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A Novel Hybrid Scheme for Lya Line Transfer

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Collaborators

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Monte Carlo Schemes

Zheng & Miralda-Escude 2002, ApJ, 578, 33

Ahn et al, 2002, ApJ, 567, 922

Tasitsiomi, 2006, ApJ, 645, 792

Verhamme et al, 2006, A&A, 460, 397

Dijkstra et al. 2006, ApJ, 649, 14

Hansen & Oh, 2006, MNRAS, 367, 979

Semelin et al. 2007, A&A 474, 365

Laursen et al., 2009, ApJ, 696, 853

Pierleoni et al. 2009, MNRAS, 393, 872

Zheng et al. 2010, ApJ, 716, 574

Zheng et al. 2011, ApJ, 726, 38

Yajima et al. 2012, MNRAS, 424, 884

[Abe, MU, et al. 2018, MNRAS, 476, 2664](#)

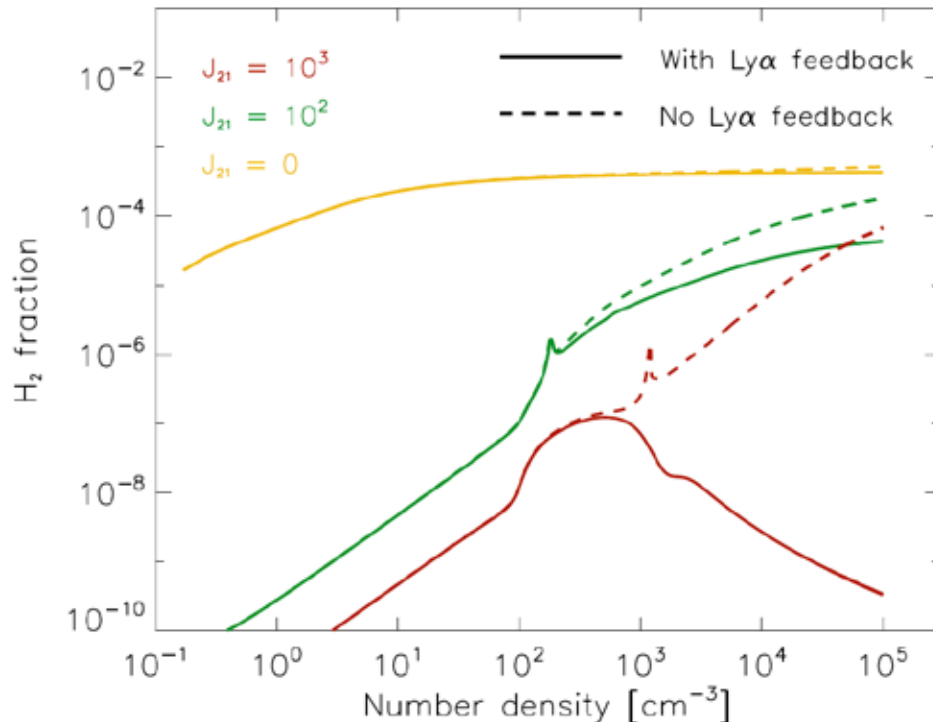
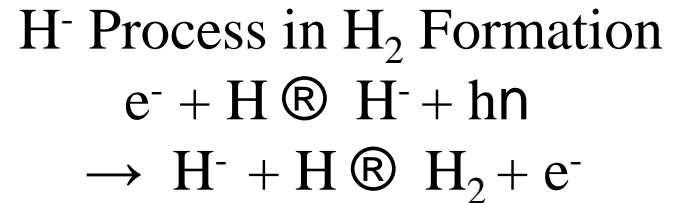
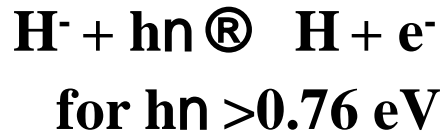
(direct SPH version: see Poster by M. Abe)

Mesh Schemes

Tasitsiomi, 2006, ApJ, 648, 762

Toward RHD with Ly α Transfer

- Cooling radiation
- Line force (eg. Dijkstra & Loeb 2008; Smith et al 2017, Kimm et al. 2018)
- H⁻ Photodetachment



Johnson & Dijkstra 2017

One-zone model

Ly α feedback

$$J_{21} = J_{\text{LW}} / 10^{-21} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ Hz}^{-1} \text{ str}^{-1}$$

Monte Carlo Schemes

- straightforward to implement
- subject to shot noises
- time consuming for all fluid elements to receive a sufficient number of photons

Mesh Schemes

- easy to couple with hydrodynamics
- transfer calculations in arbitrary optical-depth are time consuming

My talk

A novel mesh scheme coupling radiative diffusion with transfer

Hybrid Scheme

RDT: Radiative Diffusion and Transfer Scheme

∅ Context:

Consider domains containing optical thick and thin regimes.

∅ Goal:

Speed up computation by coupling diffusion and transfer equations.

Maintain accuracy of a full transfer solution.

∅ Method:

Solve the diffusion equation in optically-thick regimes.

Solve the transfer equation in optically-thin regimes.

Use diffusion solution as boundary data for the transfer equation.

radiative transfer regime
($t \gg 100$)

optically-thick
diffusion regime
($100 < t < \infty$)

Diffusion equation of resonant line scattering

$$\frac{1}{3f(x)^2} \frac{\partial^2 J(x)}{\partial t^2} = J(x) - \frac{1}{f(x)} \int R(x; x') J(x') dx'$$

$$x = \frac{n - n_0}{Dn_D} \quad f(x) : \text{line profile} \quad R(x; x') : \text{redistribution function}$$

For Lorentz profile $f(x) ; a / (\rho x^2) (t \gg 1)$

Diffusion equation
(Poisson-type)

$$\frac{\partial^2 J}{\partial t^2} + \frac{\partial^2 J}{\partial s^2} = - \frac{\sqrt{6}}{4\rho} d(t_s) d(s_s) \quad s(x) = \int^x \frac{1}{f(x)} dx ; \frac{\partial^2}{\partial s^2} = \frac{\rho}{a} \frac{\partial^2}{\partial x^2}$$

Harrington-Neufeld Solution for a Static Slab

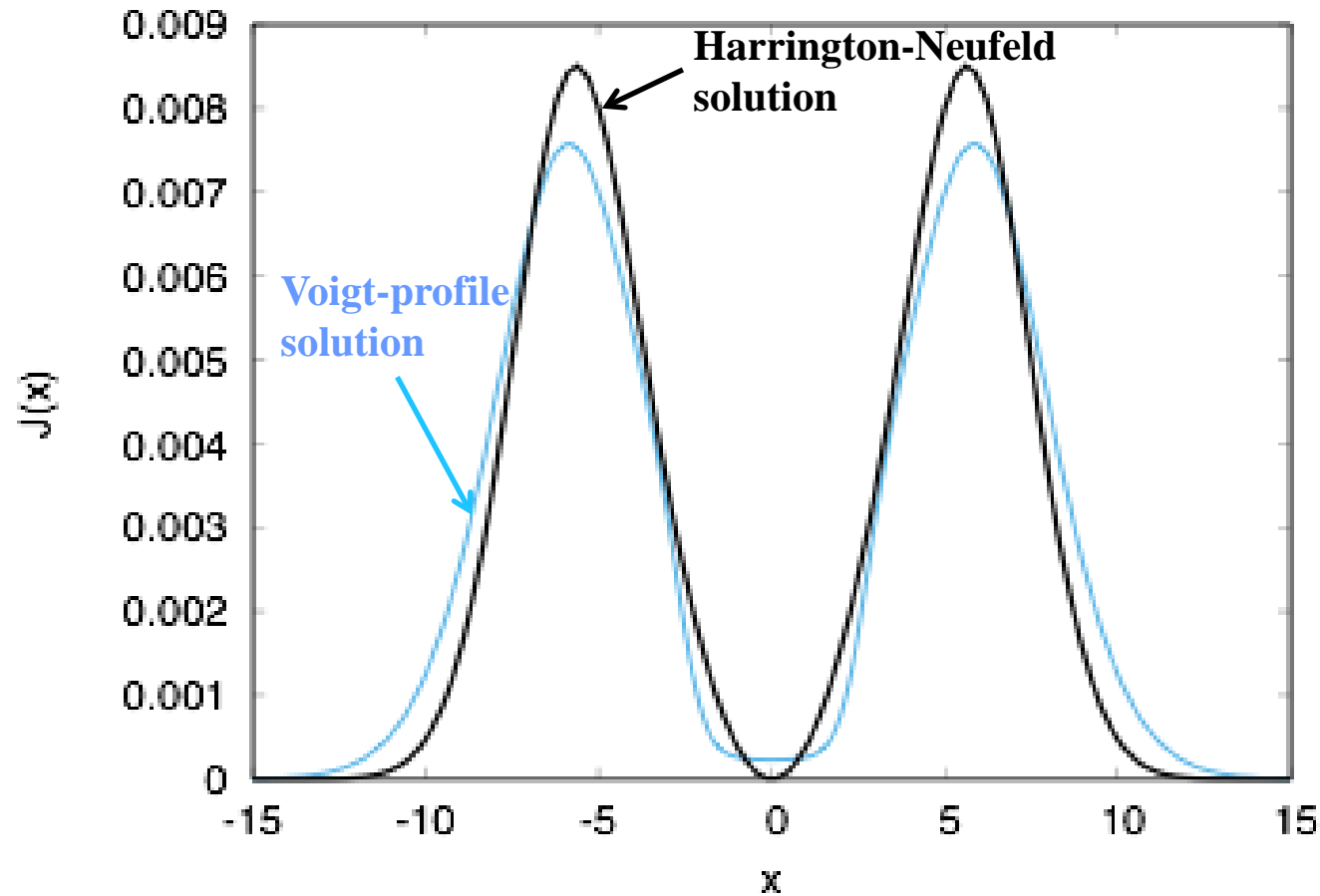
$$J(x) = \frac{\sqrt{6}}{24at_L} \frac{x^2}{\cosh \left(\sqrt{\rho^4 / 54} \left(|x^3| / at_L \right) \right)}$$

Dijkstra-Haiman-Spaans Solution for a Static Sphere

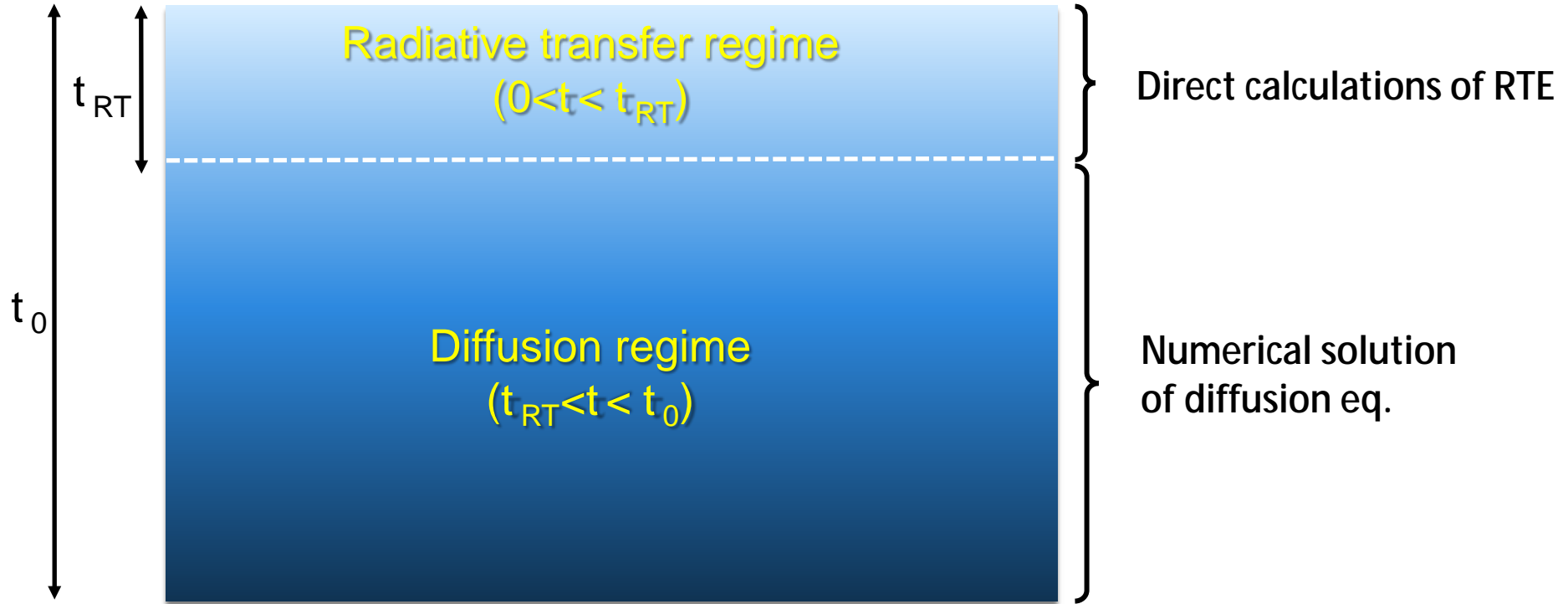
$$J(x) = \frac{\rho}{\sqrt{24at_L}} \times \frac{x^2}{1 + \cosh \left(\sqrt{2\rho^4 / 27} \left(|x^3| / at_L \right) \right)}$$

Comparison of Diffusion Solutions

$$\tau_0 = 10^4, T = 10 \text{ K}$$



Test Calculations for Ly a Transfer

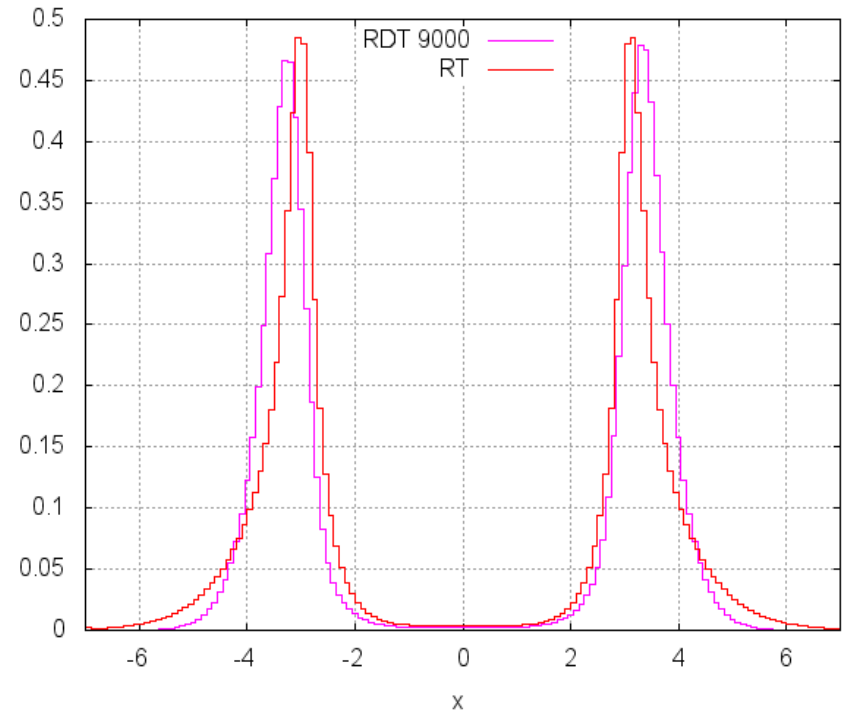
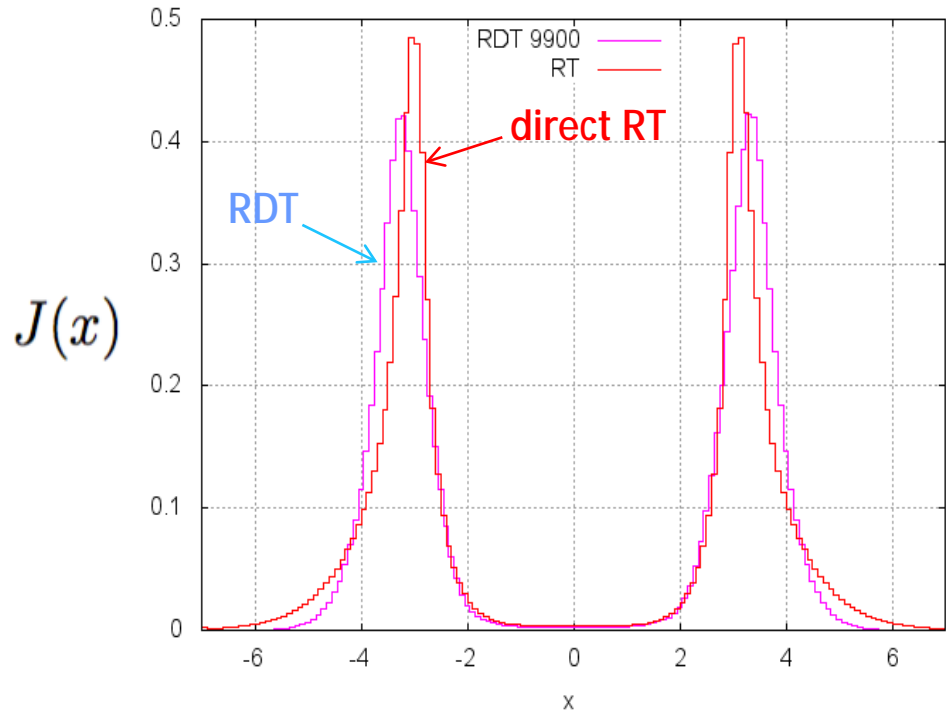


RDT vs RT

$$t_0 = 10^4, T = 10^4 \text{ K}$$

$$t_{\text{RT}} = 100 = t_0/100$$

$$t_{\text{RT}} = 1000 = t_0/10$$



- ∅ Resultant mean intensity is insensitive to t_{RT} .
- ∅ RDT with $t_{\text{RT}} = t_0/100$ give mean intensity with an accuracy of a few - 10 %.

Computational Time

$$t_0 = 10^4, Dt = 1 \text{ (} 10^4 \text{ meshes)} \quad t_{\text{diff}} = t_0 - t_{\text{RT}}$$

Method	t_{RT}	t_{diff}	Iteration #	Computational time by one core	Acceleration
RT(Direct)	10000	0	262,920	44.4days	1
RDT	1000	9000	12,664	15hrs	70
RDT	100	9900	556	8min	8000

Summary

- ∅ We've developed a novel hybrid scheme, RDT (Radiative Diffusion and Transfer), with the diffusion solution based on the Voigt profile and the exact redistribution function for non-coherent resonant line scattering in arbitrary optical-depth media on meshes.
- ∅ RDT calculations with $t_{RT} = t_0/100$ give the mean intensity with an accuracy of a few - 10 %. The accuracy of radiation force is enhanced with increasing t_0 .
- ∅ RDT scheme can reduce the computational cost dramatically and allow us to properly calculate the formation of Pop III objects or LAEs incorporating Ly α feedback.