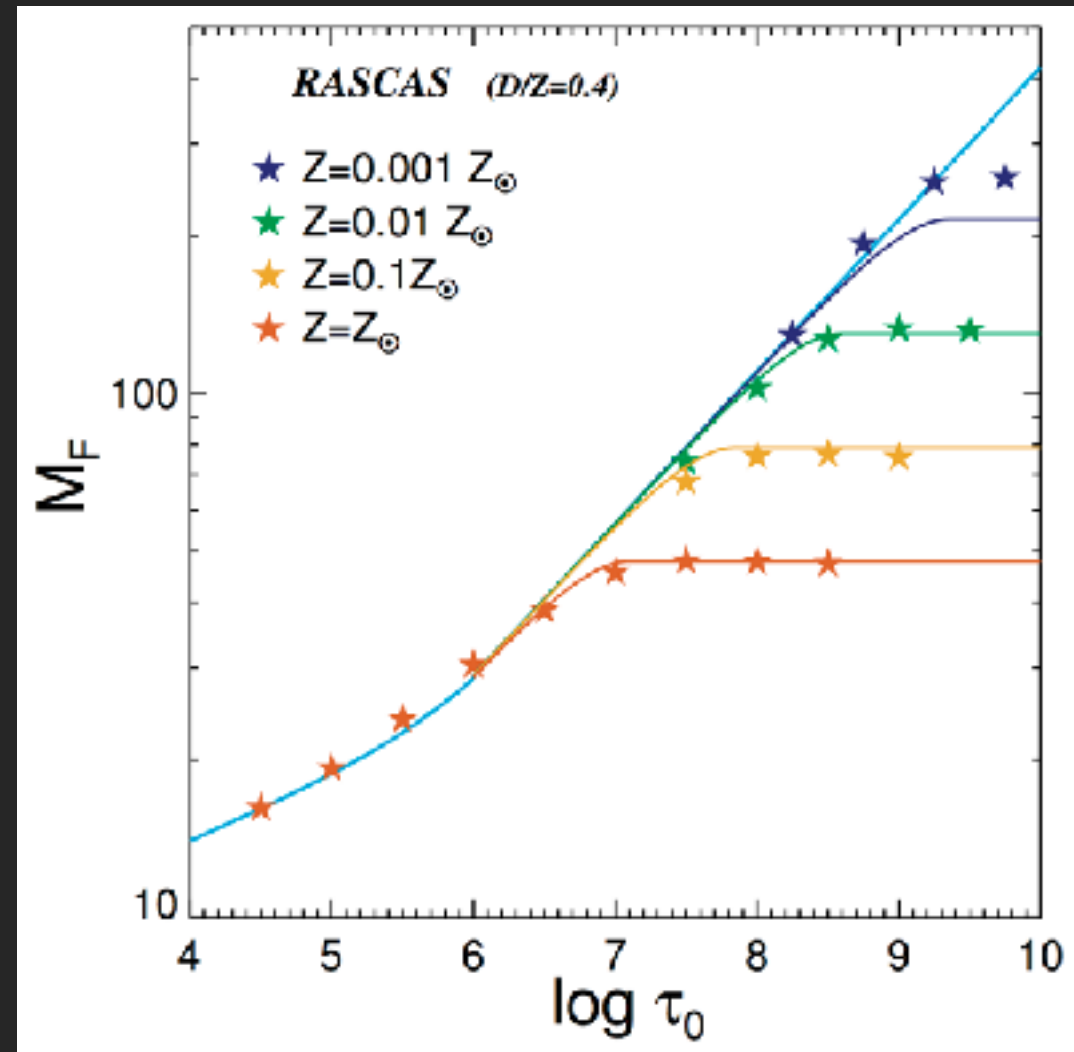
Scattering of Ly α photons transfers momentum to the surroundings

See also Dijkstra & Loeb (08), Smith et al. (17)

Dust-free



Dusty



Multiplication factor

$$F_{Ly\alpha} = M_F \frac{L_{Ly\alpha}}{c}$$

Dust-to-metal ratio: Remy-Ruyer+(14)

$$\log f_{d/m} = 0 \quad (x > 8.10),$$

$$= 1.25 - 2.10 \times (x_{\odot} - x) \quad (x \leq 8.10),$$

$$x \equiv 12 + \log(O/H) \text{ and } x_{\odot} = 8.69$$

Momentum budget from Ly α pressure

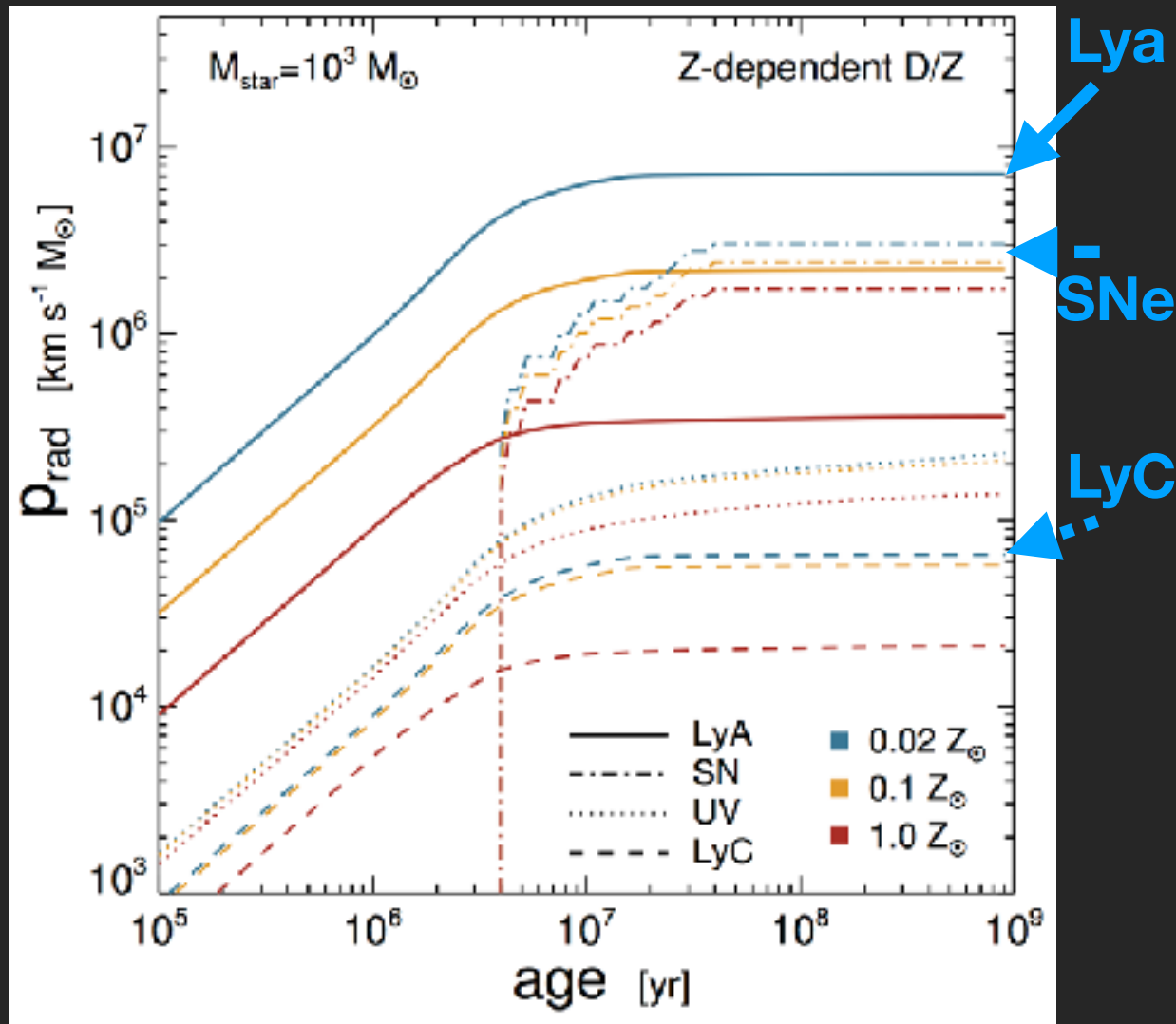


Photo-ionization heating

$$n_{\text{H,ion}} T_{\text{ion}} = n_{\text{H,0}} T_0$$

$$r_{\text{PH}} \approx 26 \text{ pc} \left(\frac{m_{\text{star}}}{10^3 M_{\odot}} \right)^{1/3} \left(\frac{P/k_B}{10^5 \text{ cm}^{-3} \text{ K}} \right)^{-3/2} \left(\frac{T_{\text{ion}}}{10^4 \text{ K}} \right)^{2/3}$$

Lyman alpha pressure

$$n_{\text{H}} k_B T = \frac{M_F L_{\alpha}}{4\pi r_{\alpha}^2 c}$$

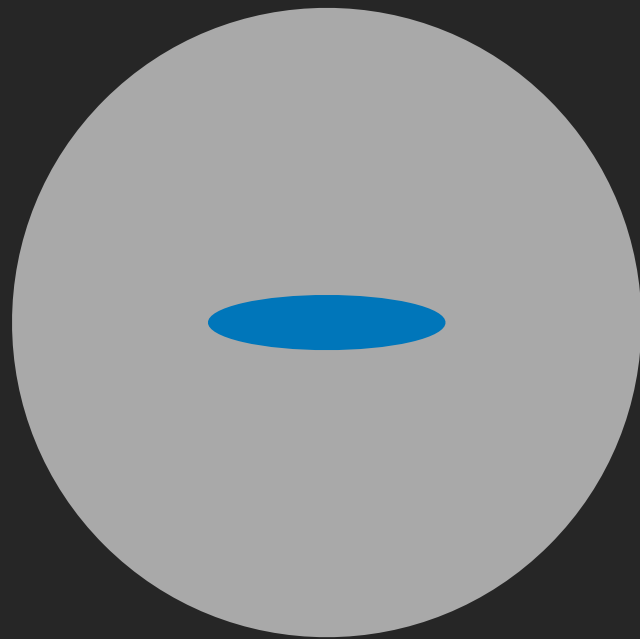
$$r_{\alpha} = 37 \text{ pc} \left(\frac{M_F}{100} \right)^{1/2} \left(\frac{m_{\text{star}}}{10^3 M_{\odot}} \right)^{1/2} \left(\frac{P/k_B}{10^5 \text{ cm}^{-3} \text{ K}} \right)^{-1/2}$$

- The momentum from Ly α is comparable or more significant than that of SNe
- Ly α pressure is advantageous from a computational viewpoint as well

Radiation-hydrodynamic simulations of an isolated disk

Simulation set-up

$$M_{\text{DMH}} \sim 10^{10} M_{\text{sun}}$$



$$Z = 0.02 Z_{\text{sun}}$$

$$M_{\text{star}} = 2 \times 10^8 M_{\text{sun}}$$

$$M_{\text{gas}} = 1.7 \times 10^8 M_{\text{sun}}$$

Max Resolution: 2 - 5 pc

Input physics

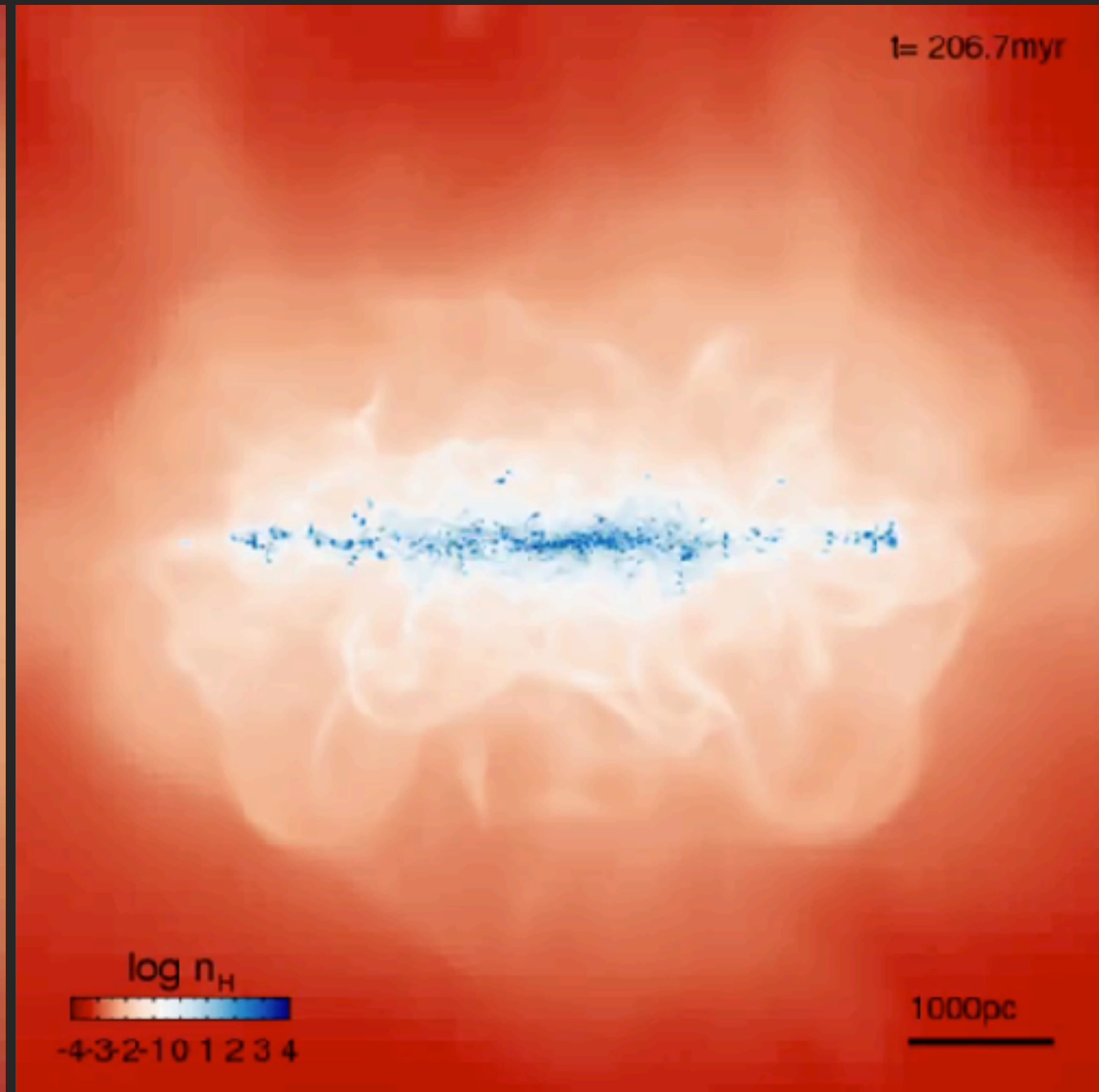
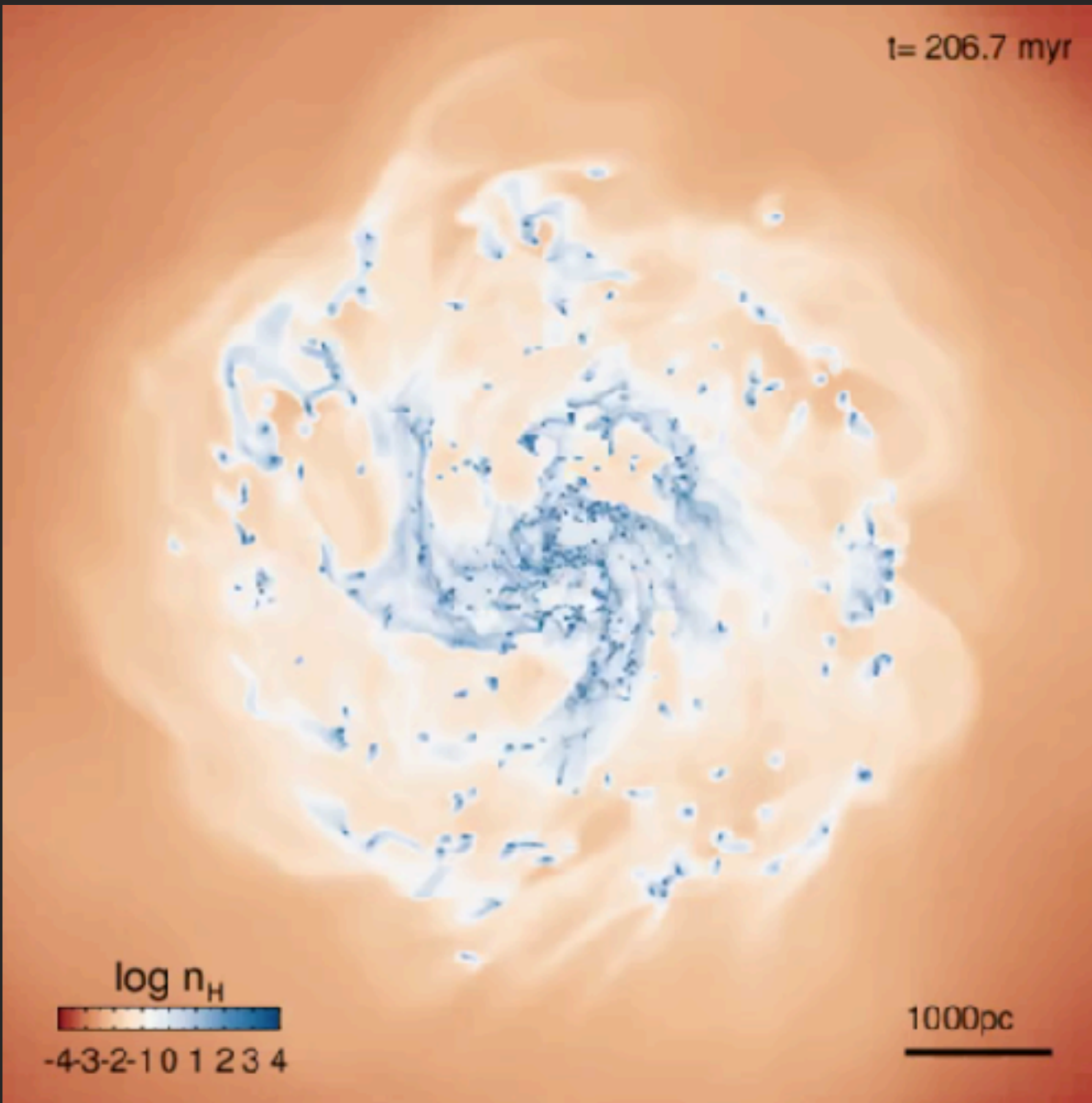
RAMSES-RT (Teyssier 02; Rosdahl+13)

- Thermo-turbulent star formation scheme (Kimm+17)
- Momentum-conserving SNe (Kimm & Cen 14, Kimm+15)
- Non-equilibrium photo-chemistry with H₂ (Katz, Kimm+17)
- Photo-ionisation heating (Rosdahl+13)
- Direct radiation pressure (Rosdahl+13)
- RP by reprocessed IR photons (Rosdahl & Teyssier 15)
- Photoelectric heating on dust (Kimm+17)
- Ly α pressure (Kimm+18)

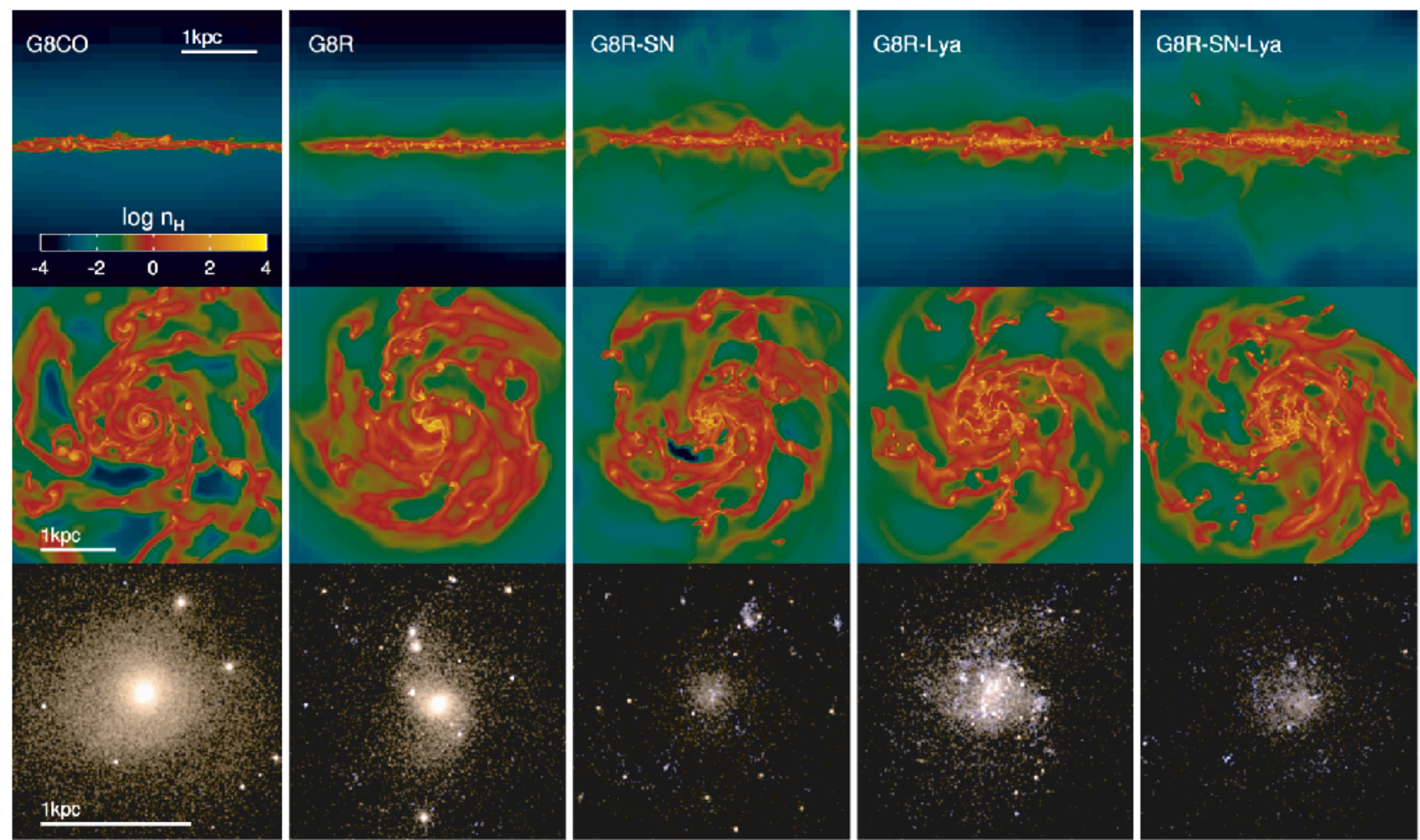
Photon group	ϵ_0 [eV]	ϵ_1 [eV]	κ [cm ² /g]	Main function
EUV _{HeII}	54.42	∞	10 ³	HeII ionisation
EUV _{HeI}	24.59	54.42	10 ³	HeI ionisation
EUV _{III,2}	15.2	24.59	10 ³	HI and H ₂ ionisation
EUV _{HI,1}	13.6	15.2	10 ³	HI ionisation
LW	11.2	13.6	10 ³	H ₂ dissociation
FUV	5.6	11.2	10 ³	Photoelectric heating
Optical	1.0	5.6	10 ³	Direct RP
IR	0.1	1.0	5	Radiation pressure (RP)

Radiation-hydrodynamic simulations of an isolated disk

PHOTOHEATING+DIRECT RADIATION PRESSURE +IR PRESSURE
+SN EXPLOSIONS + Lya PRESSURE

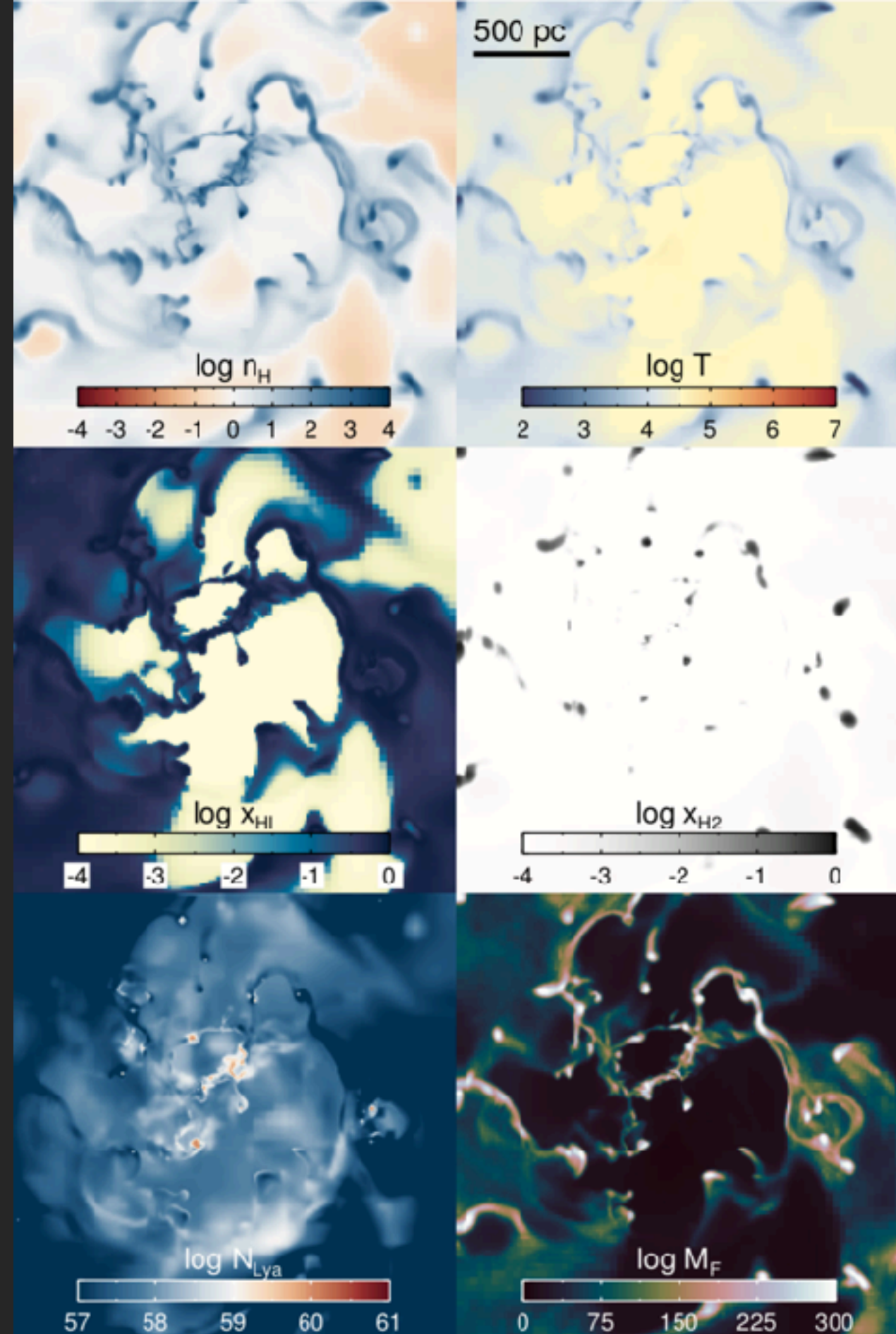
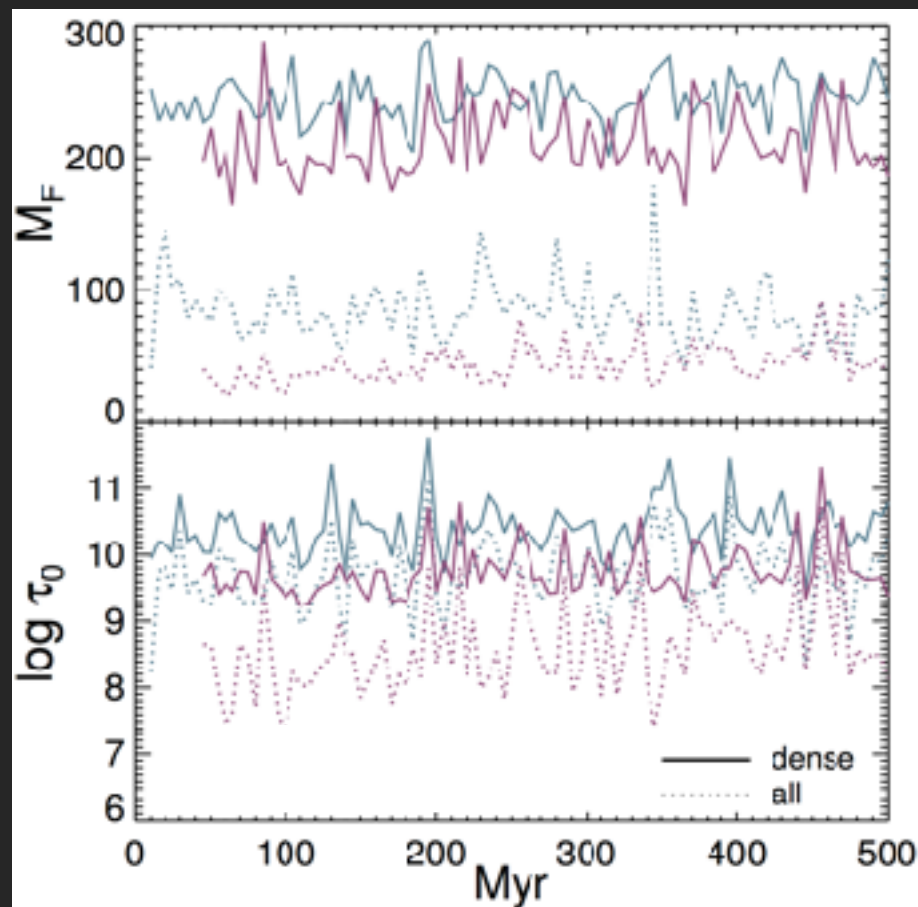


Radiation-hydrodynamic simulations of an isolated disk

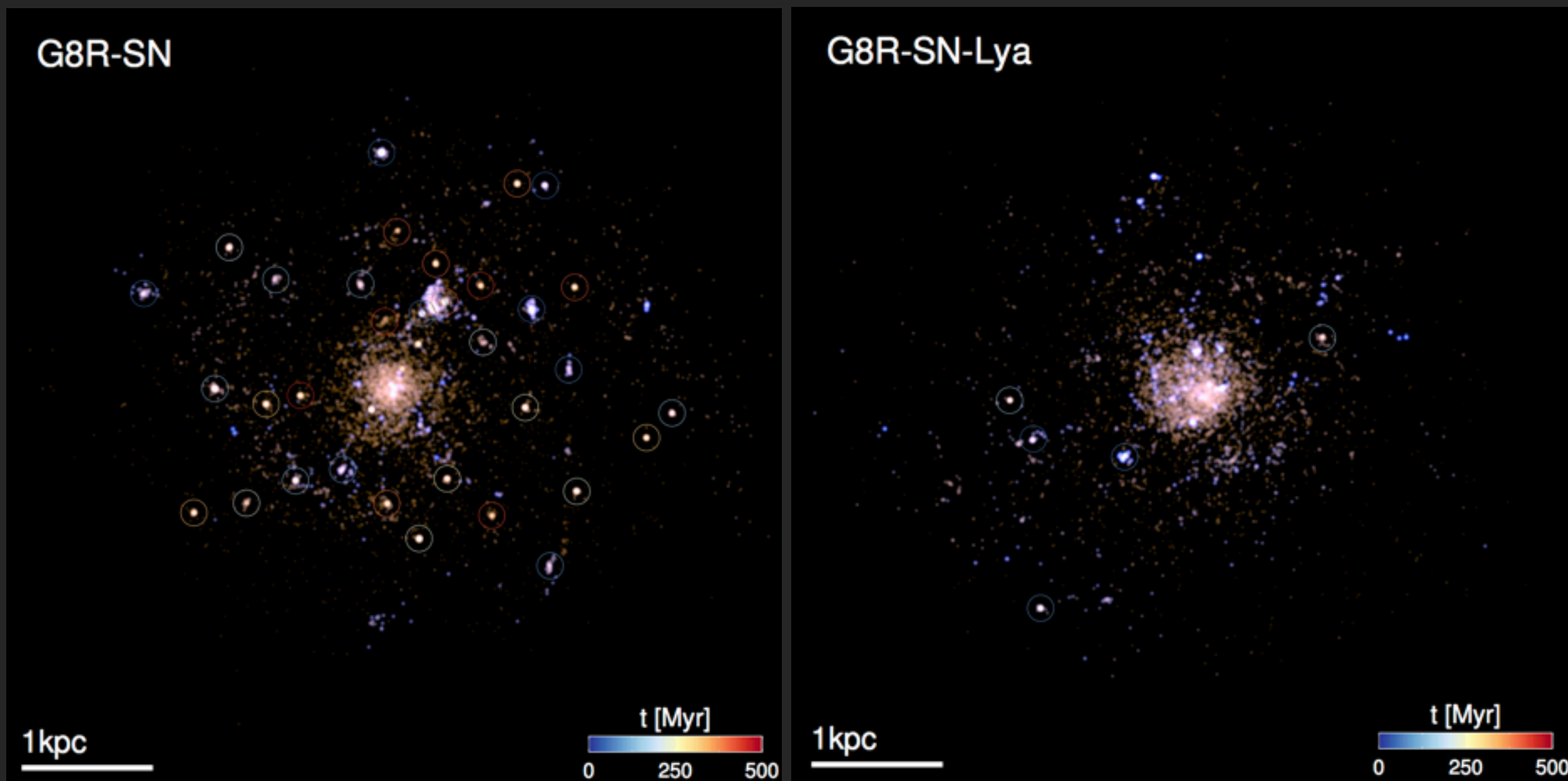


Where does Lya operate?

- Requirement for strong Lya pressure
 - Luminous ionizing source
 - Large N_{HI} density
- **AROUND YOUNG STARS**
- **INTERRUPT SF QUICKLY (<5MYR)**
- Effective $M_F \sim 200\text{-}300$ in dense regions



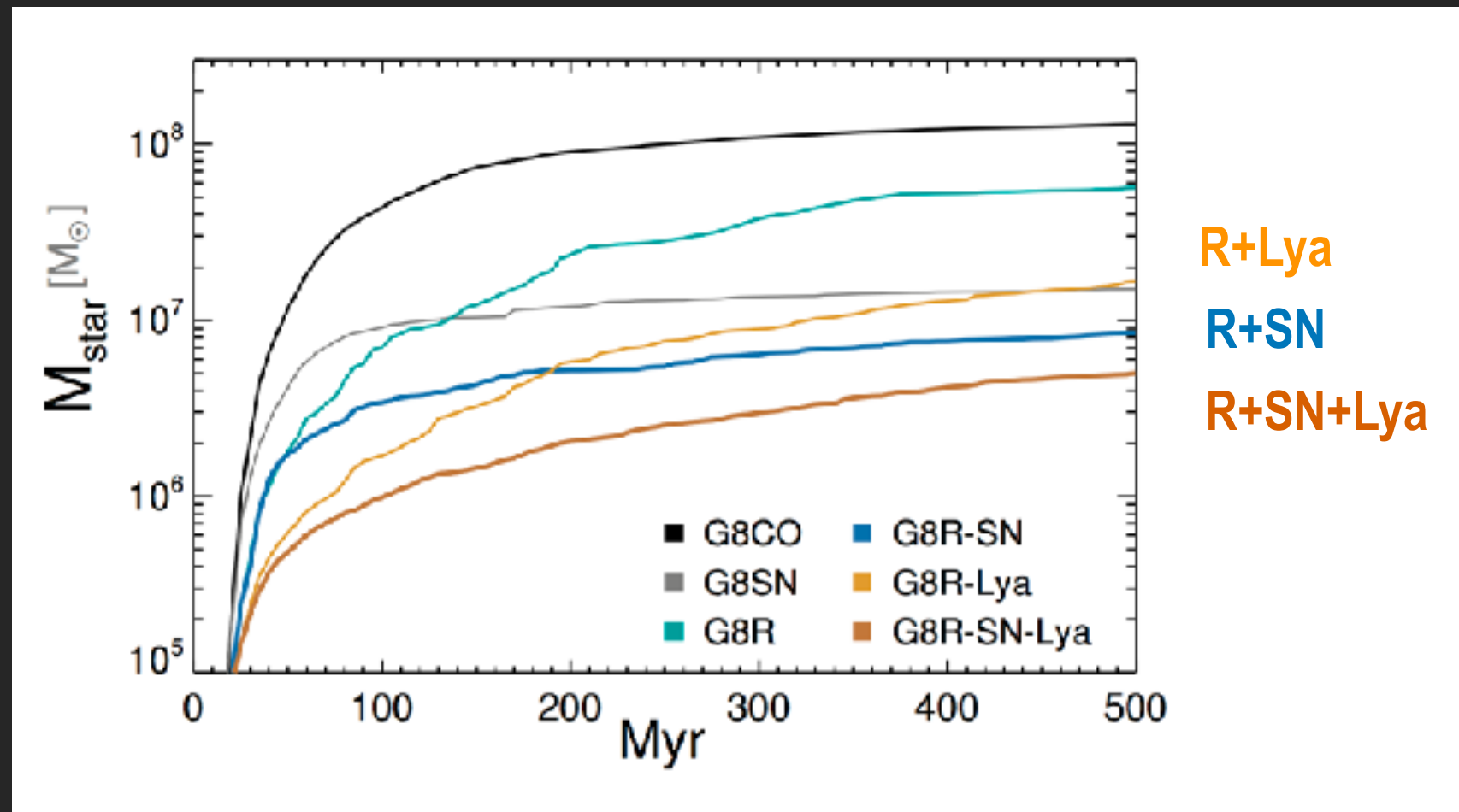
Cluster Formation with LyA Feedback



- Fewer clusters form and survive when strong radiation feedback is present (caution: cluster formation in HD simulations...)

[see also Abe & Yajima 18]

Star formation histories of a gas-rich, metal-poor dwarf

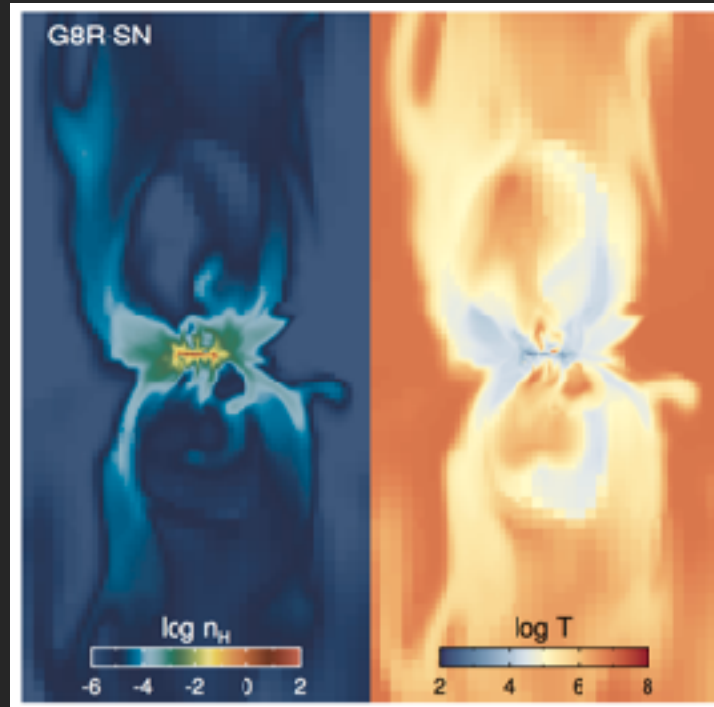


- Suppression of Star formation

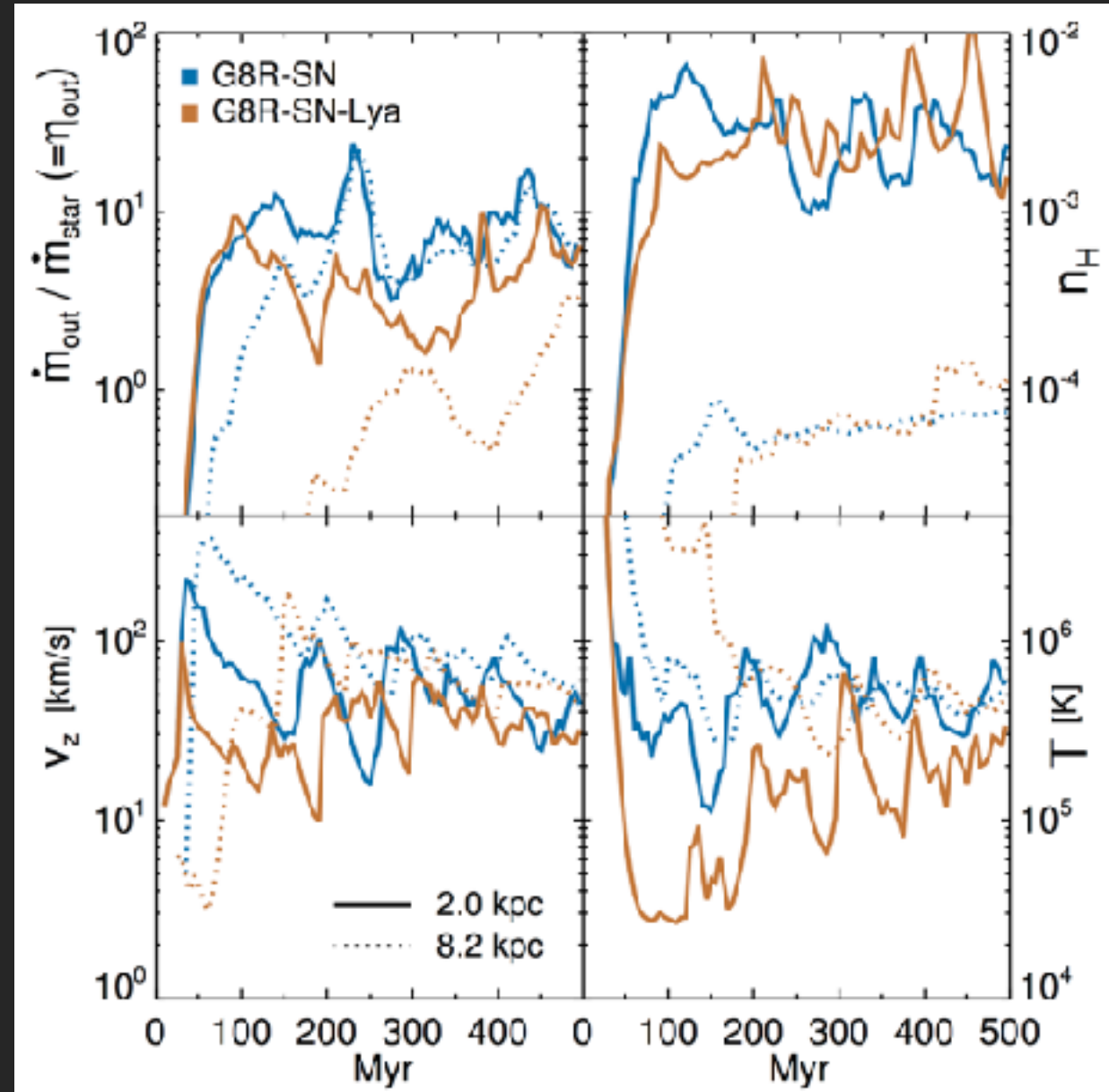
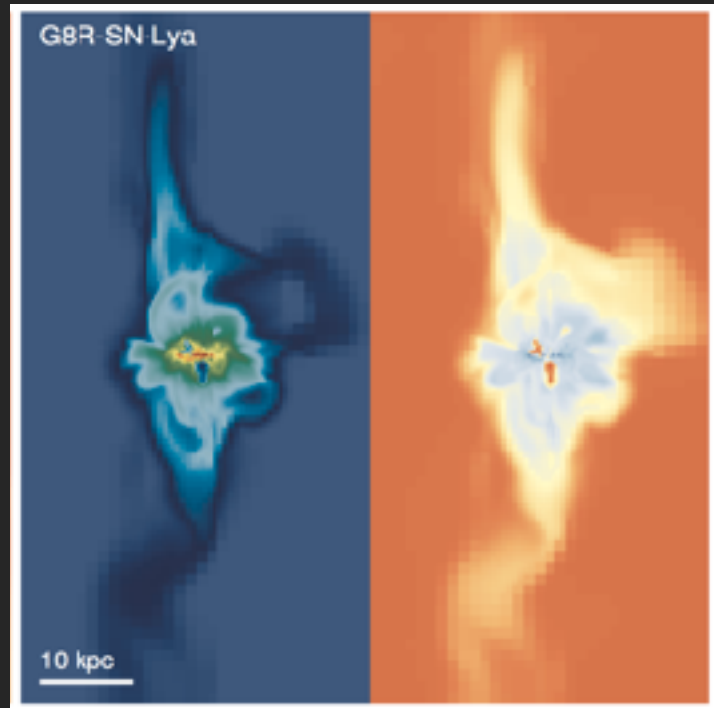
$$\text{LYA} < \text{SN} < \text{SN+LYA}$$

Weaker outflows with Lya pressure

w/o Lya



w/ Lya

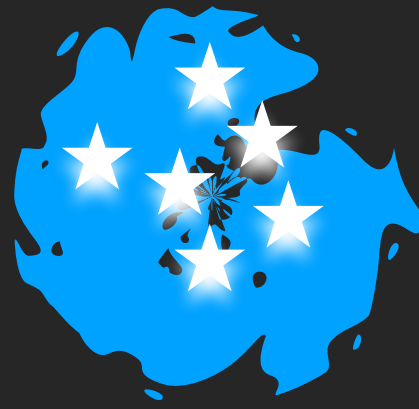


WITH LYA PRESSURE

- Mass-loading factor is decreased
- Outflows become cooler and slower

A picture with strong radiation feedback

No or Weak Radiation Feedback



Coherent Supernova Feedback

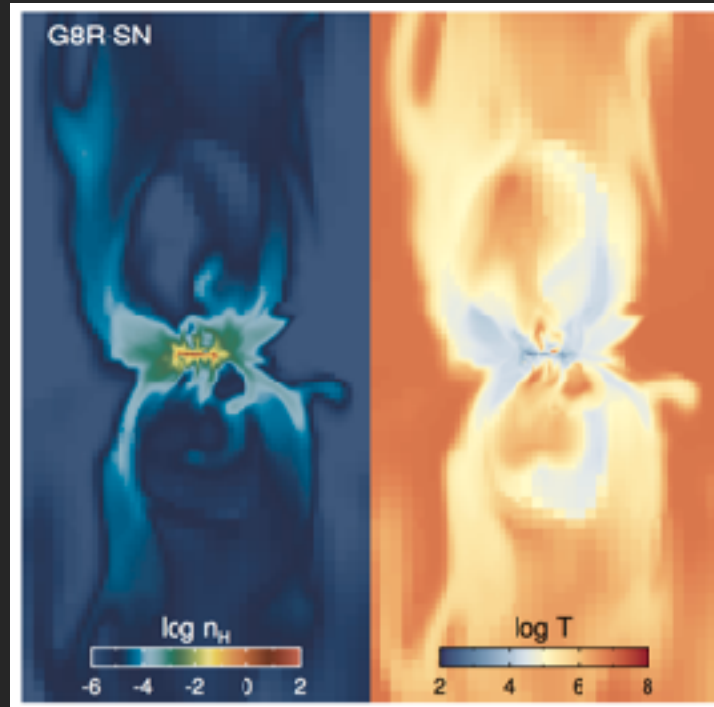
Strong Radiation Feedback



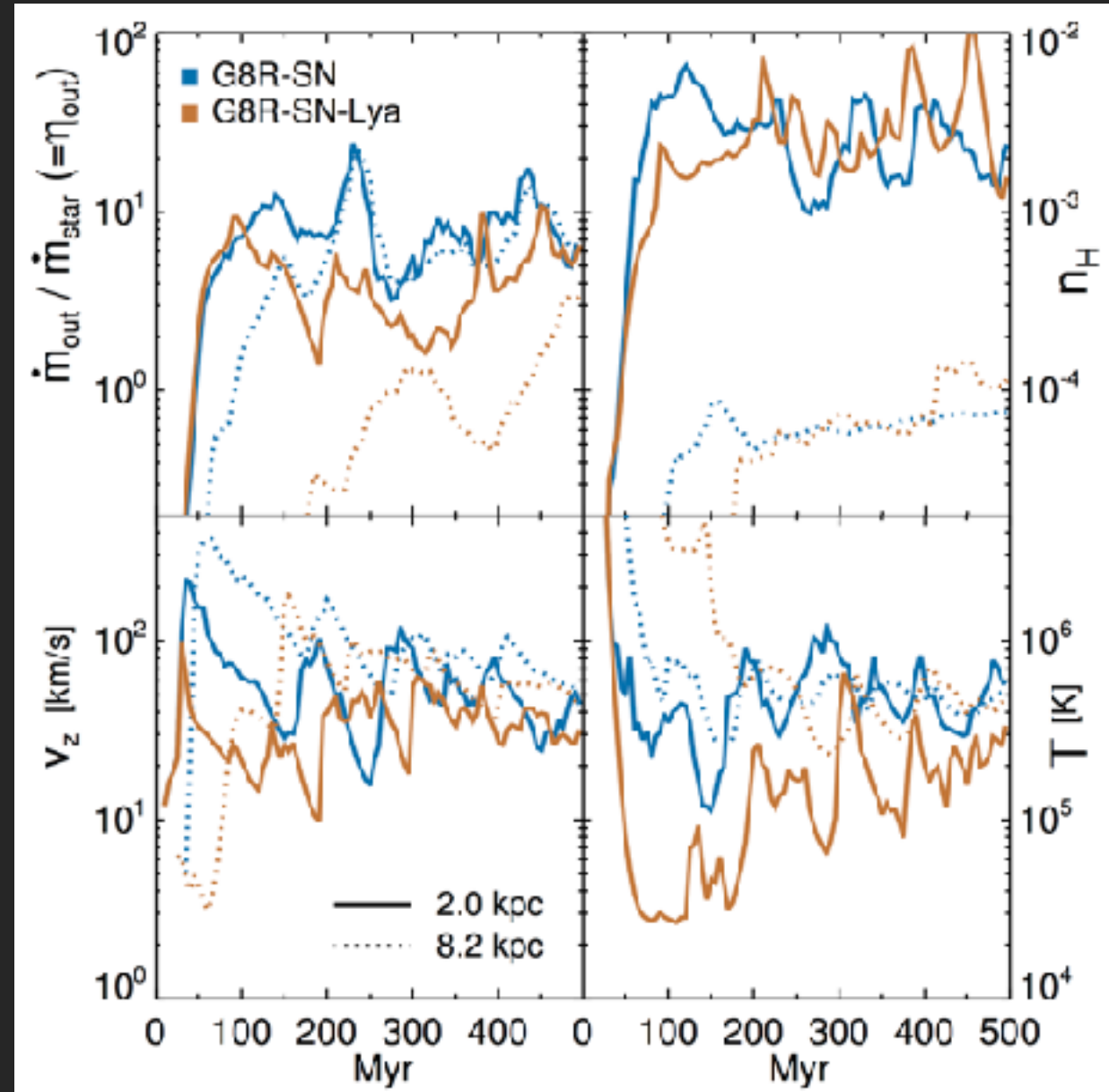
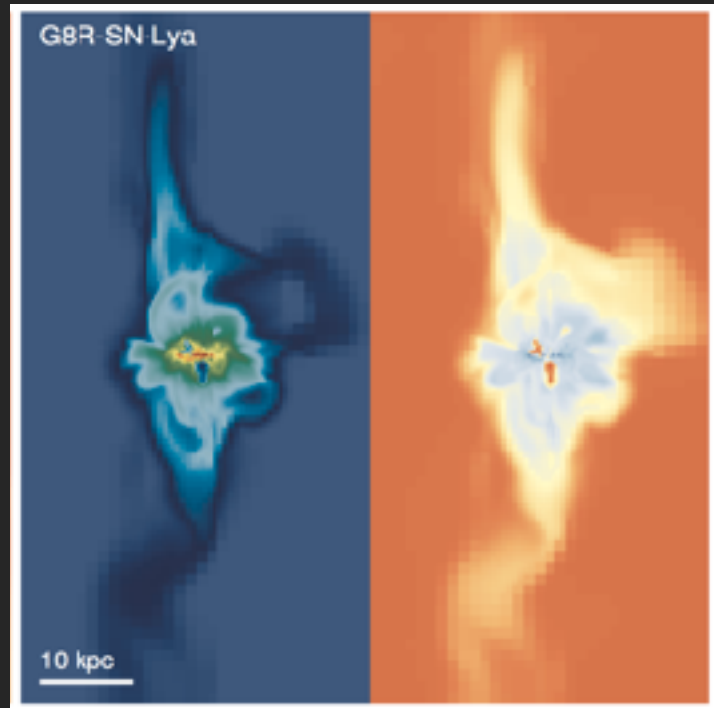
Less coherent Supernova Feedback

Weaker outflows with Lya pressure

w/o Lya



w/ Lya



WITH LYA PRESSURE

- Mass-loading factor is decreased
- Outflows become cooler and slower

Summary

- LyA photons resonantly scatter with HI, and impart 100-300 times more momentum than the single-scattering case ($L_{\text{Ly}\alpha}/c$) in the metal-poor regime
- Isolated gas-rich, metal-poor dwarf galaxy test:
 - Total stellar mass : **suppressed** by a factor of ~ 2
 - **weaker outflows** (mass loading \sim a few at $0.2 R_{\text{vir}}$)
 - **Star clusters** are more difficult to form and survive \rightarrow **important for GC formation**
 - Strong RP does not necessarily lead to stronger outflows (due to self-regulated SF)
- (Partial) Solution to the **over-cooling problem** in galaxy formation simulations