

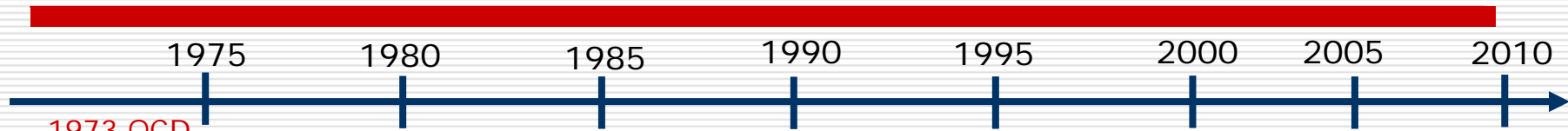


Akira Ukawa
University of Tsukuba

Present and Future of Lattice QCD

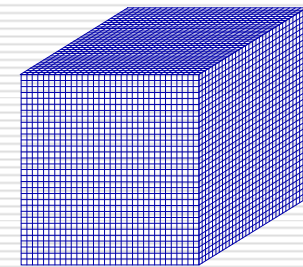
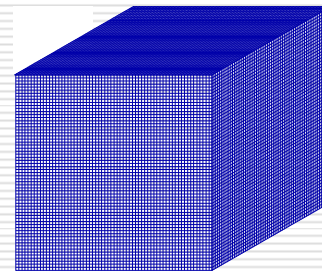
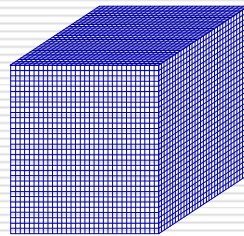
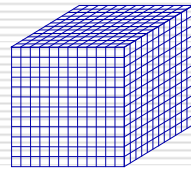
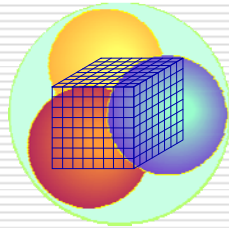


Lattice QCD over the years...



Physics

1st spec calculation
1981
Hamber-Parisi
Weingarten



Lattice size L

0.8fm
 $8^3 \times 16$

1.6fm
 $16^3 \times 32$

2.4fm
 $24^3 \times 48$

3.0fm
 $64^3 \times 118$

3.0fm
 $32^3 \times 64$

N_f=0 quenched

Algorithms

N_f = #sea quarks

N_f=2 u,d

N_f=2+1 u,d,s

Machines

1st generation 1Gflops
2nd generation 10Gflops
3rd generation 1Tflops
4th generation 10Tflops 100Tflops



APE1

QCDPAX



APE100



QCDSP



QCDDOC



BlueGene/L,P



CP-PACS



PACS-CS



My personal feeling

- Lattice QCD is finally turning a corner in the last couple of years.
 - Previously, despite the premise, it remained an approximate method requiring extrapolations in a number of ways (quenching, unphysically heavy quark masses, etc).
 - Progress over the years has been removing these restrictions, and it is now becoming a *real first principle method*, not only in principle but also *in practice*, for actually calculating physical quantities *at the physical point on physically large lattices, i.e., Nature*.
-



What I wish to do today

Review recent progress and try to share this feeling with you

- Algorithmic progress
- Flavor physics
- High temperature/density QCD
- No more nuclear physics?
- Computer trends -if time allows-
- Conclusions



Algorithmic progress



Quantum Chromodynamics

Gross-Wilczek-Politzer 1973

- Quantum field theory of quarks and gluons

$$\left. \begin{array}{ll} q_f(x) & \text{Quark field} \\ A_\mu(x) & \text{Gluon field} \end{array} \right\} \text{ defined over 4-dim space time}$$

$$L_{QCD} = \frac{1}{8\pi\alpha_s} \text{Tr}(F_{\mu\nu}F_{\mu\nu}) + \sum_f \bar{q}_f (\gamma_\mu \cdot (\partial_\mu - iA_\mu) + m_f) q_f \quad \text{QCD lagrangian}$$

$$\langle O(A, \bar{\psi}, \psi) \rangle = \frac{1}{Z} \int dA d\bar{q} dq O(A, \bar{q}, q) e^{-\int d^4x L_{QCD}} \quad \text{Physical quantities by Feynman path integral}$$

- Knowing

1 coupling constant
and
6 quark masses

$$\alpha_s = \frac{g_s^2}{4\pi}$$

$$m_u, m_d, m_s, m_c, m_b, m_t$$

will allow full understanding of hadrons and their strong interactions

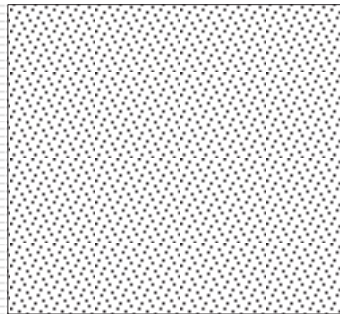
"fulfilling Yukawa's dream of 1934 in a refined way"



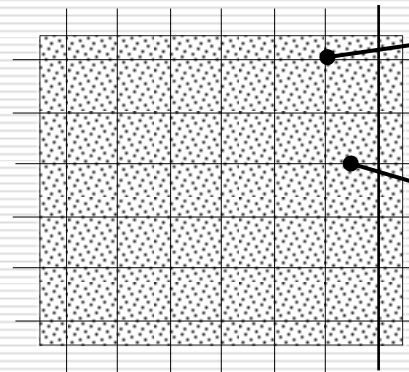
QCD on a space-time lattice

K. G. Wilson 1974

Space-time continuum



Space-time lattice



q_n

quark fields on
lattice sites

$U_{n\mu}$

gluon fields on
lattice links

□ Feynman path integral

■ Action $S_{QCD} = \frac{1}{g_s^2} \sum_P \text{tr}(UUUU) + \sum_f \bar{q}_f (\gamma \cdot U + m_f) q_f$

■ Physical quantities as integral averages



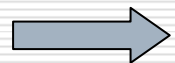
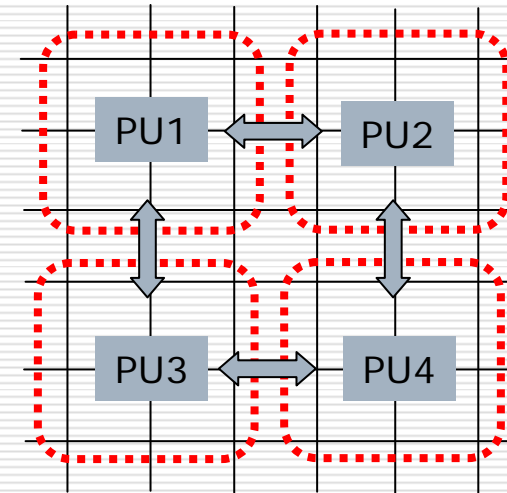
Monte Carlo calculation
of the integral average

$$\langle O(U, \bar{q}, q) \rangle = \frac{1}{Z} \int \prod_{n\mu} dU_{n\mu} \prod_n d\bar{q}_n dq_n O(U, (\bar{q}, q)) e^{-S_{QCD}}$$



Lattice QCD as computation(I)

- QCD is a local field theory; only nearest neighbor interactions
- Natural mapping of space-time lattice to processor array
 - Each compute node carries a sub-lattice
 - Only nearest neighbor communication needed



Highly parallelizable and scalable



Lattice QCD as computation(II)

- Quarks are fermions, so their field, being anti-commuting, needs a special trick

$$\int \prod_n d\bar{q}_n dq_n e^{-\sum_{n,m} \bar{q}_n D_{nm}(U) q_m} = \det D(U) = \int \prod_n d\bar{\phi}_n d\phi_n e^{-\sum_{n,m} \bar{\phi}_n \left(\frac{1}{D(U)} \right)_{nm} \phi_m}$$

- Need to invert the lattice Dirac operator $D(U)$
 - Sparse but large matrix
 - Large condition number $\sim 1/m_q$ for quarks in nature

$$\sum_m D_{nm}(U) x_m = \phi_n \Rightarrow x_n = \left(\frac{1}{D(U)} \right)_{nm} \phi_m \quad \text{Core calculation of QCD}$$



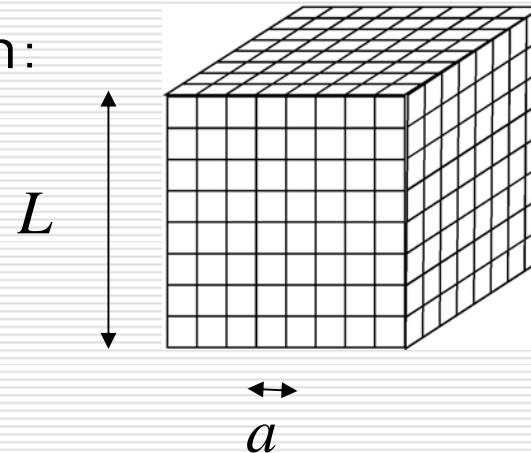
Computationally very intensive



Difficulties with light quark masses

- Parameters of lattice QCD simulation:

- Quark mass m_q or $m_\pi \propto \sqrt{m_q}$
- Lattice size L (fm)
- Lattice spacing a (fm)



- #arithmetic ops of hybrid Monte Carlo (HMC) algorithm for $N_f=2$ flavor full QCD(2001)

$$\#FLOP's \approx 1.9 \left[\frac{\#conf}{1000} \right] \cdot \left[\frac{m_\pi}{500MeV} \right]^{-6} \cdot \left[\frac{L}{3fm} \right]^5 \cdot \left[\frac{a}{0.1fm} \right]^{-7} Tflops \cdot year$$

- Severe scaling toward small pion mass /large volume/small lattice spacing



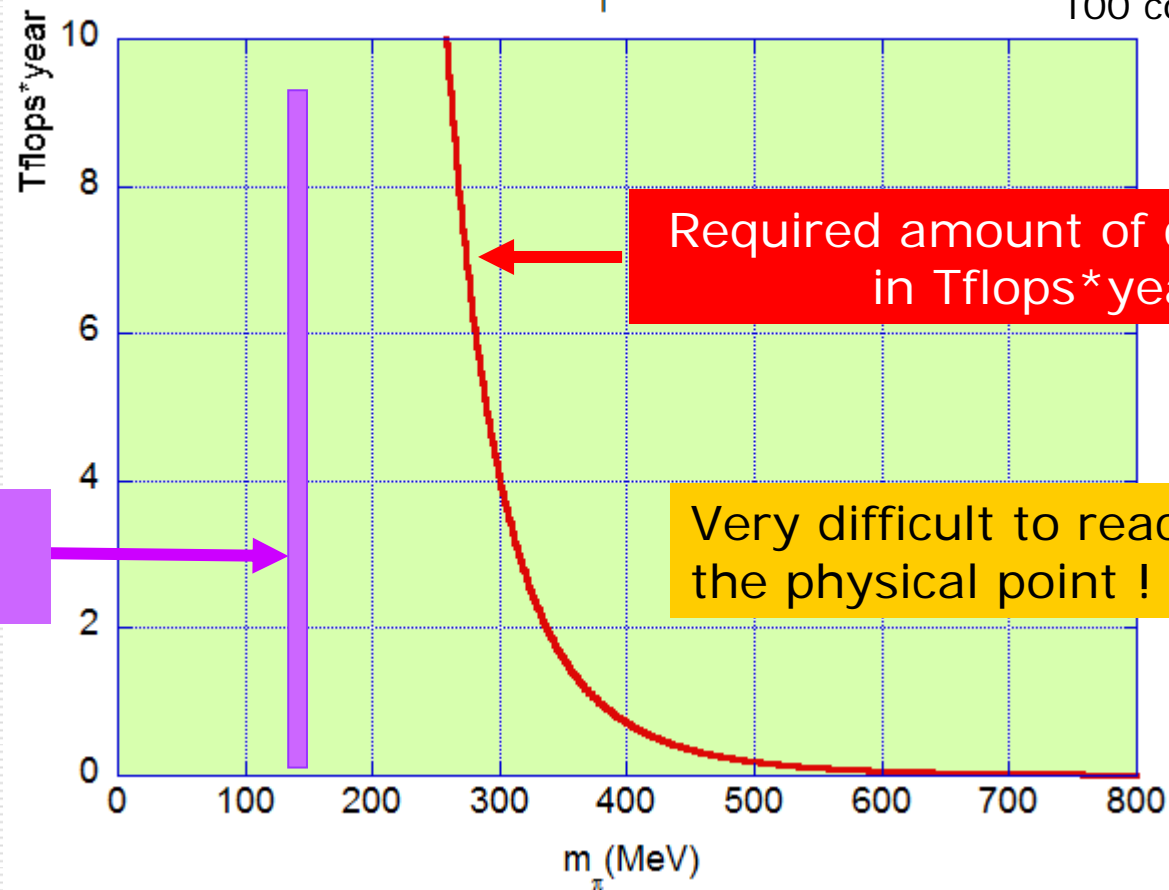
“Berlin wall” at Lattice 2001@Berlin

A. Ukawa for CP-PACS and JLQCD

L=3fm QCD with $N_f=2+1$ dynamical quarks

$a=0.1\text{fm}$

100 configurations

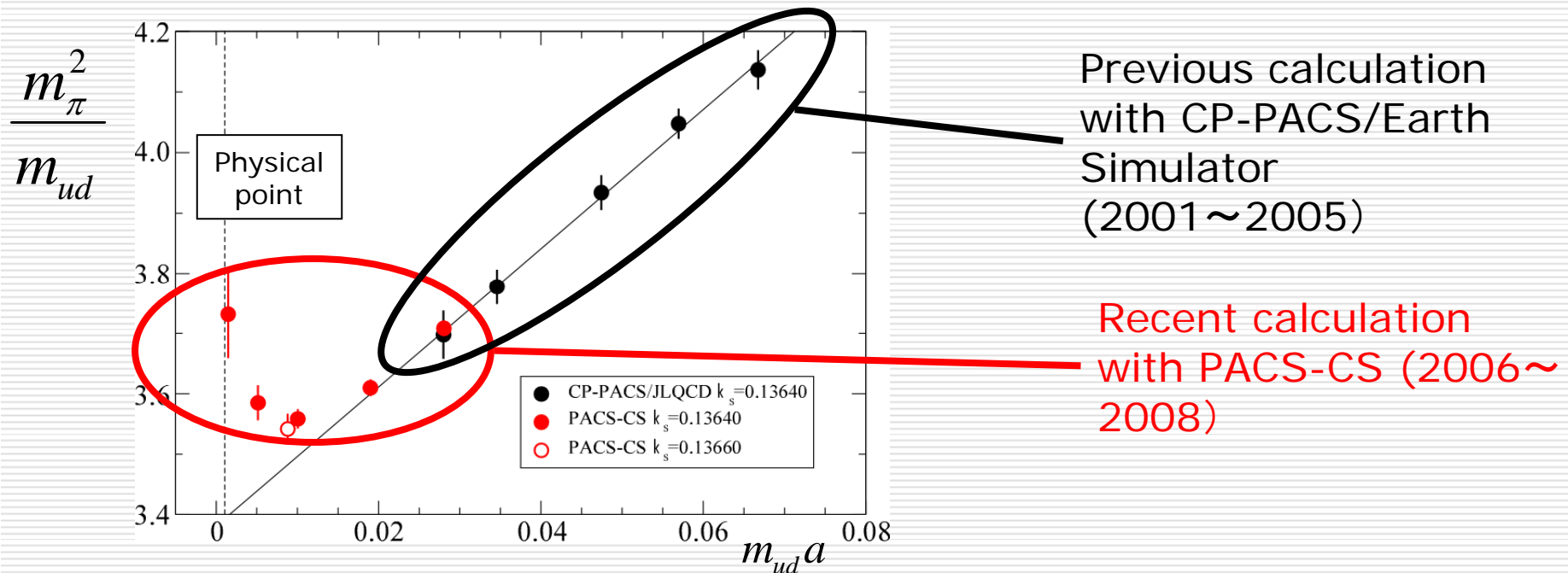


Physical point
i.e., $m_\pi = 135\text{MeV}$

Very difficult to reach
the physical point !



Why so important to go physical?



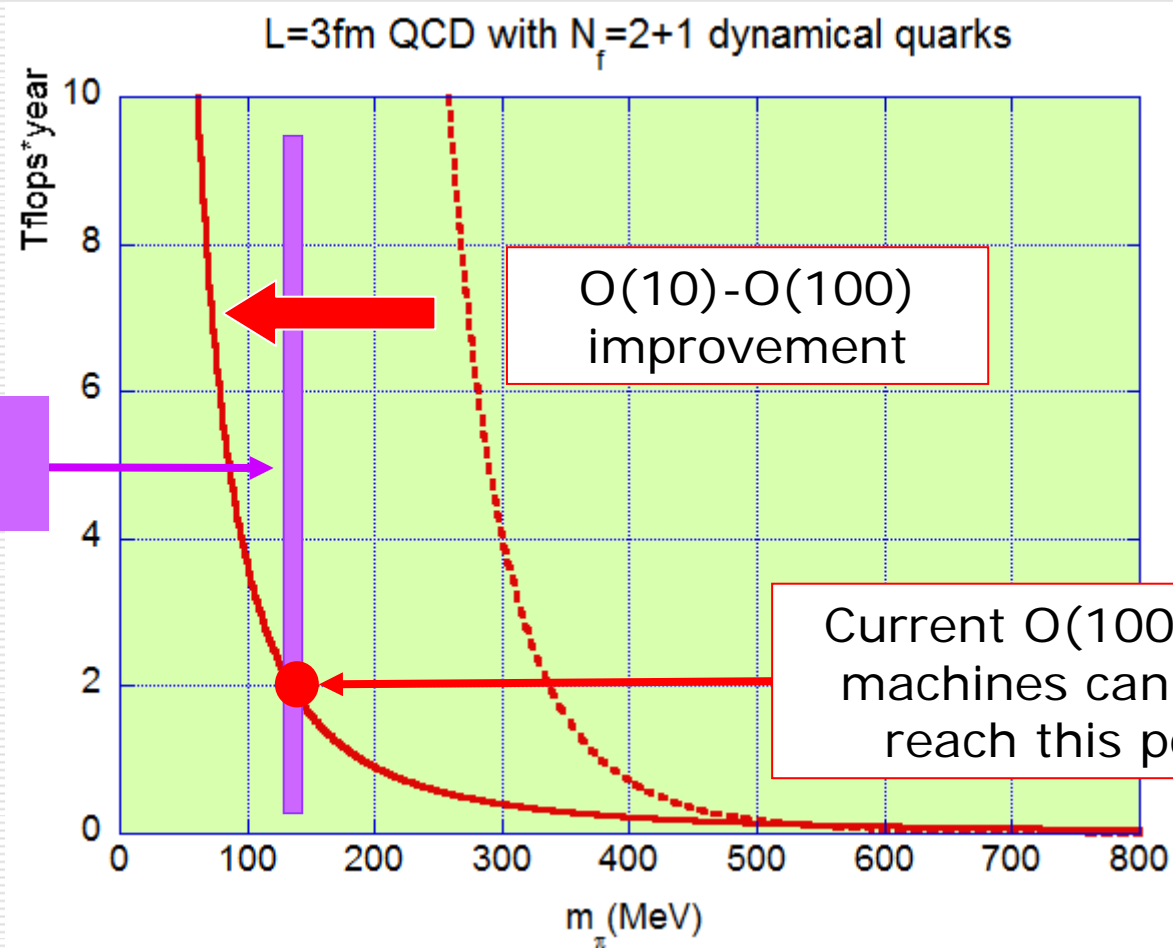
- Anticipated effect of *chiral logarithm* at zero quark mass

$$\frac{m_\pi^2}{m_{ud}} \propto 1 + \frac{2Bm_{ud}}{(4\pi f)^2} \log \frac{2Bm_{ud}}{\mu^2} + \dots$$

- However, extrapolation difficult to control since
 - Convergence radius a priori not known
 - Have to determine a number of unknown constants



Improved HMC algorithm



Physical point
i.e., $m_\pi=135\text{MeV}$

Physical point simulation has become reality



How that progress came about?

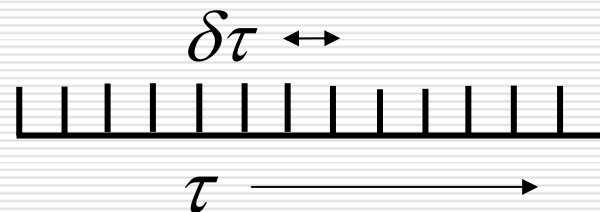
- Molecular dynamics equation of hybrid Monte Carlo algorithm

$$\frac{d}{d\tau} U_{n\mu} = -iU_{n\mu} P_{n\mu}$$

$$\frac{d}{d\tau} P_{n\mu} = F_{n\mu} = \frac{1}{g^2} \overset{\text{gluon force}}{(UUUU)_{n\mu}} + \overset{\text{quark force}}{\bar{\phi} \left(\frac{1}{D(U)} \right) \frac{\partial D(U)}{\partial U_{n\mu}} \left(\frac{1}{D(U)} \right) \phi}$$

Most time-consuming part of computation

- Molecular dynamics equation is integrated in discrete steps, so a larger time step is better!





Key observation

M. Luescher (2005)

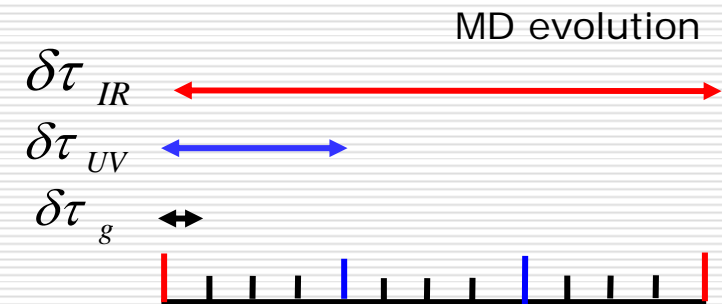
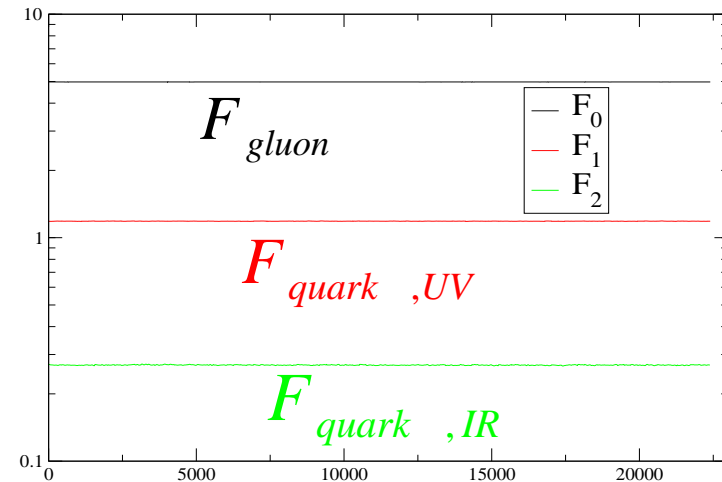
- Strategy :
 - Separate UV and IR modes of quark fluctuations
 - Use separate time step to UV and IR modes
- Numerically found:

$$F_{gluon} \gg F_{quark,UV} \gg F_{quark,IR}$$

so take

$$\delta\tau_{gluon} \ll \delta\tau_{quark,UV} \ll \delta\tau_{quark,IR}$$

i.e., one can enlarge the time step for the most compute intensive IR quark force, leading to large acceleration of the algorithm.



This is physics.



Recent large-scale $N_f=2+1$ calculations

□ Features

- Fully incorporates dynamical effects of up, down, strange sea quarks, hence called “ $N_f=2+1$ ”
- Pion mass reaching down to even attempting the physical point
- Lattice size to avoid finite size effects

$$m_\pi \approx 200 - 300 \text{ MeV}$$

$$m_\pi \approx 140 \text{ MeV}$$

$$m_\pi L \approx 3 - 4$$

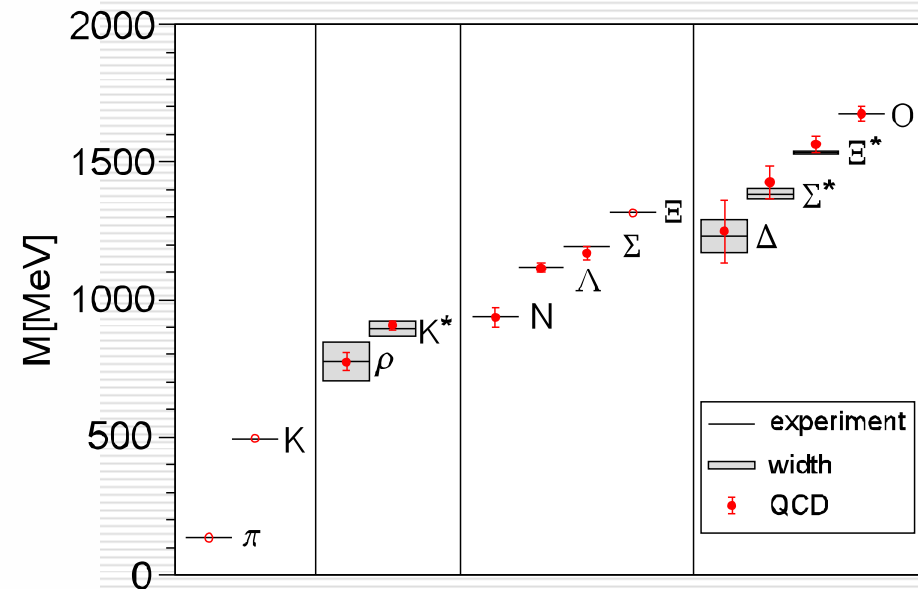
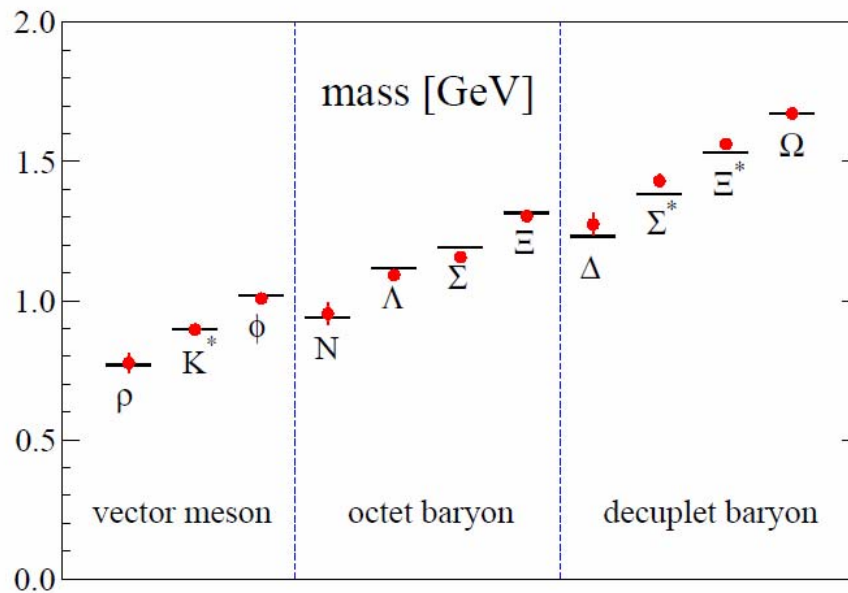
□ Collaborations	action	a (fm)	L (fm)	m_π (MeV)
■ MILC	staggered	0.06	4.0	180
■ PACS-CS	wilson-clover	0.09	2.3	155
■ BMW	wilson-clover	0.09	4.0	190
■ RBC-UKQCD	domain-wall	0.09	4.0	290
■ JLQCD	overlap	0.11	2.8	320
■ ETMC ($N_f=2$)	twisted mass	0.07	3.0	250



Hadron spectrum 2008

PACS-CS Collaboration
(Tsukuba, Japan)
Phys. Rev. D79 034504 (2008)

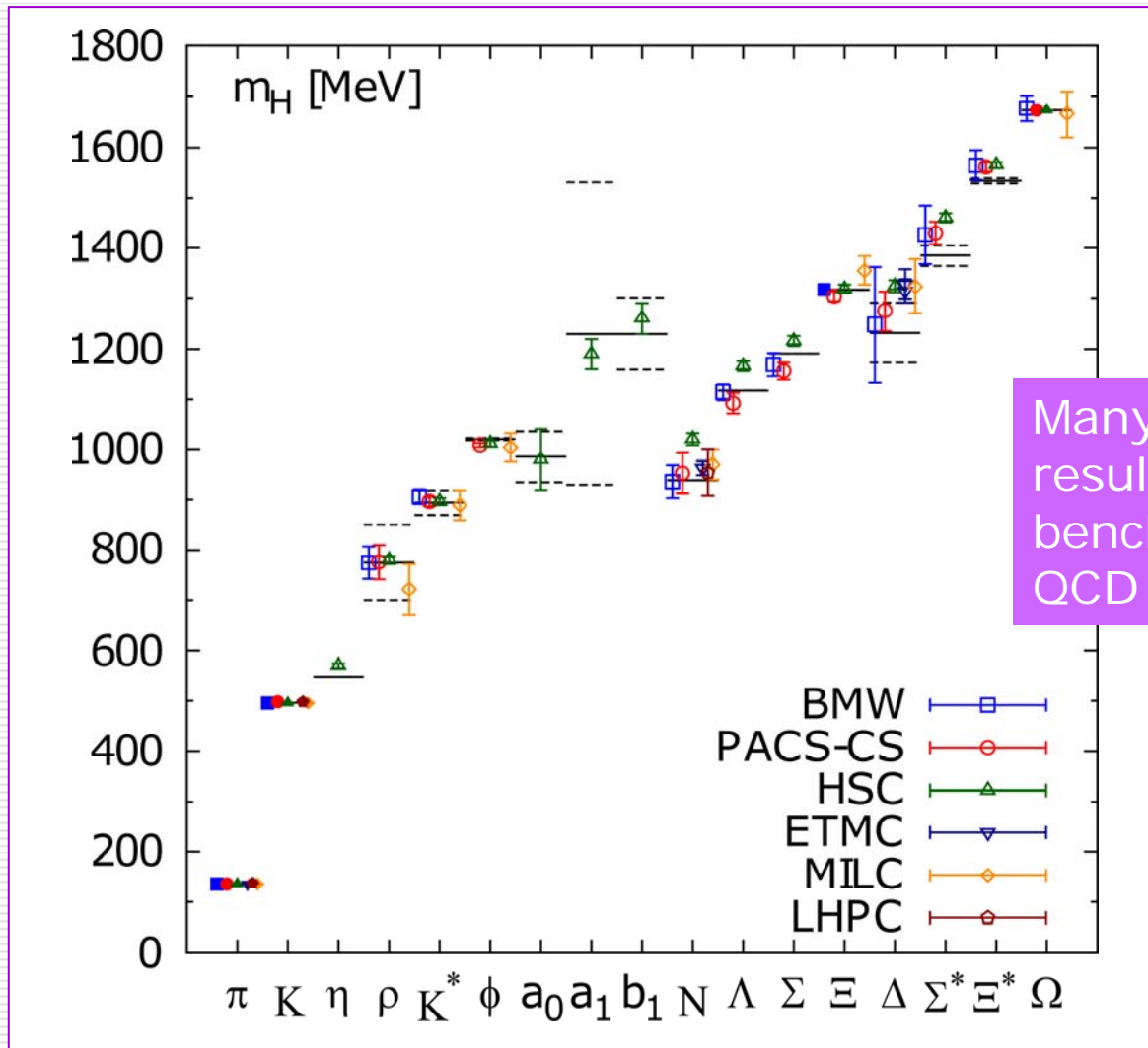
BMW Collaboration
(Butapest-Marseille-Wuppertal)
Science 322(2008) 1224
Continuum extrapolated





Hadron spectrum 2009

From E. Scholz@Lattice 2009



Many more $N_f=2+1$ flavor results; becoming a basic benchmark of any lattice QCD calculations

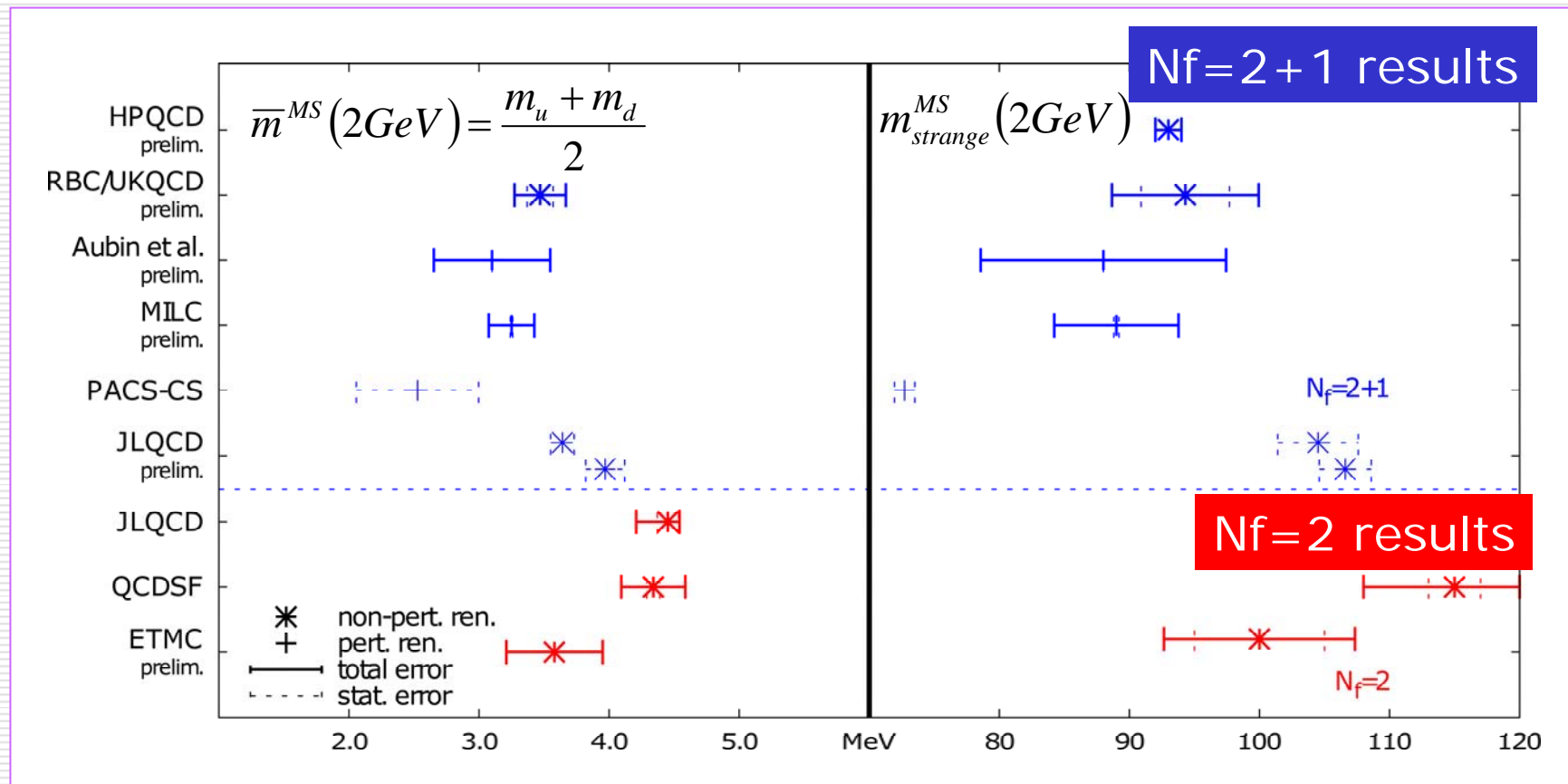


Light quark masses – status –

From E. Scholz@Lattice 2009

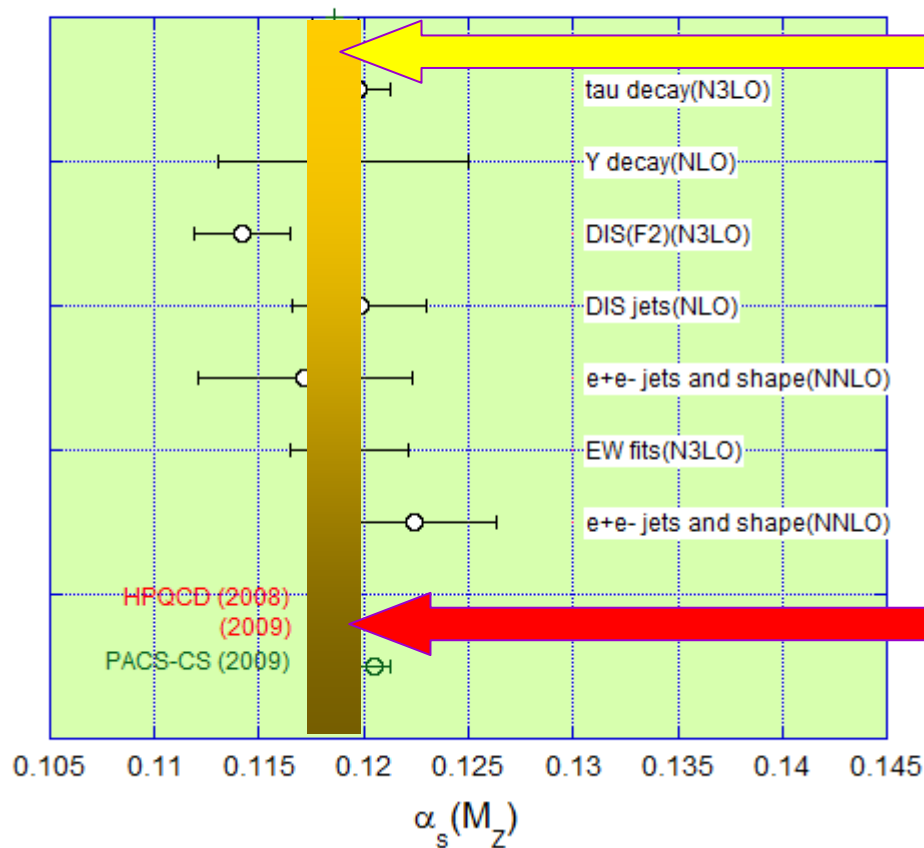
- $N_f=2+1$ continuum estimation indicates

$$\bar{m}^{MS}(2GeV) \approx 3MeV, \quad m_s^{MS}(2GeV) \approx 90MeV$$





Strong coupling constant – status -



- Experimental average

$$\alpha_s^{MS}(M_Z) = 0.1186 \pm 0.0011$$

S. Bethke, ArXiv.0908.1135

- Nf=2+1 Lattice QCD

$$\alpha_s^{MS}(M_Z) = 0.1184 \pm 0.0004$$

HPQCD 2009 update

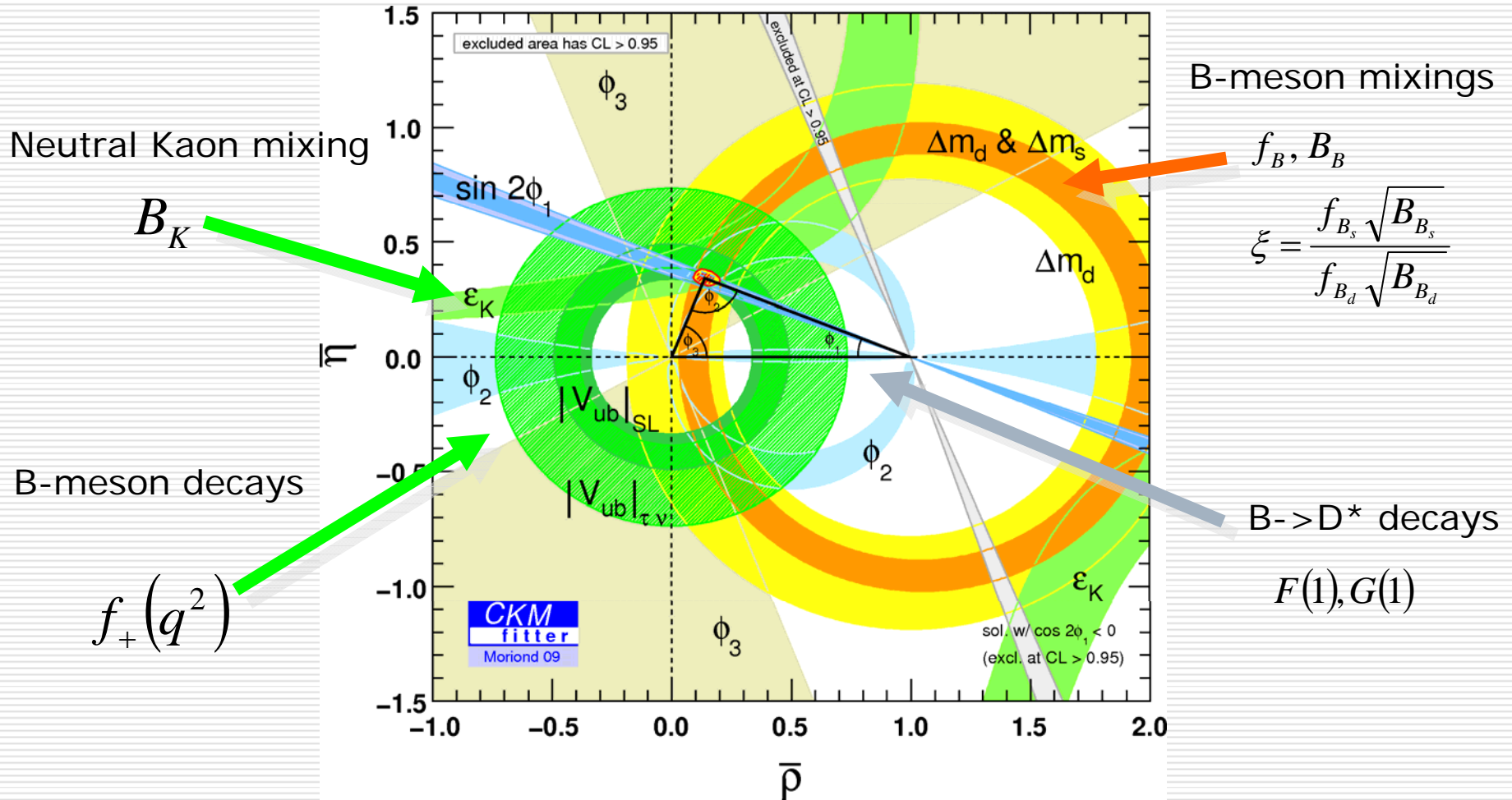
(private communication from C. Davies)



Flavor physics



CKM matrix and lattice QCD





Inputs from lattice QCD 2009

Van de Water@Lattice 2009

<i>Quantity</i>	<i>Value</i>	<i>Error</i>
\hat{B}_K	0.725 ± 0.028	4%
ξ	1.243 ± 0.028	2%
$ V_{ub} _{excl.}$	$3.42 \pm 0.37 \times 10^{-3}$	11%
$ V_{cb} _{excl.}$	$38.6 \pm 1.2 \times 10^{-3}$	3%
f_K	$155.8 \pm 1.7 MeV$	1%

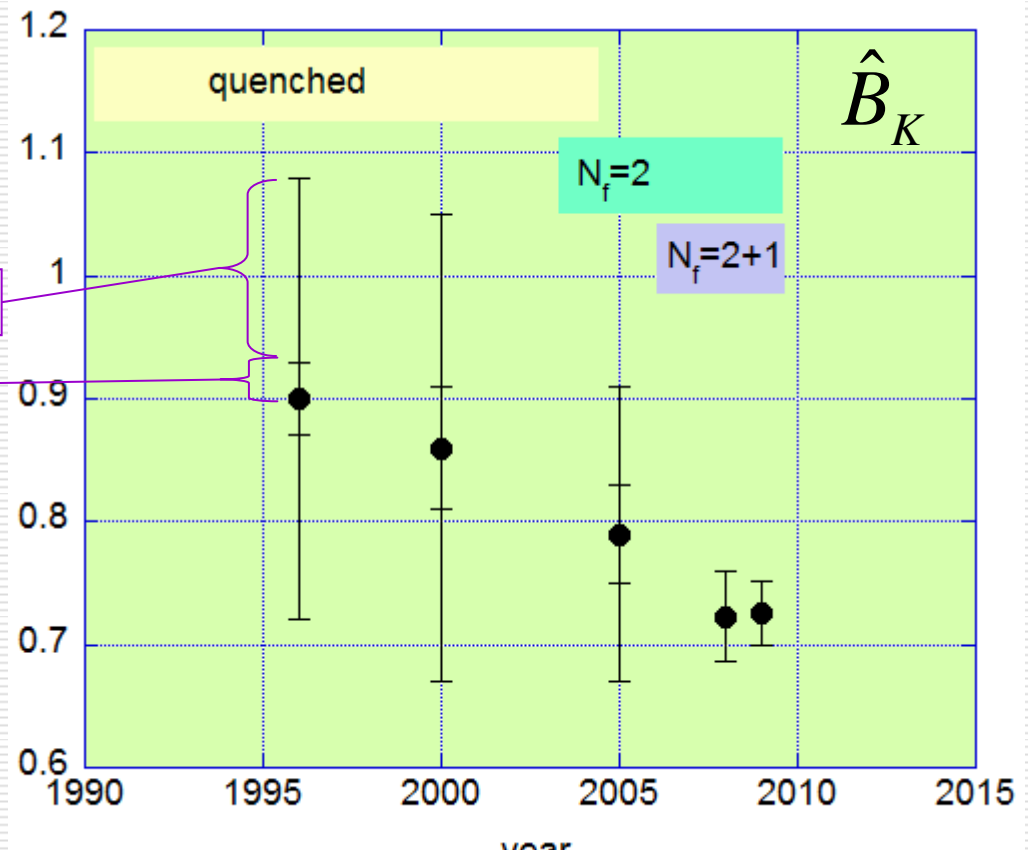
All values

- From $N_f=2+1$ simulations
- All errors calculated/estimated
Statistical/Chiral extrapolation/Finite volume/Continuum extrapolation



B_K over the years

- 1996: JLQCD
 - Quenched
 - Continuum extrapolated
- 2008: RBC/UKQCD
 - ArXiv 0710.5136
 - $N_f=2+1$
 - Chiral action
 - DWF on DWF sea
 - one lattice spacing
- 2009: Aubin-Laiho-Van de Water
 - ArXiv 0905.3947
 - $N_f=2+1$
 - Chiral action
 - Overlap on staggered sea
 - Two lattice spacing and Continuum extrapolated

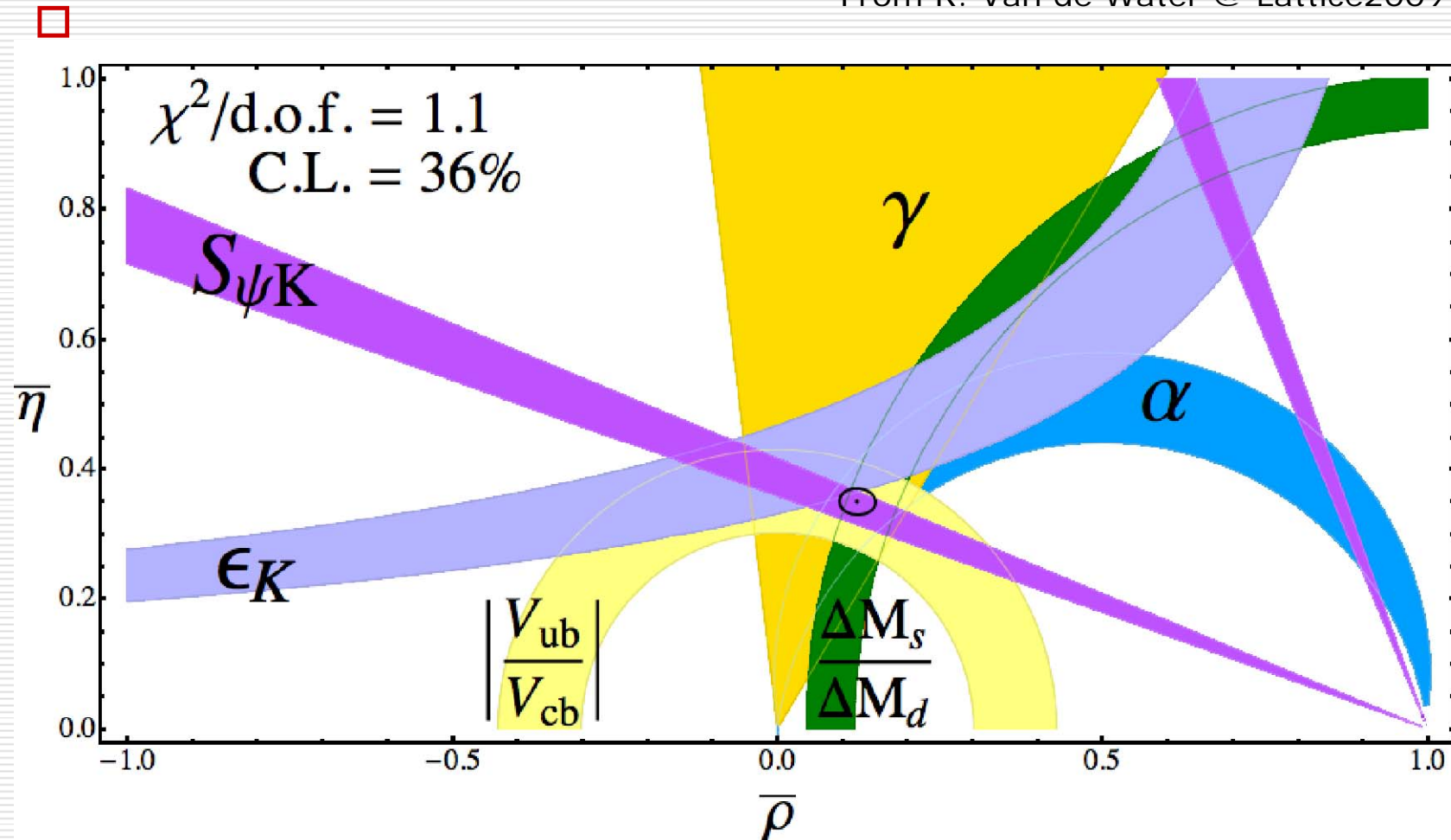


4% error of B_K is now smaller than 10% error due to $|V_{cb}|^4$ in ε_K



CKM status with lattice inputs 2009

From R. Van de Water @ Lattice2009





CKM unitarity with lattice inputs 2009

- *First row unitarity* holds to 0.1% accuracy

V_{ud}	V_{us}	V_{ub}	$\sum V_{ij} ^2 - 1$
0.97425	0.2246	0.00342	-0.0004
± 0.00022	± 0.0012	± 0.00037	± 0.0013
Nuclear transitions Hardy-Towner ArXiv 0812.12.02	K \rightarrow pi FlaviA Net + Nf=2+1 lattice QCD RBC/UKQCD Lattice'09	B \rightarrow pi HFAG + Nf=2+1 lattice QCD FNAL/MILC, HPQCD Lattice'08	

- *Second row unitarity* requires much improvement

V_{cd}	V_{cs}	V_{cb}	$\sum V_{ij} ^2 - 1$
0.239	0.969	0.039	-0.002
± 0.032	± 0.105	± 0.001	± 0.110
2004 number from FNAL/MILC/HPQCD 2005 number no better $V_{cs}=1.015 \pm 0.107$		B \rightarrow D, D* HFAV + Nf=2+1 lattice QCD FNAL/MILC, Lattice'08	

Charm physics on the lattice has to improve!

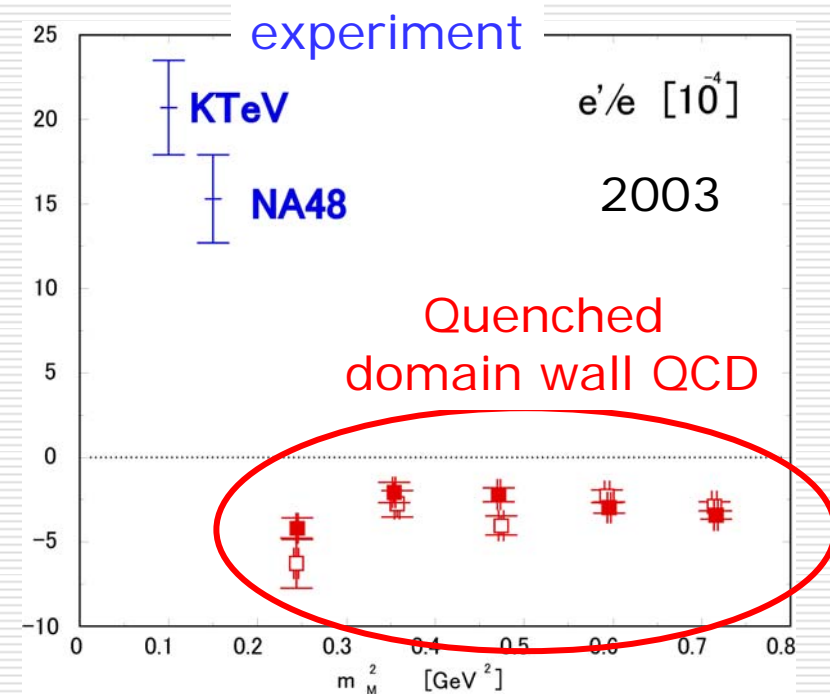


Comment on $\text{Re}(\varepsilon' / \varepsilon)$

- Failure of the previous lattice calculation (2003) indicates
 - Inadequacies of Quenched approximation
 - Failure of SU(3) chiral perturbation theory

- Steady progress since then

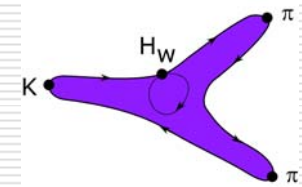
$$\frac{\varepsilon'}{\varepsilon} = \frac{\omega}{\sqrt{2}|\varepsilon|} \left[\frac{\text{Im} A_2}{\text{Re} A_2} - \frac{\text{Im} A_0}{\text{Re} A_0} \right]$$





Finite volume framework for $K \rightarrow \pi \pi$ amplitude

C. Lellouche and M. Luescher (2001)



- Finite-size formula for direct $K \rightarrow \pi \pi$ amplitude

$$\underbrace{\left| A_{\text{physical}}(K \rightarrow \pi\pi) \right|^2}_{\text{Physical amplitude}} = 8\pi \left(\frac{E_{\pi\pi}}{p} \right)^3 \underbrace{\left\{ p \frac{\partial \delta(p)}{\partial p} + q \frac{\partial \phi(q)}{\partial q} \right\} \left| \langle K | H_W | \pi\pi \rangle_{\text{lattice}} \right|^2}_{\text{Finite volume lattice amplitude}}$$

Physical amplitude

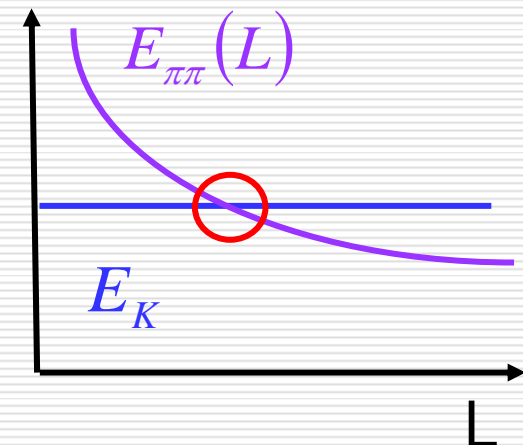
Finite volume lattice amplitude

$$p^2 = E_{\pi\pi}^2 / 4 - m_\pi^2, \quad q^2 = (pL / 2\pi)^2$$

$$\tan \phi(q) = -\frac{q\pi^{3/2}}{Z_{00}(1; q^2)}, \quad Z_{00}(1; q^2) = \frac{1}{\sqrt{4\pi}} \sum_{\vec{n} \in \mathbb{Z}^3} \frac{1}{\vec{n}^2 - q^2}$$

$$\delta(p) = n\pi - \phi(q) \quad \text{Phase shift}$$

$$\text{Requires } E_K = E_{\pi\pi}(L)$$

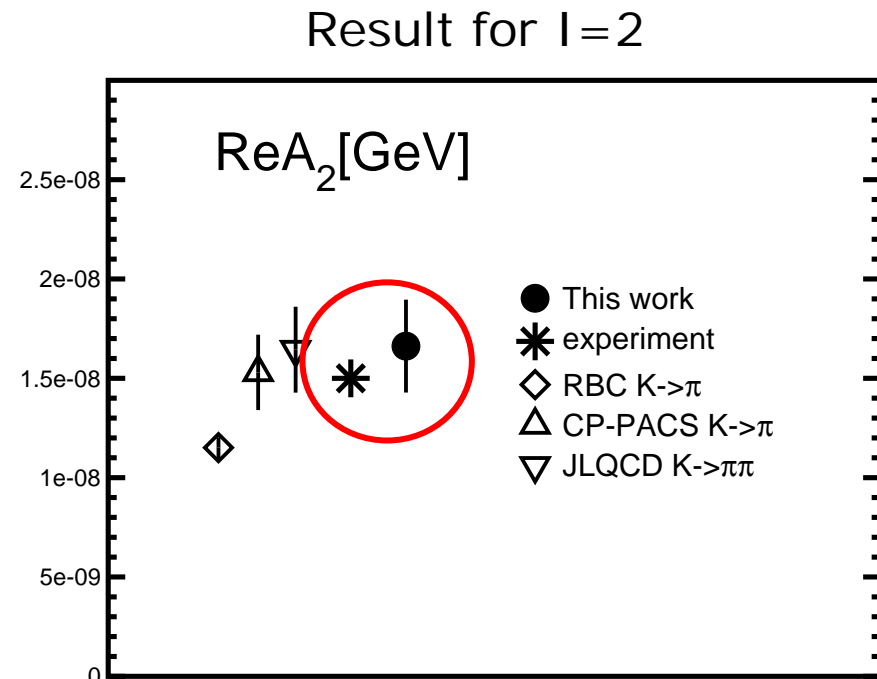
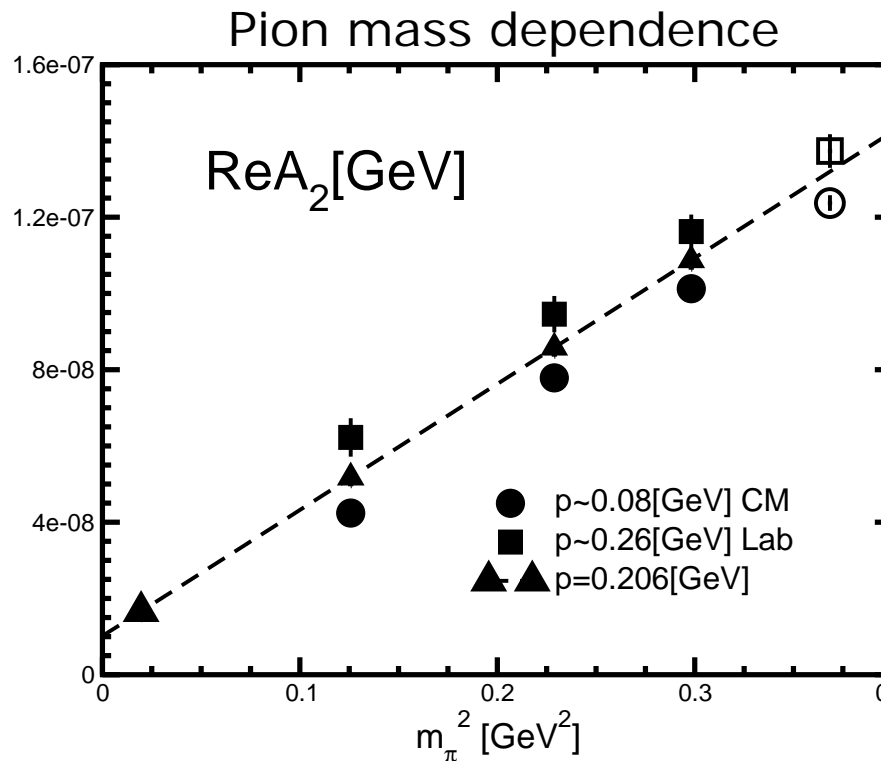




Application for $I=2$ channel with domain wall QCD

T. Yamazaki et al, Archive 0807.3130 (2008)

- Only $I=2$ channel at present, for which previous attempts yielded reasonable results
- But an encouraging start toward a direct $K \rightarrow \pi \pi$ calculation in $N_f=2+1$ full QCD; expect such a calculation in a few year's

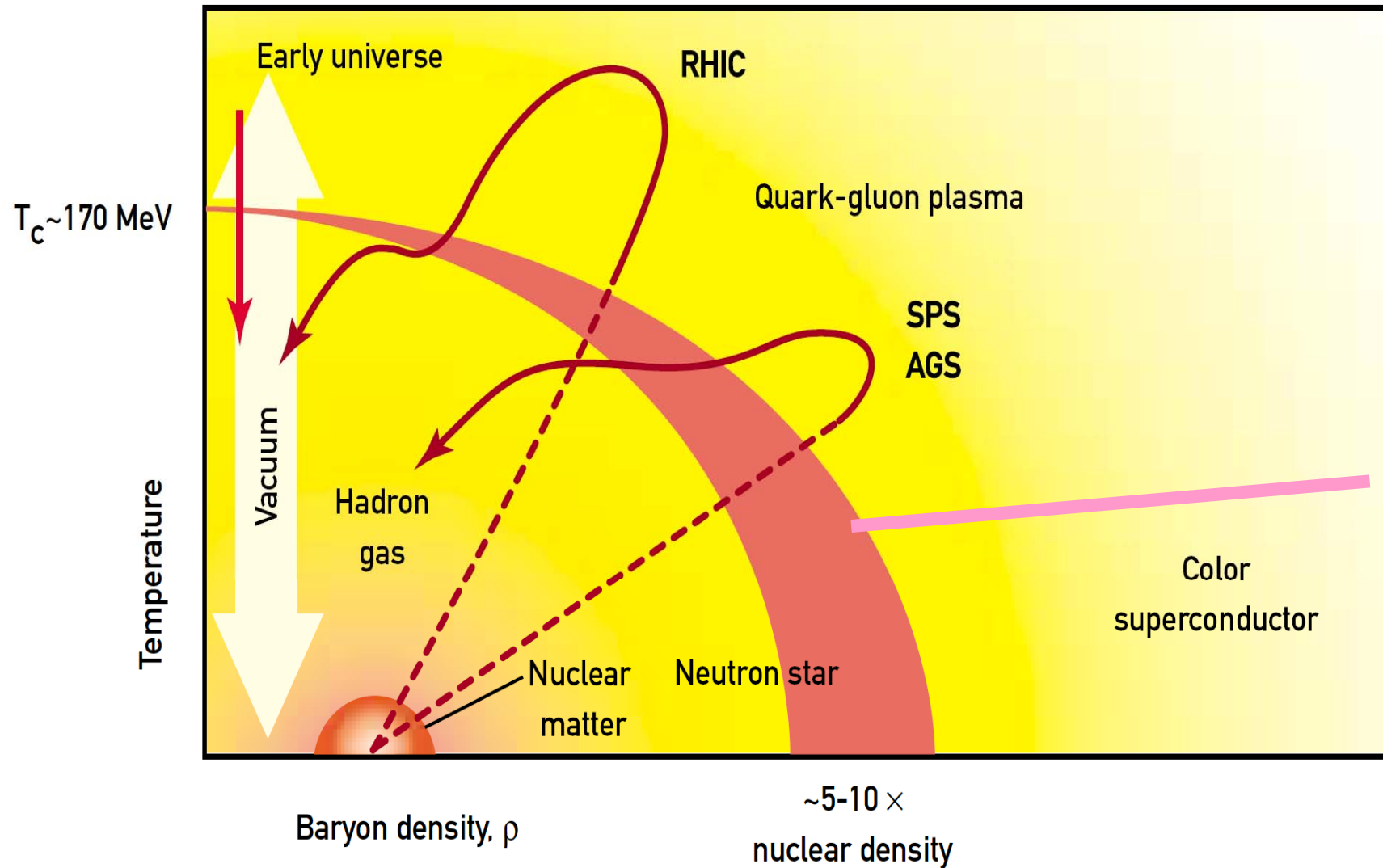




Hot/Dense QCD



Schematic QCD phase diagram





$N_f=2+1$ Phase diagram at zero density

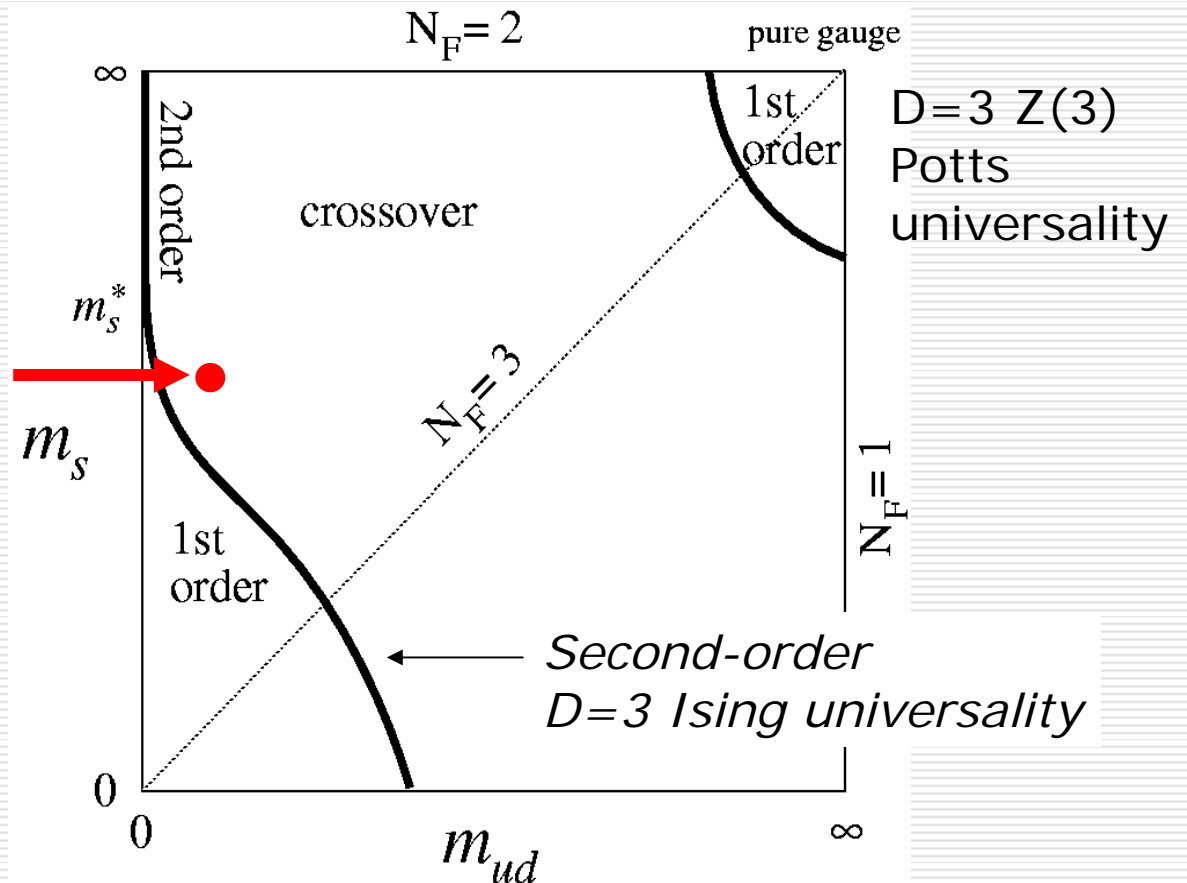
QGP transition at the physical point is a *crossover* according to staggered action (2006)

Wuppertal Group, Y. Aoki et al,
Nature 443 (2006) 675
Bielefeld-RBC-BNL Collaboration,
M. Cheng et al
HotQCD Collaboration

Transition temperature

$$T_c \approx 150 - 170 \text{ MeV}$$

- depends on the physical quantity
- still some debate on the value (Wuppertal vs Hot QCD)



No major change since 2006

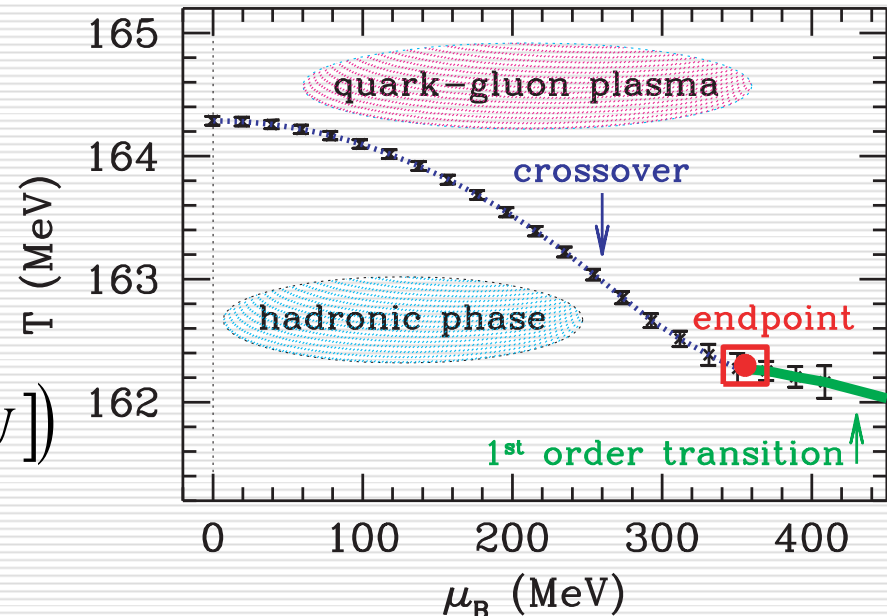


Status of finite-density QCD

- The “*sign problem*”, i.e., large phase fluctuation of the quark determinant $\det D$ for non-zero chemical potential

$$Z_{QCD} = \int \prod dU_{n\mu} \det D[U] \exp(-S_{gluon}[U])$$

- Slow but steady progress over the years for *not too large baryon density*:
 - Estimate of the end point of the 1st order line on the T - μ plane



2-parameter reweighting method:

Z. Fodor, S. Katz, JHEP 0404 (2004) 050

$N_f=2+1$, $N_t=4$

$$(T_E, \mu_E) = (162 \pm 2, 360 \pm 40) \text{ MeV}$$

Taylor expansion method:

C. Allton et al, Phys.Rev. D71 (2005) 054508

$N_f=2$, $L_t=4$



Canonical ensemble simulation ?

$$Z_{grand\ canonical}(T, m_q, \mu_B) = \sum e^{n_B \mu_B / T} Z_{canonical}(T, m_q, n_B)$$

$$Z_{canonical}(n_B, T, m_q) = \int [dU] \left[\int_0^{2\pi} d\varphi e^{-i3n_B \varphi} \det D(U, m_q, \mu = i\varphi T) \right] e^{-S_{gluon}(U)}$$

Projects out the states with baryon number n_B

- Vital to accurately estimate the projection of the quark determinant
- Some previous attempts:
 - Exact evaluation of the projection;
computationally expensive for large lattices
Forcrand-Kratochvila
hep-lat/0602024
 - Saddle point approximation;
controlling the approximation is not easy
S. Ejiri
ArXiv 0804.3227



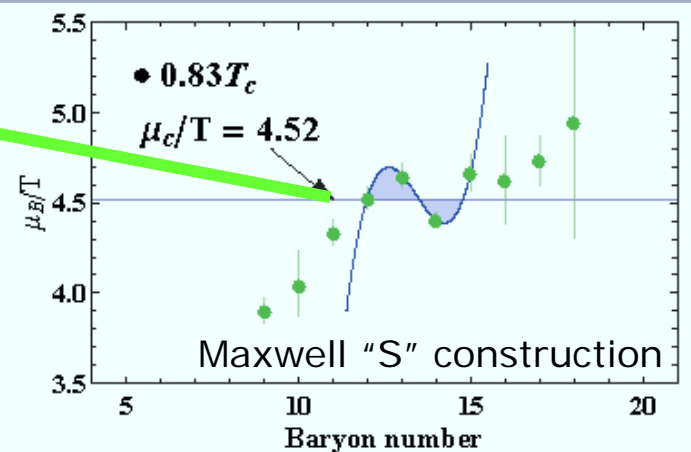
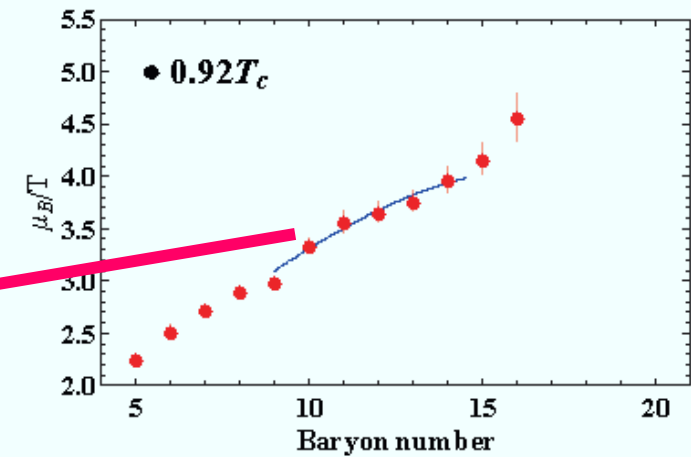
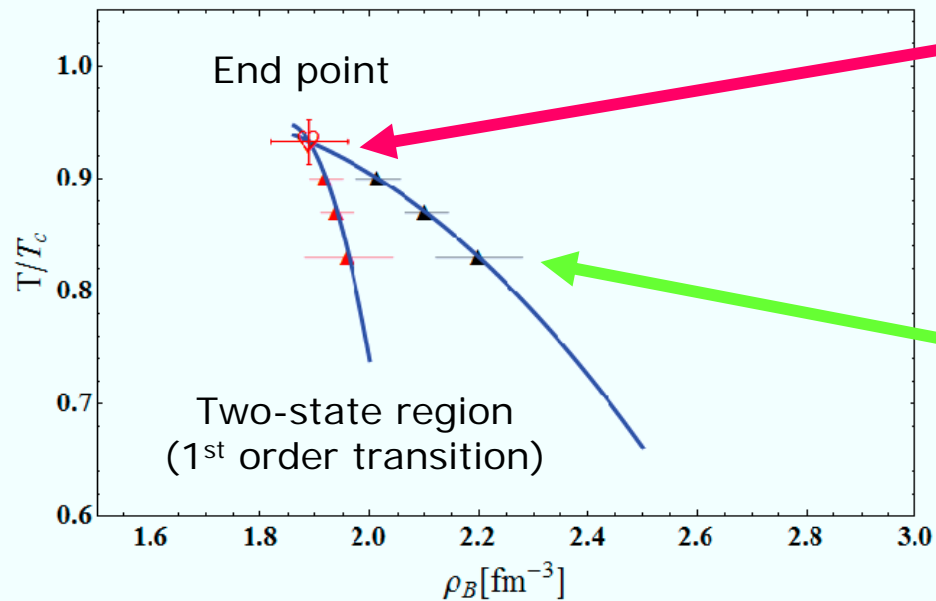
This year's attempt

Anyi Li et al @Lattice 2009

- Expand logdetD since detD is an extensive quantity

$$\ln \det D(U, m_q, \mu) = \sum_{n_B} e^{n_B \mu_B / T} A(U, m_q, n_B)$$

phase diagram for $N_f = 3$ $m_\pi \approx 700 \text{ MeV}$ $6^3 \times 4$





Nuclear physics



Two avenues from quarks to nuclei

- Effective theory approach
 - Extract nucleon two-body, three-body, ..., potentials via lattice QCD simulations
 - Use the potentials in conventional nuclear physics calculational schemes
 - S. Aoki, T. Hatsuda, N. Ishii (2007) based on the method developed by N. Ishizuka et al (2005)

- Direct approach
 - Calculate multi-quark Green's functions and directly extract the properties of nuclei, e.g., binding energies etc
 - Y. Kuramashi, M. Fukugita, A. Ukawa et al for nucleon-nucleon scattering lengths (1993)
 - T. Yamazaki, Y. Kuramashi, A. Ukawa for He4, He3 (2009)



Nuclear force from lattice QCD(2007)

N. Ishii, S. Aoki, T. Hatsuda, PRL 99, 022001 (2007)

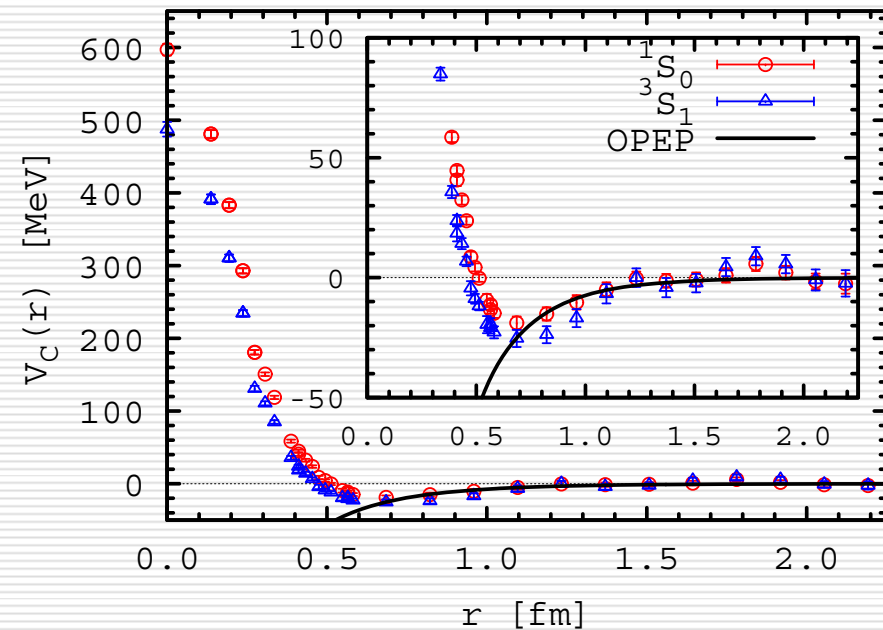
- 2-nucleon BS amplitude from lattice QCD

$$\phi(r) = \frac{1}{L^3} \sum_{\vec{x} \in L^3} \langle 0 | N(\vec{x} + r) N(\vec{x}) | NN \rangle$$

- Extraction of potential from an effective Schrodinger eq.

$$V(r) = E + \frac{1}{2\mu} \frac{\nabla^2 \phi(r)}{\phi(r)}$$

- Impact and prospects
 - Derivation of the hard core
 - Extension to hyperon-nucleon potential etc



Quenched QCD 32^4 lattice

$$m_\pi / m_\rho = 0.595$$



He nuclei directly from lattice QCD

T. Yamazaki, Y. Kuramashi, A. Ukawa, arXiv:0912.1383 (2009)

□ Methodology

- Define He operator in terms of quark fields

$$p(x) = \varepsilon^{abc} (u^a(x) C \gamma_5 d^b(x)) u^c(x)$$

$$n(x) = \varepsilon^{abc} (d^a(x) C \gamma_5 u^b(x)) d^c(x)$$

$$He^4(x) = \text{spin projection} (p(x)p(x)n(x)n(x)) \quad \text{J. E. Beam, Phys. Rev. 158, 907 (1968)}$$

- Calculate He Green's function to extract the binding energy

$$G_{He^4}(t) = \frac{\langle 0 | He^4(t) He^4(0) | 0 \rangle}{\langle 0 | p(t) p(0) | 0 \rangle^2 \langle 0 | n(t) n(0) | 0 \rangle^2}$$

$$\xrightarrow{t \rightarrow \infty} Z \exp\left(-\underbrace{(M_{He^4} - 2m_p - 2m_n)}_{-B_{He^4}} t\right)$$

$-B_{He^4}$ Binding energy of He4 nuclei



Three difficulties

- Statistical error

$$\frac{\text{noise}}{\text{signal}} \propto \frac{1}{\sqrt{N_{\text{meas}}}} \exp\left(4\left(m_N - \frac{3}{2}m_\pi\right)t\right)$$

Use heavy pion and large statistics in quenched QCD

- Factorially large number of Wick contractions of quark operators

$$He^4 = p^2 n^2 = (udu)^2 (dud)^2 = u^6 d^6$$

$$N_u! \times N_d! = (2N_p + N_d)! \times (N_p + 2N_d)!$$

Use symmetries to remove identical contractions

$$\text{He3 } 5! \times 4! = 2880$$

$$\text{He4 } 6! \times 6! = 518400$$

- Identification of bound states
Bound state or scattering state?

Use a set of spatial sizes



Comments on Wick contractions

- Symmetries
 $p \leftrightarrow p$, $n \leftrightarrow n$ in He operator
Isospin all $p \leftrightarrow$ all n
- Simultaneous calculation of two contractions
 $u \leftrightarrow u$ in p or $d \leftrightarrow d$ in n

Result of reduction:

□ He

$$\frac{518400}{2_{ISO} \cdot 2_{u \leftrightarrow u}^2 \cdot 2_{d \leftrightarrow d}^2 \cdot \left(2_{p \leftrightarrow p} \cdot 2_{n \leftrightarrow n}\right)_{src, sink}^2} \xrightarrow{\text{w/o double count}} 1107$$

□ He3

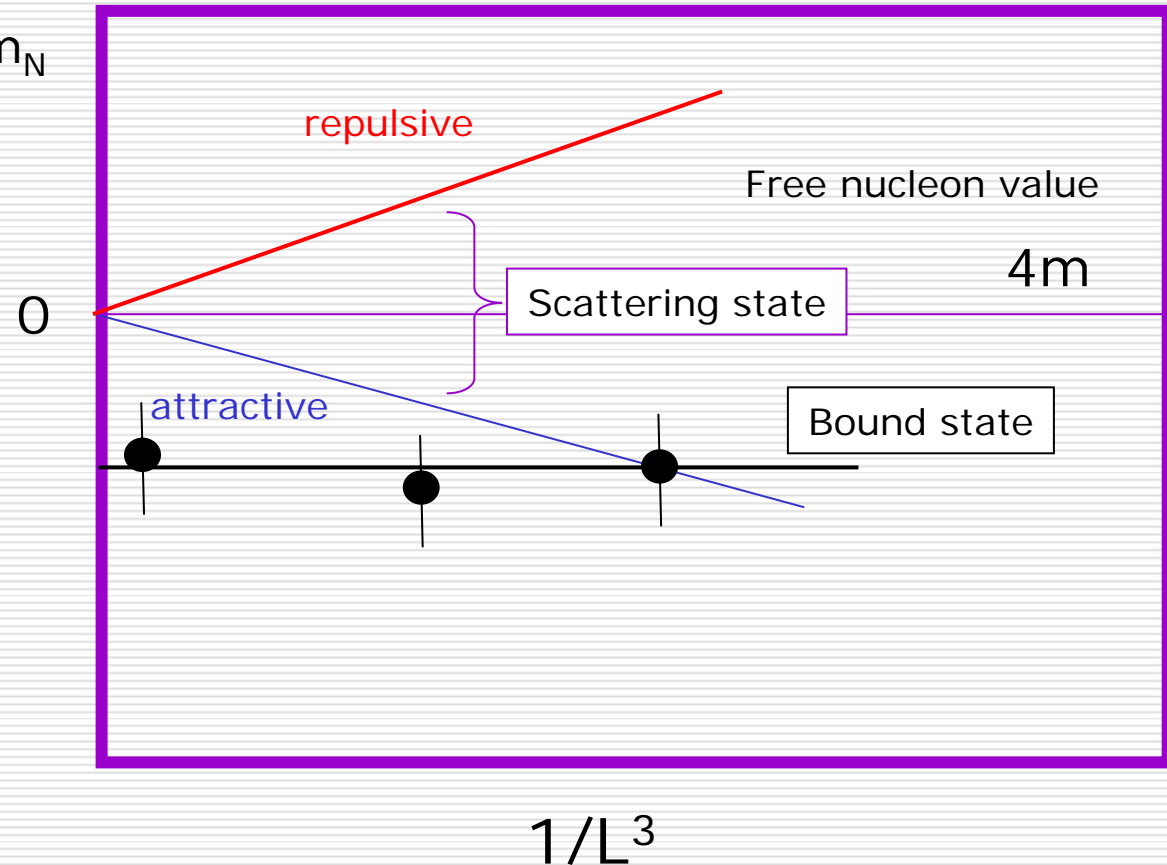
$$\frac{2880}{2_{u \leftrightarrow u}^2 \cdot 2_{d \leftrightarrow d} \cdot \left(2_{p \leftrightarrow p}\right)_{src, sink}^2} \xrightarrow{\text{w/o double count}} 93$$



Bound state or scattering state?

- Measurement for a single spatial volume cannot distinguish a bound state from scattering states
- Use multiple volumes for measurements

$$M_{\text{He}} - 4m_N$$





simulation

□ Quenched calculation

- Iwasaki gauge action $\beta = 2.416$ $1/a = 1.54 \text{ GeV}$
- Tadpole-improved Wilson quark action
 $m_\pi = 0.8 \text{ GeV}$, $m_N = 1.62 \text{ GeV}$

□ Lattice sizes and #measurements

L	L(fm)	#conf	#meas
24	3.1	2500	2
48	6.1	400	12
96	12.3	200	12

□ Source smearing

(A, B) = (0.5, 0.5), (0.5, 1.0) for L=24

(0.5, 0.5), (1.0, 0.4) for L=48, 96

$$q(\vec{x}) = A \exp(-B|\vec{x}|)$$

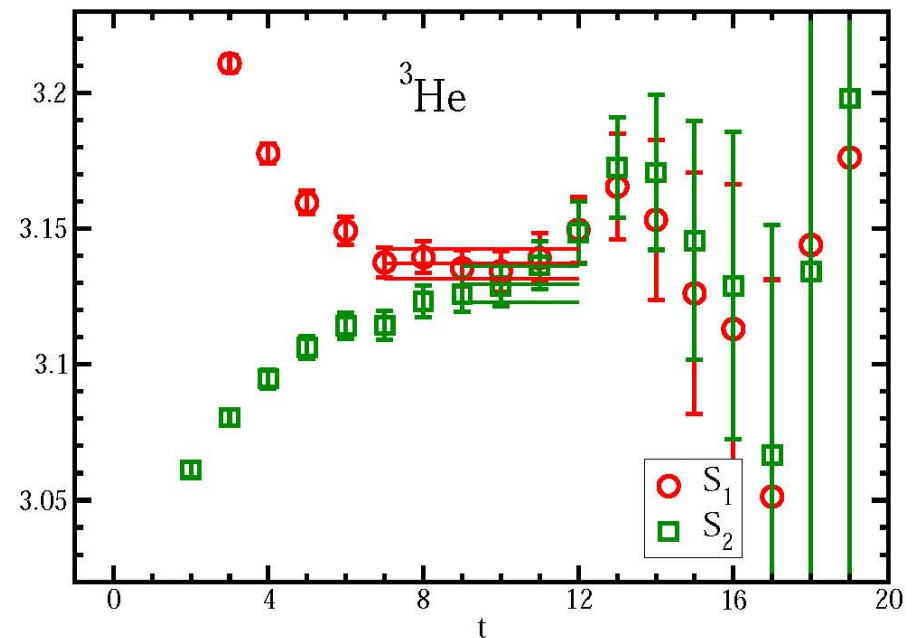
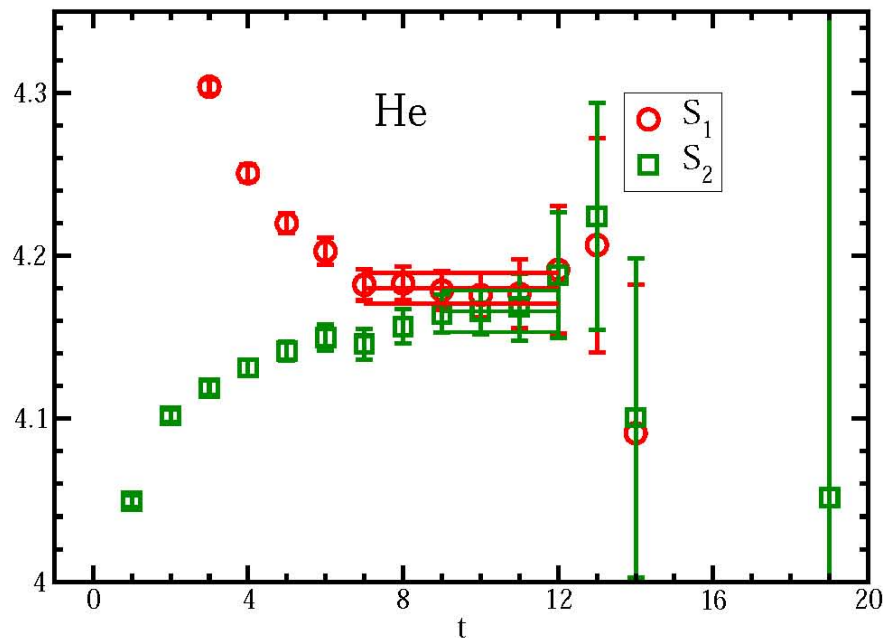
S_1

S_2



Results for effective mass $L=48$

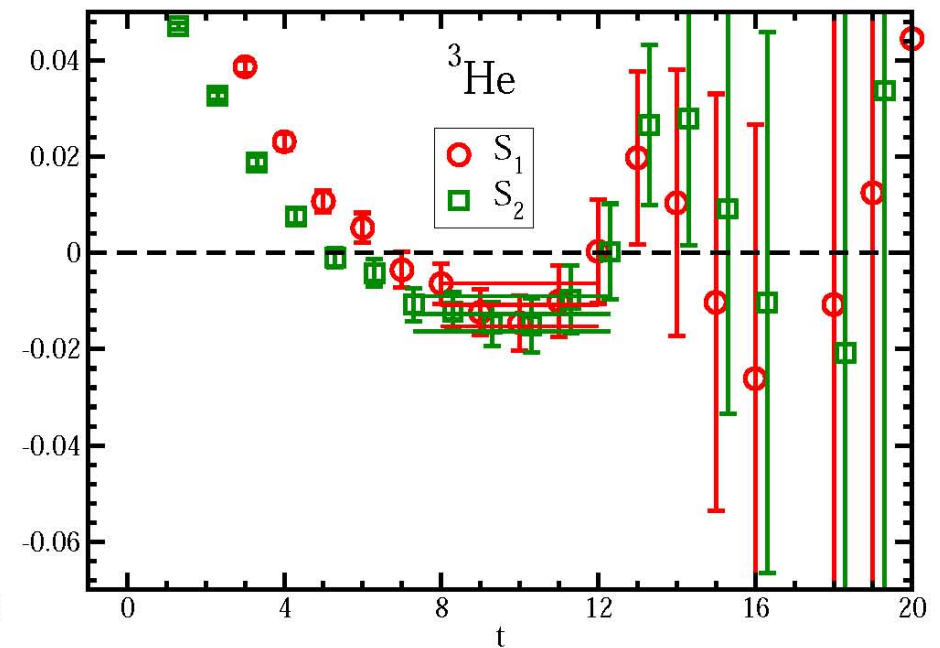
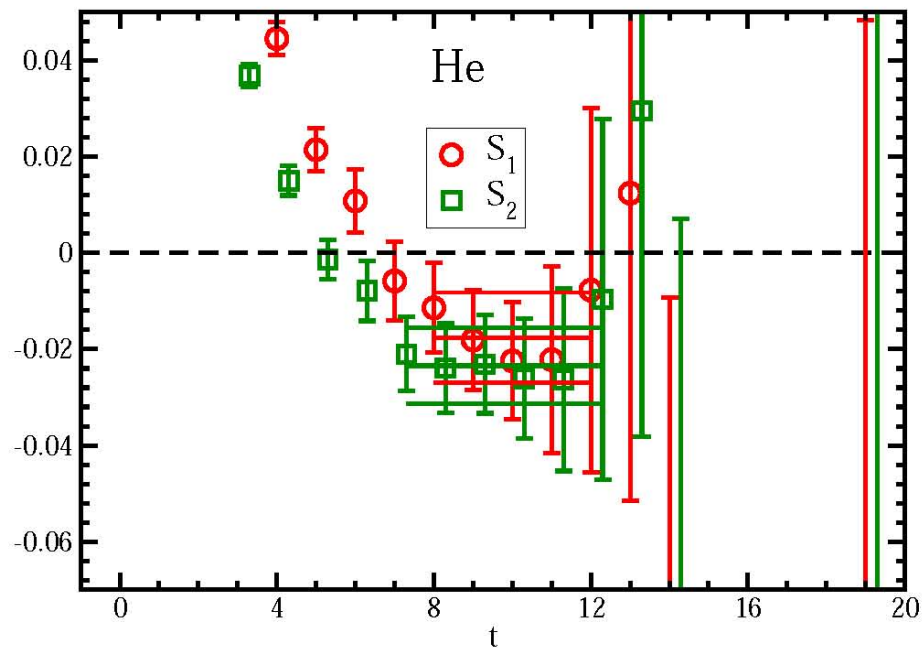
$$m_{eff}(t) = \log \frac{C(t)}{C(t+1)} \xrightarrow{t \rightarrow \infty} m$$





Results for effective energy shift $L=48$

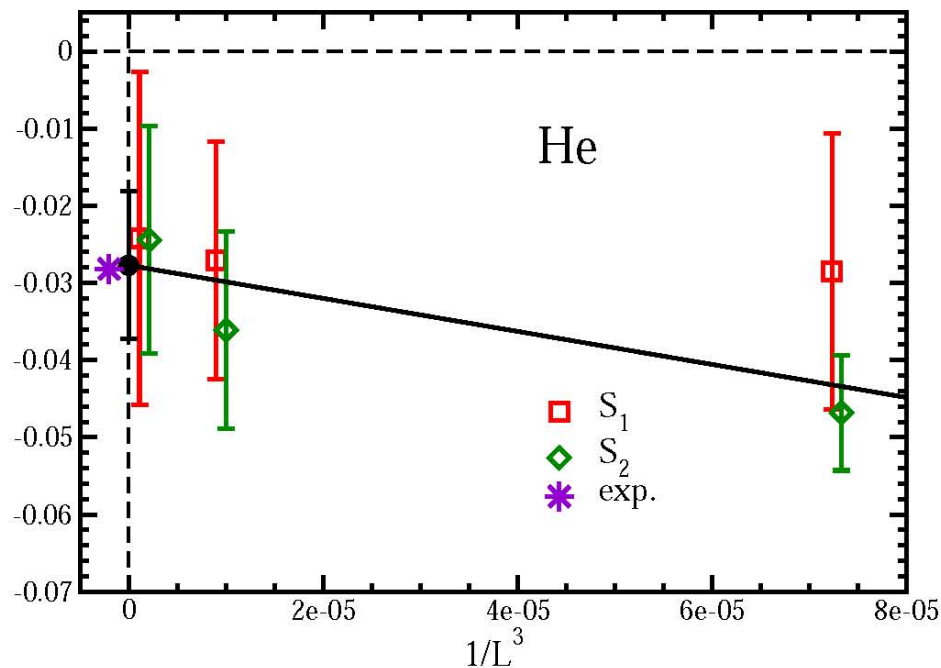
$$B_{eff}(t) = \log \frac{G_{He}(t)}{G_{he}(t+1)} \xrightarrow{t \rightarrow \infty} -B_{He}$$



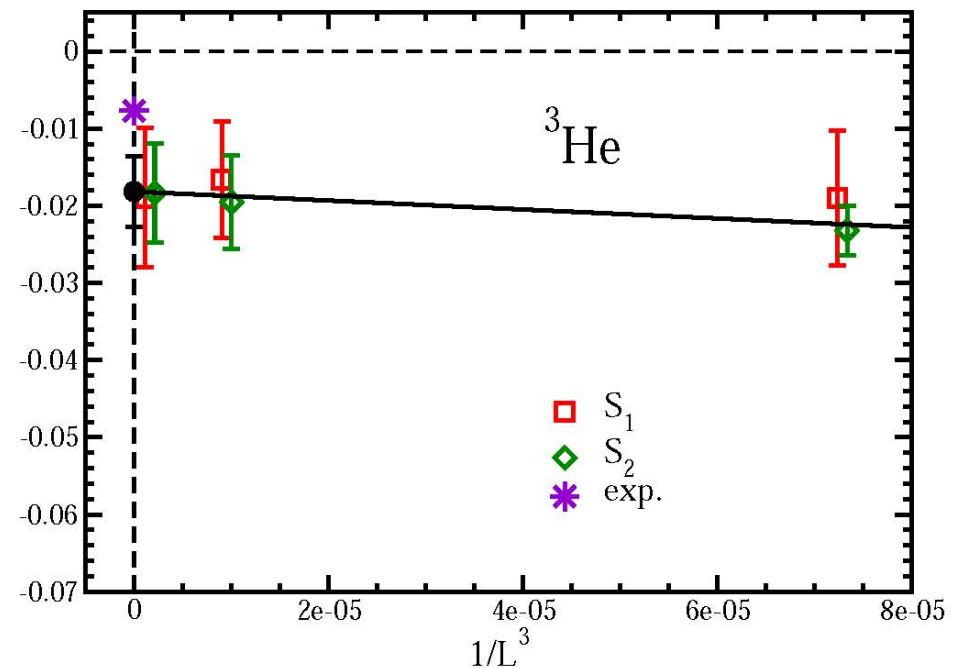


Results for size dependence

- Small volume dependence
- Non-zero value(!) in the infinite volume limit



$$\Delta E_{\text{He}} = 27.7(7.8)(5.5) \text{ MeV}$$



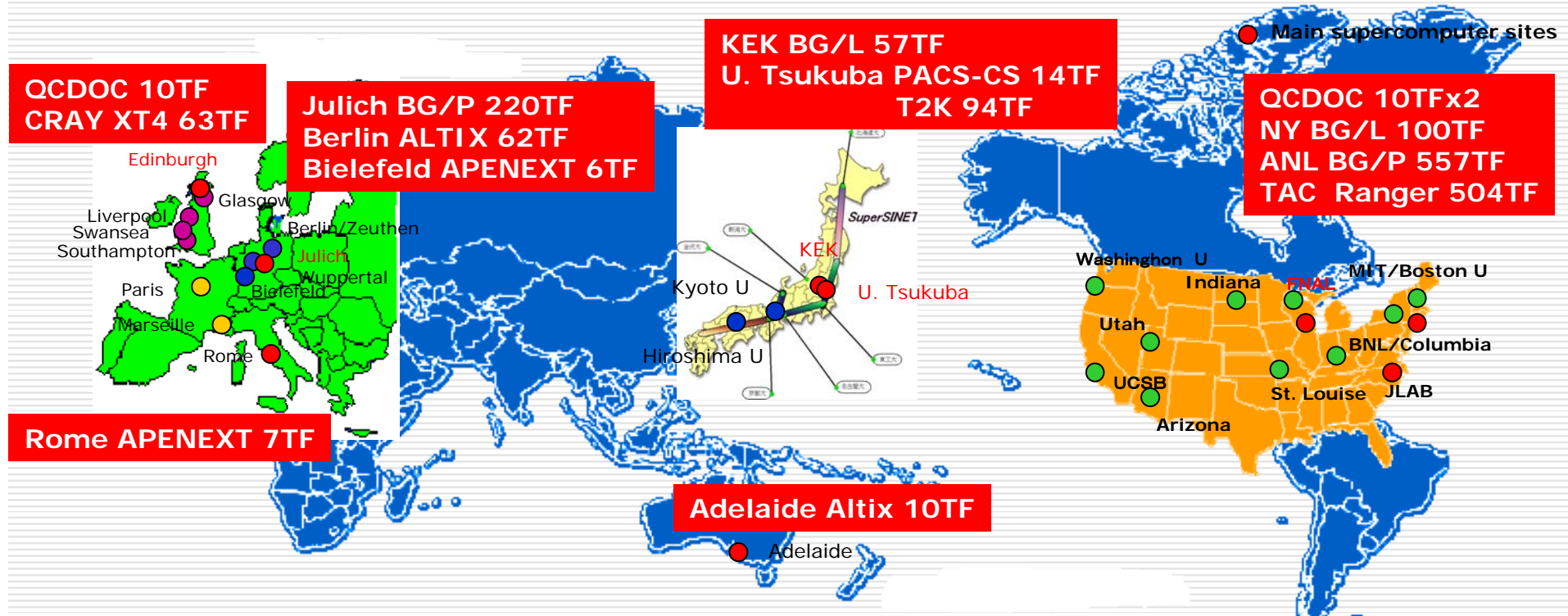
$$\Delta E_{^3\text{He}} = 18.2(3.5)(2.9) \text{ MeV}$$



Computer trends and ILDG



Current computing resources for lattice QCD in the World



- About a dozen major sites scattered in USA, EU(UK, Germany, Italy etc), Japan
- In total 500~600Tflops in peak speed (US300Tf, EU150Tf, Japan100Tf); about 3% of World HPC resources



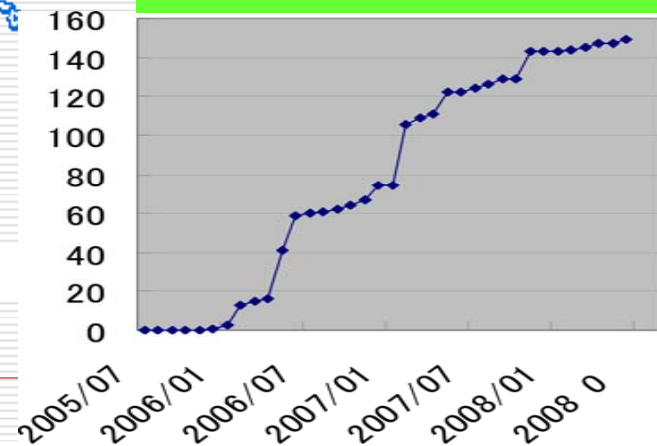
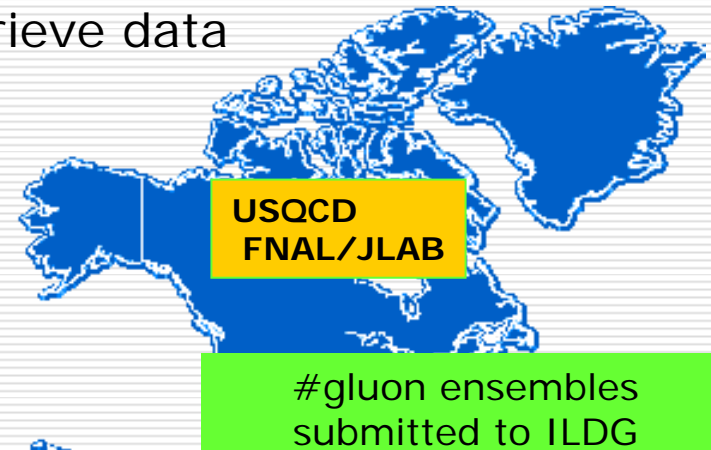
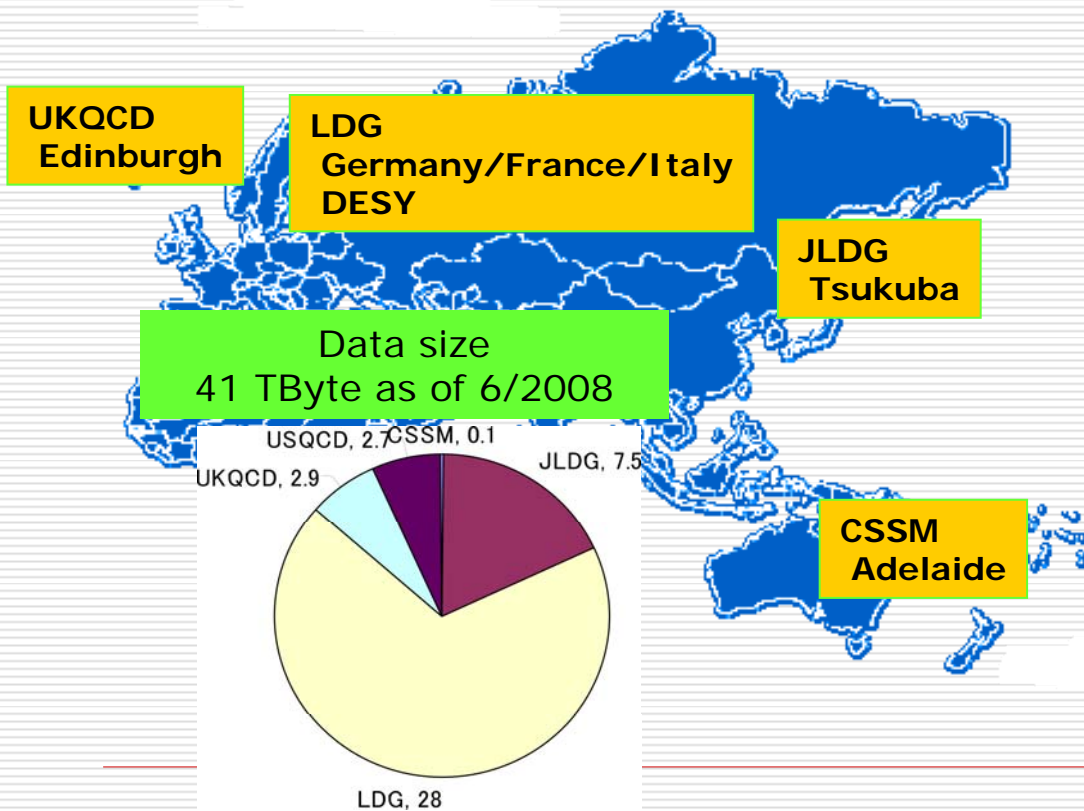
International Lattice Data Grid

In preparation since 2002 / in operation since 2006

- World-wide sharing of gluon configurations
- Grid of regional grids
 - Standardized xml to describe data
 - Common interface to search/retrieve data

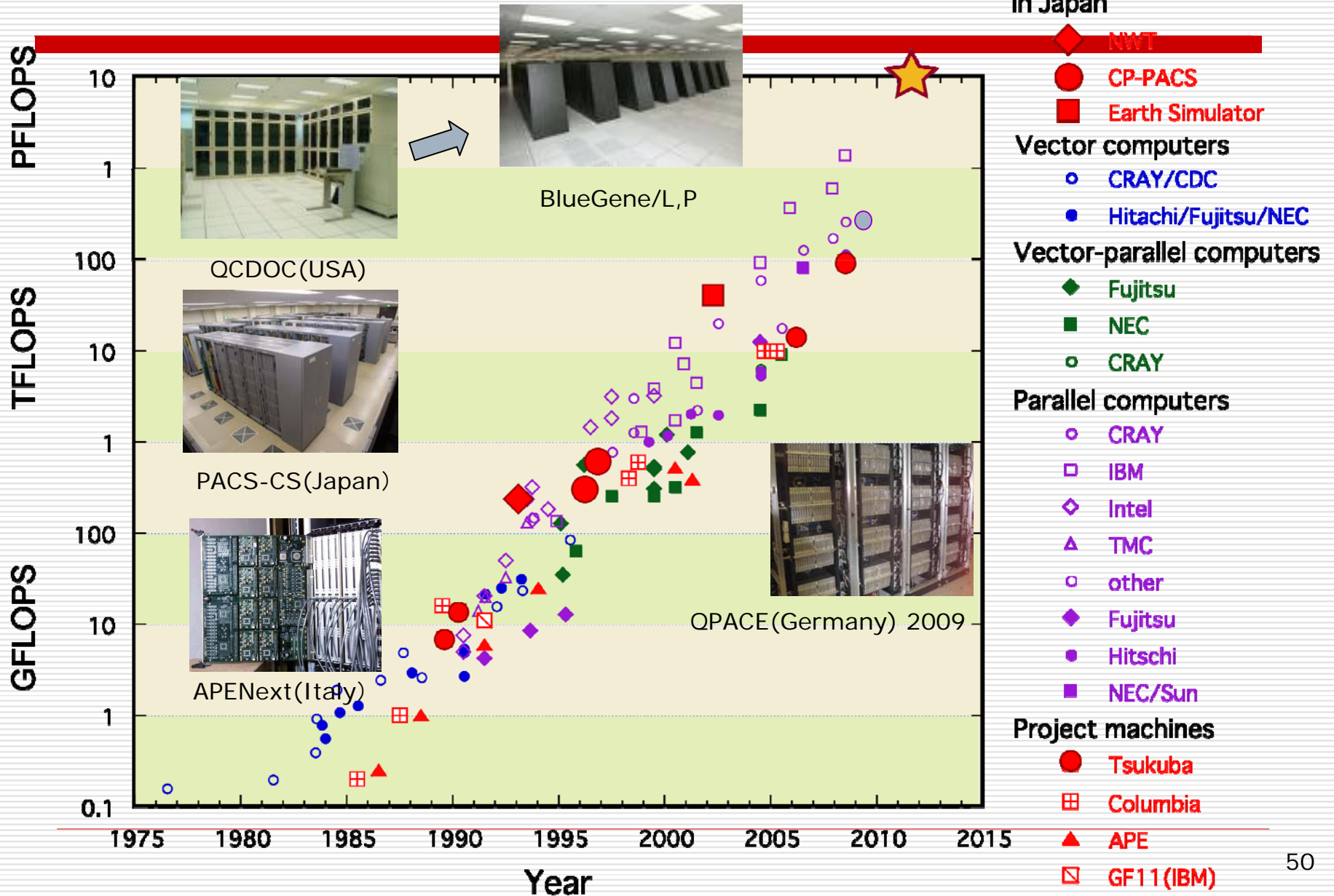


<http://ildg.sasr.edu.au/Plone>





Development of supercomputers and QCD dedicated computers





Future: petascale computing (2010-2015)

- Peta-scale computing is around the corner
 - “National” projects
 - USA: Bluewater, BlueGene/Q, ...
 - Japanese Petaflops Project
 - clusters based on commercial multi-core CPU (Intel, AMD)
- New projects for lattice QCD
 - QPACE Project (QCD Parallel Computing on the CELL)
 - CELL-based cluster/200Tflops in 2009
 - GPGPU?
 - Many-core high speed graphic cards/software development



O(10-100) enhancement will allow physical point simulation on larger volumes ($L=3\text{fm} \rightarrow 6\text{fm}$) / smaller lattice spacings ($a=0.1\text{fm} \rightarrow 0.05\text{fm}$)



Status of the Japanese Next Generation Supercomputer Project



Political turmoil last year

16 September:

Hatoyama Cabinet took office (first real change of power since 1951)

18 September:

“Government Revitalization Unit (GRU) ” is set up; starts reexamination of F2010 budget

13 November:

GRU Working Group, after an 1-hour public hearing, recommends suspension of the Supercomputer Project ; many science & technology budget also recommended cut.

Late November :

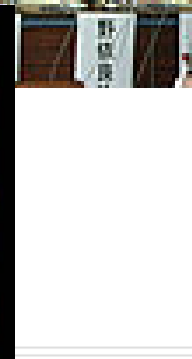
appeals by many academic societies and universities against the cut

8 December:

Science and Technology Council recommends continuation of the Project

16 December:

Government decides to proceed with the Project





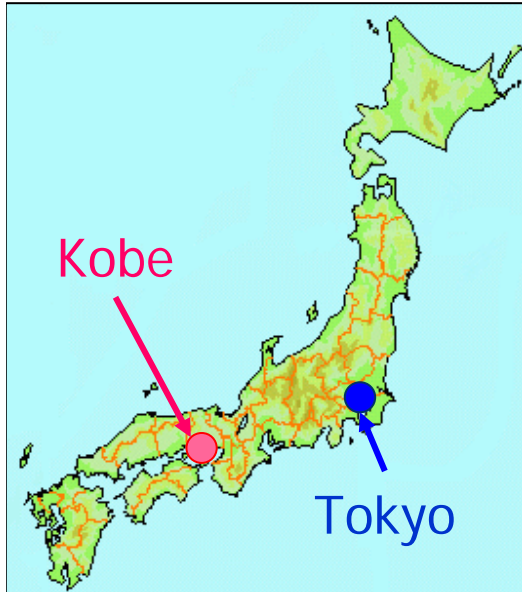
Modification of the project

- Machine R&D schedule
 - Original target date: November 2011
 - New target date: June 2012

- New target of the project
 - “High Performance Computing Infrastructure (HPCI) Project”
 - Buildup of HPCI in Japan by connecting the next-generation supercomputer and other supercomputers in Japan
 - Buildup of “Consortium” for the best effective use of HPCI resources in Japan



Cite for the Next Generation Supercomputer



450km (280miles)
west from Tokyo





Building construction

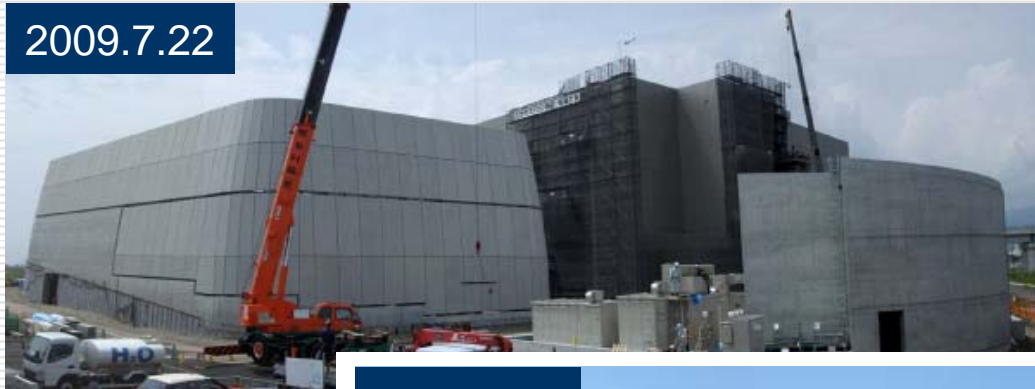
2008.11.12



2009.3.26



2009.7.22



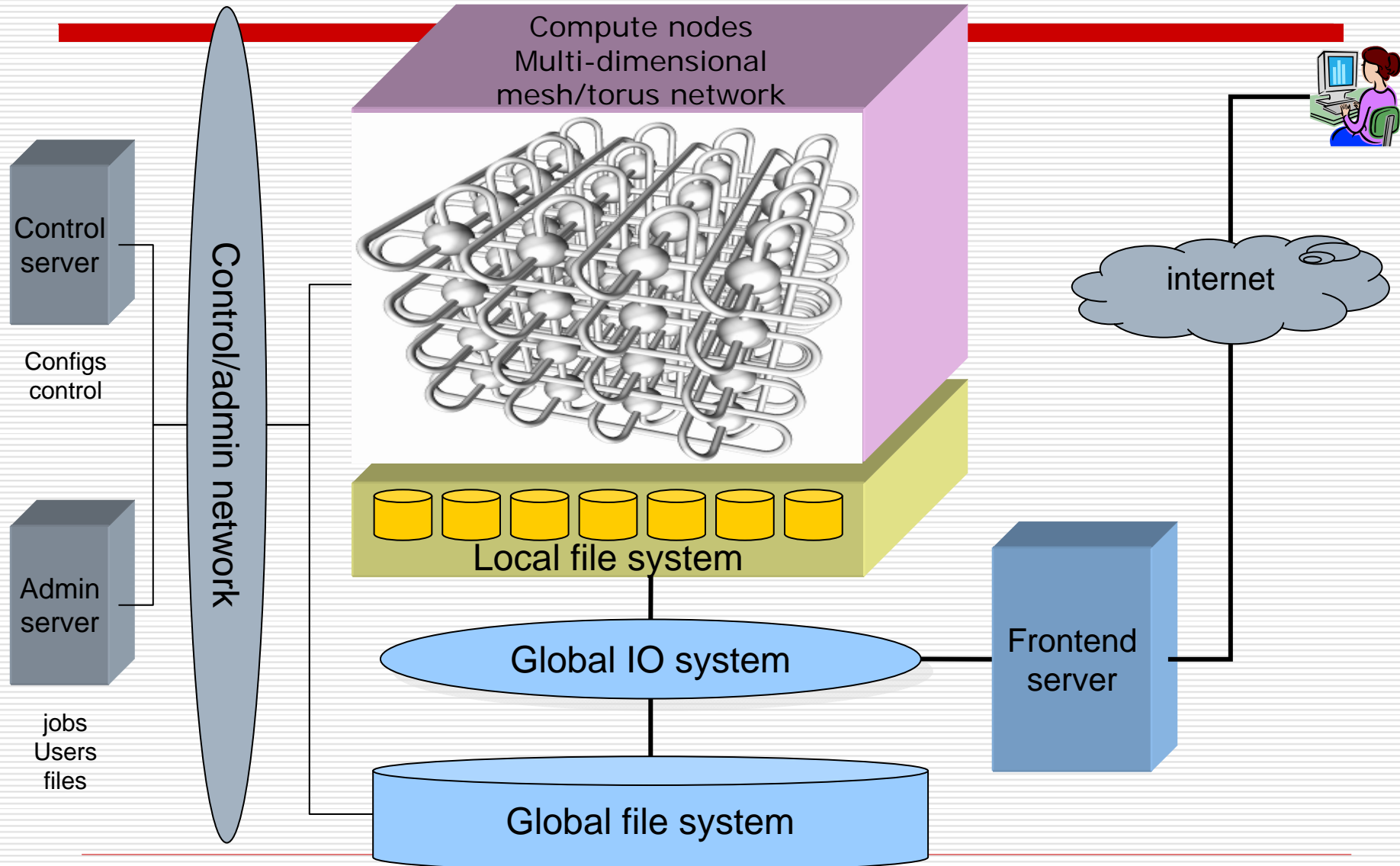
2009.9.1



To be completed
in May 2010



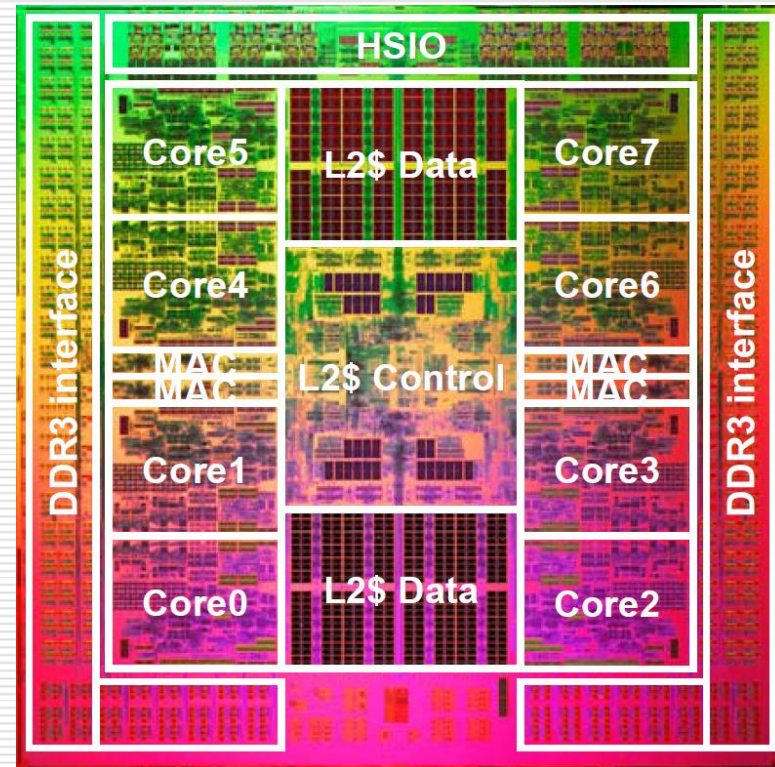
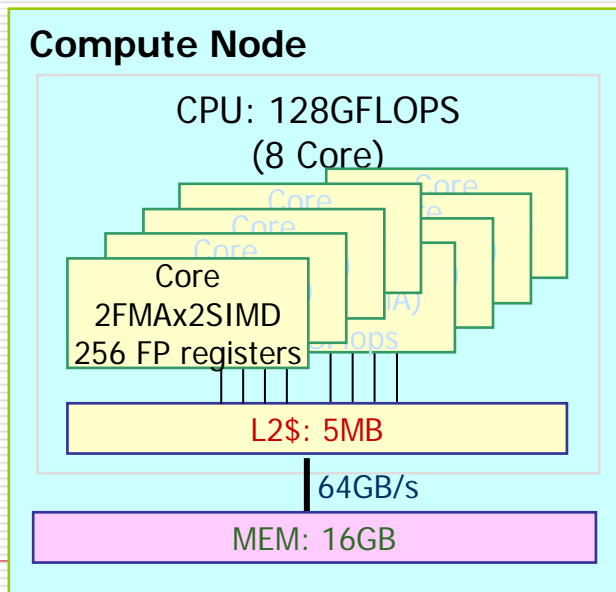
System overview





Compute node

- CPU: SPARC64™ VIIIfx
 - 8 core 128Gflops
 - DDR3 Memory 64GB/sec
 - Power consumption 58W (water cooled at 30°C)
 - 45nm CMOS

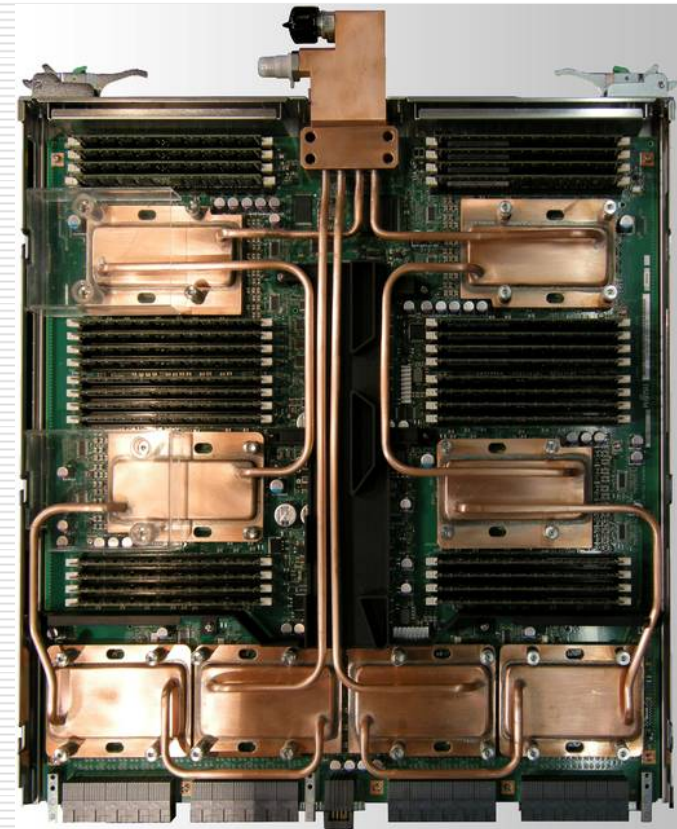
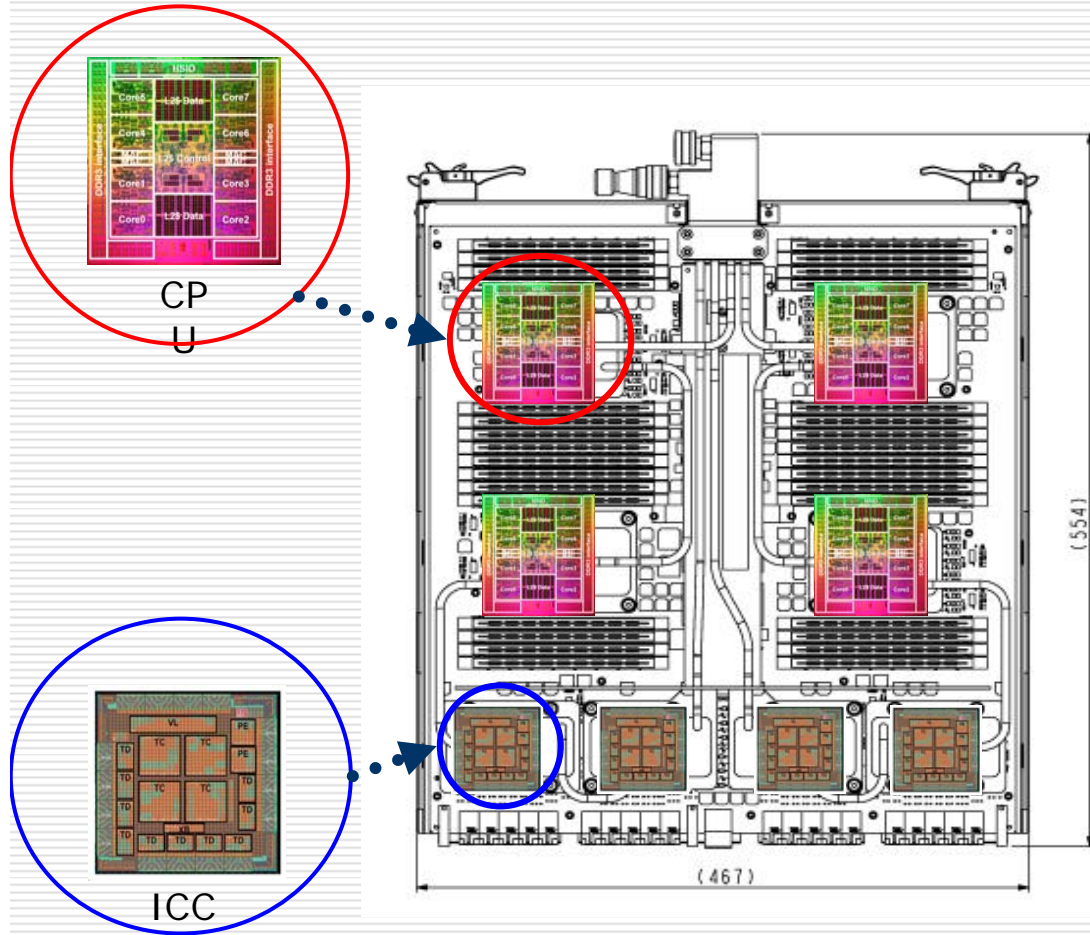


22.7mm x 22.6mm
760 M transistors

Photo courtesy FUJITSU Limited



System board

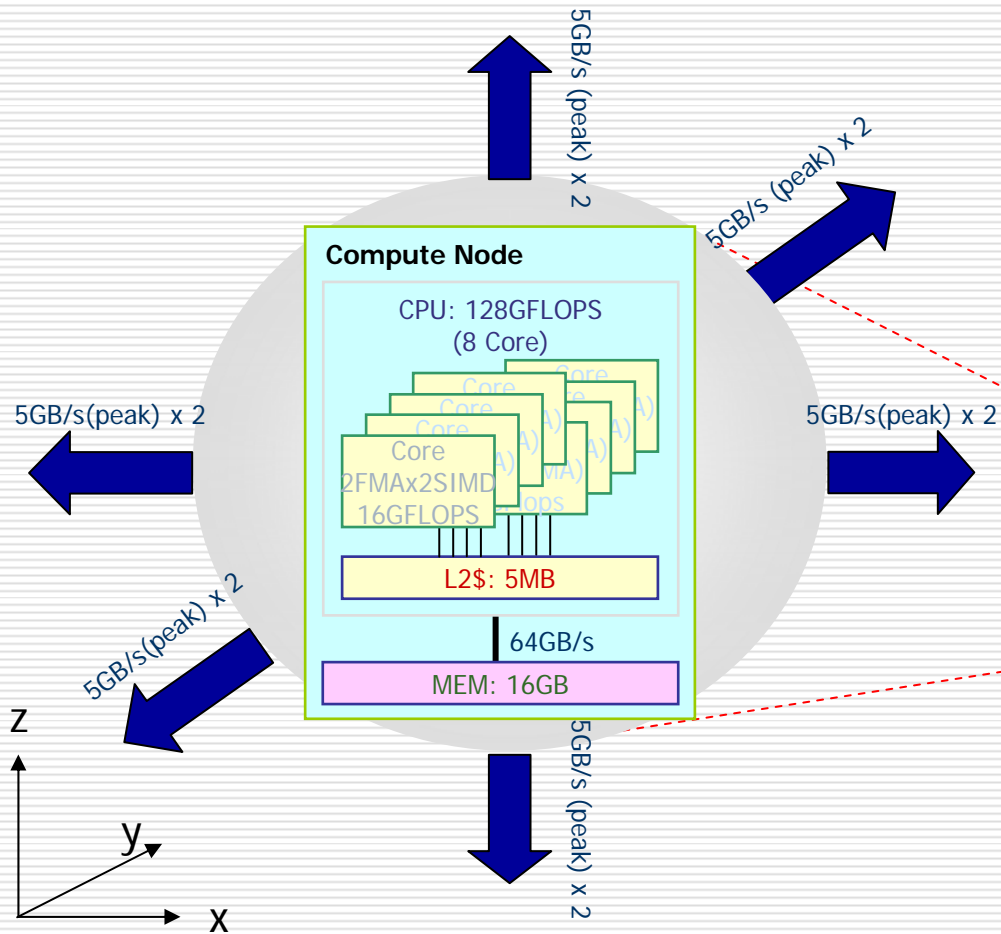


4 CPU and Network Interface/board

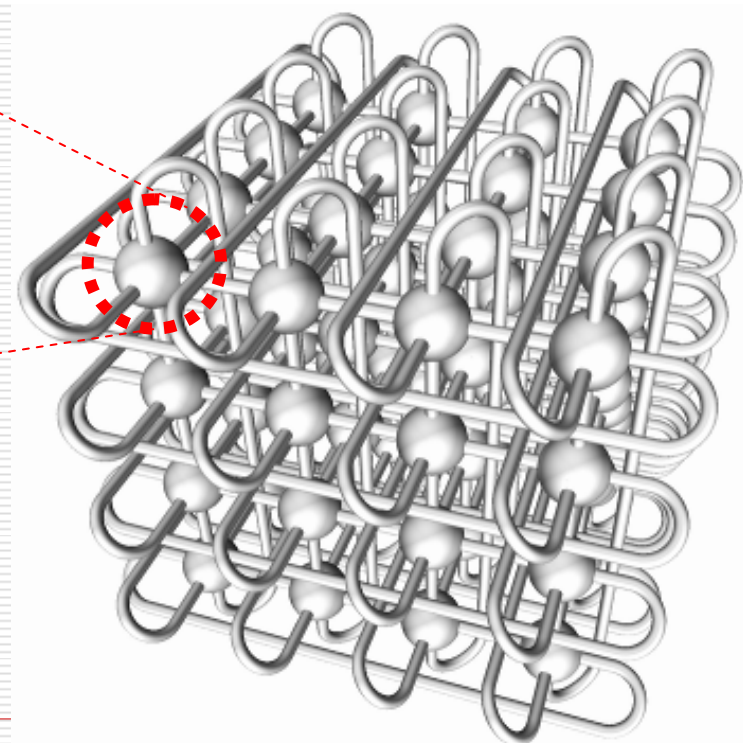
Photo courtesy FUJITSU Limited



Network configuration



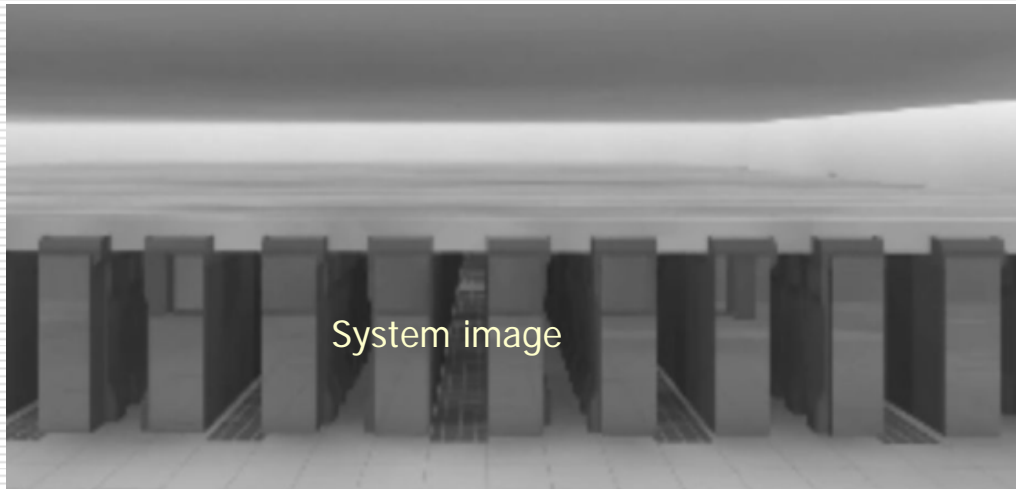
- ❑ Logical 3d torus/physical multi-d torus/mesh
- ❑ 5GB/s x2 for each 3d link





Rack prototype

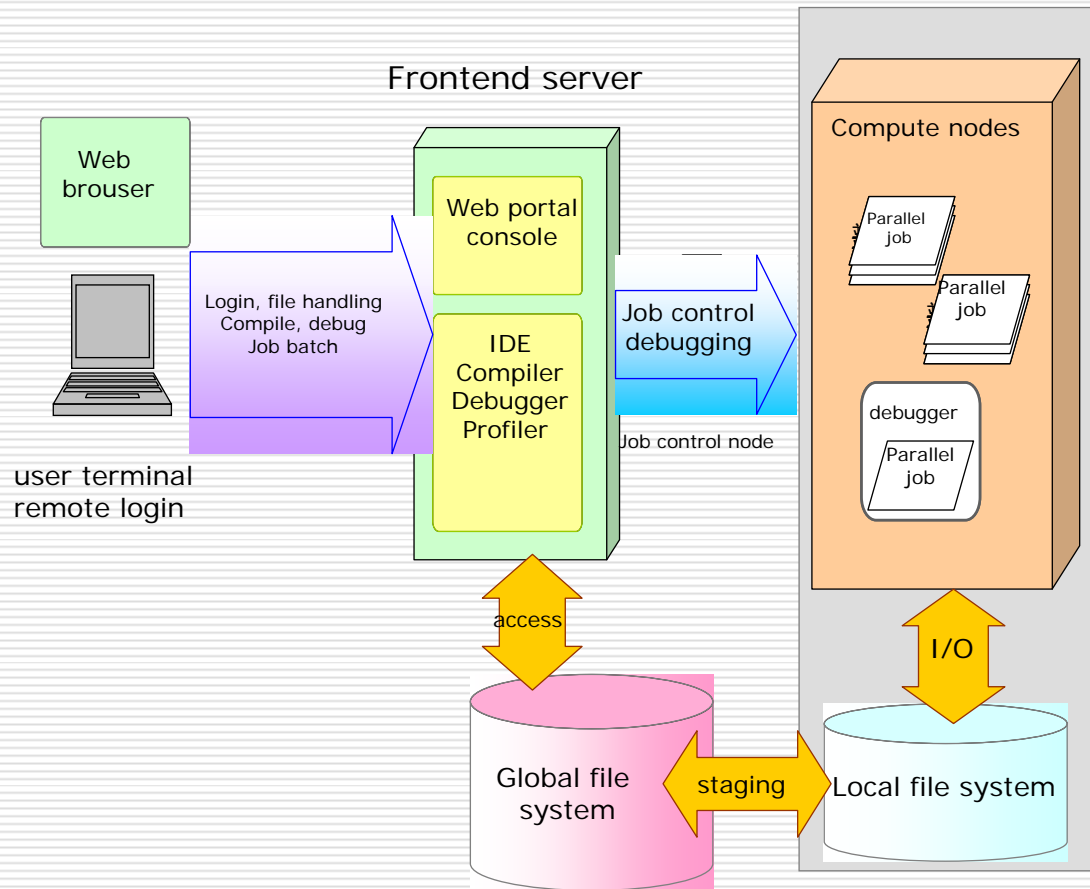
- ❑ 96 nodes (24 system boards) and system disk
- ❑ 796mm × 750mm × 2060mm





Software and job environment

- System software
 - OS: Linux
 - Programming
 - Fortran, C, C++
 - MPI library for communication
- Distributed file system
 - Stage in/out to/from local disk
 - File sharing
- Batch-based job execution
 - Interactive debugger (planned)





Research system of the Project

MEXT

Japanese Government



Strategy Committee

■ 5 strategic areas have priority of use

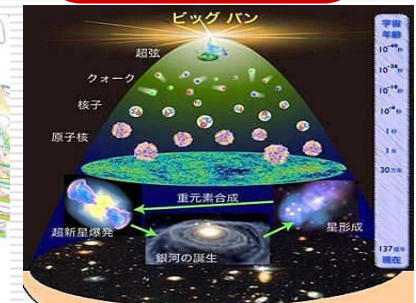
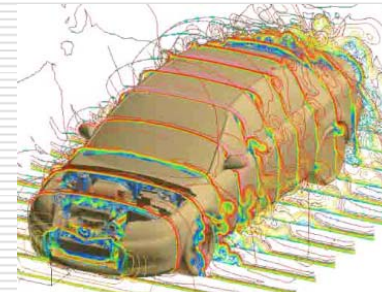
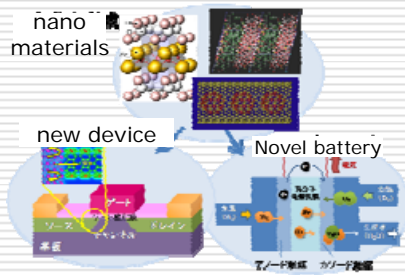
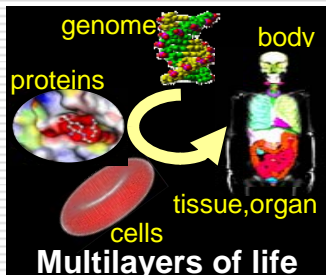
Life Science & Medicine

New materials & Energy

Global change prediction

Next generation Engineering

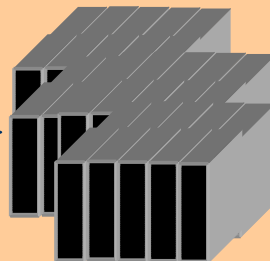
Matter & Universe



Next generation supercomputer

RIKEN

operation



Advanced Institute for Computational Science (under planning)



Conclusions



Where we stand now

- *Realistic calculation directly at the physical point is finally reality*
 - Fruit of continuous effort over 25 years toward:
Better physics understanding
Better algorithms
More powerful machines

- *Change of philosophy from “simulation” to “calculation”*

If lattice spacing sufficiently small,

 - No more approximations/extrapolations
 - Gluon configuration produced is Nature itself



Where do we go

- *Expect that the fundamental issues of lattice QCD as particle theory makes major progress over the next five year range*
 - Single hadron properties and fundamental constants
 - Precision flavor physics ($<1\%$) and old issues such as $K \rightarrow \pi \pi$ decays
 - Hot/dense QCD with chiral lattice action on large lattices
- *Vast area of multi-hadron systems/atomic nuclei lies in wait for nuclear physics colleagues to explore*
 - Nuclear force from lattice QCD
 - Exotic nuclei with unusual n/p ratios/strangeness etc
 - Even direct computation bypassing nuclear theory!

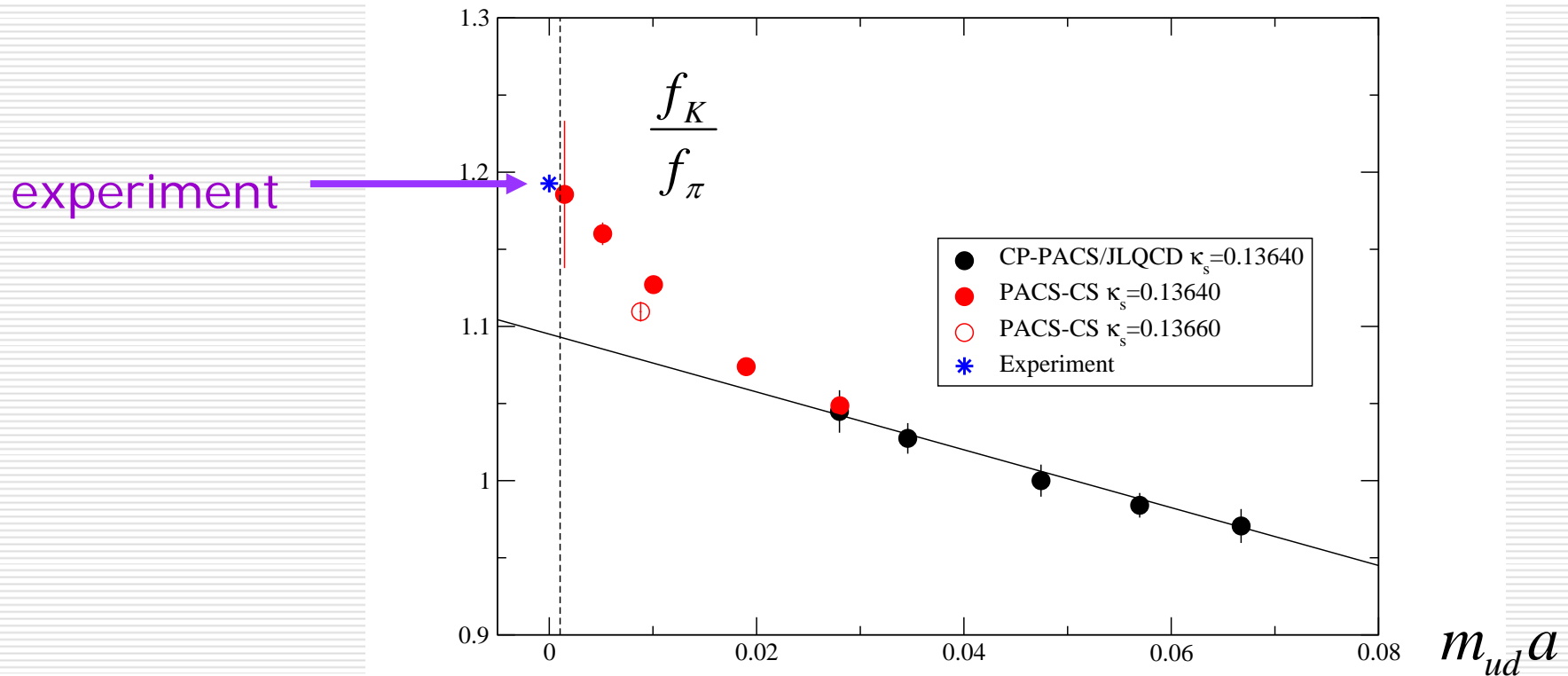


Supplementary slides



Another example

Decay constant of K and π



$$\frac{f_K}{f_\pi} \propto 1 + \frac{5}{4} \frac{2Bm_{ud}}{(4\pi f)^2} \log \frac{2Bm_{ud}}{\mu^2} + \dots$$



u/d quark mass difference

Handling of EM effects is the issue

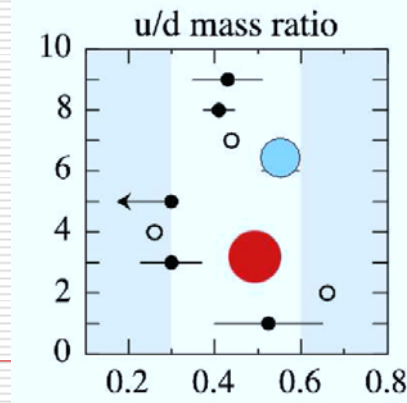
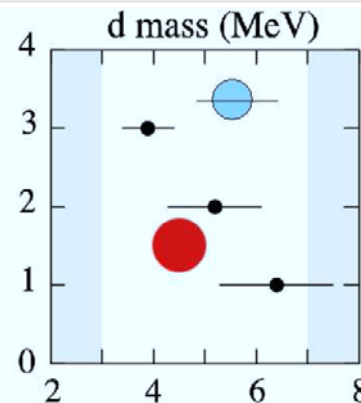
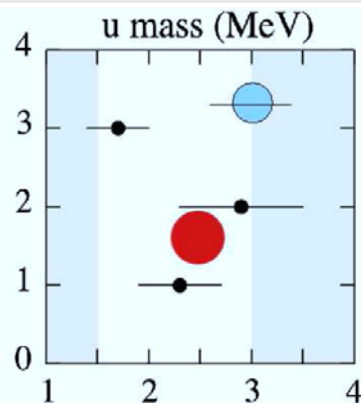
- Use Dashen's theorem to estimate the EM effects, and obtain m_u/m_d from K^0, K^+ masses

Lattice09	m_u (MeV)	m_d (MeV)	m_u / m_d
MILC	1.96 ± 0.17	4.53 ± 0.27	0.432 ± 0.040
Aubin et al	1.7 ± 0.3	4.4 ± 0.5	0.39 ± 0.5

- Calculate EM effects by coupling (quenched)QED
RBC/UKQCD@Lattice 2009 "study of isospin breaking effects"

Blue: $N_f=2$

Red: $N_f=2+1$





(non-Lattice) comment on ϵ_K

Buras-Guadagnoli ArXiv 0805.3887

$$\epsilon_K = e^{i\phi_\epsilon} \sin \phi_\epsilon \left(\frac{\text{Im } M_{12}^K}{\Delta M_K} + \frac{\text{Im } A_0}{\text{Re } A_0} \right) = \kappa_\epsilon \frac{\text{Im } M_{12}^K}{\Delta M_K}$$

- Usually neglected since small. However, with improved estimates of B_K and V_{cb} , this correction is significant
- Estimate by Buras-Guadagnoli

$$\kappa_\epsilon \approx \sqrt{2} \sin \phi_\epsilon \left(1 - \frac{1}{\omega} \text{Re} \left(\frac{\epsilon_K'}{\epsilon_K} \right) + \frac{1}{\sqrt{2} |\epsilon_K|} \frac{\text{Im } A_2}{\text{Re } A_2} \right)$$

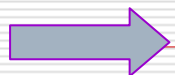
$$= 0.92 \pm 0.02$$

- Using quenched lattice QCD estimate for $\text{Im} A_2$ (RBC, PACS-CS etc) yields the same value

Van de Water @ Lattice2009

$$-6.4 \pm 6.4 \times 10^{-5}$$

Quenched lattice QCD estimate; assign 100% error



ϵ_K band might move up

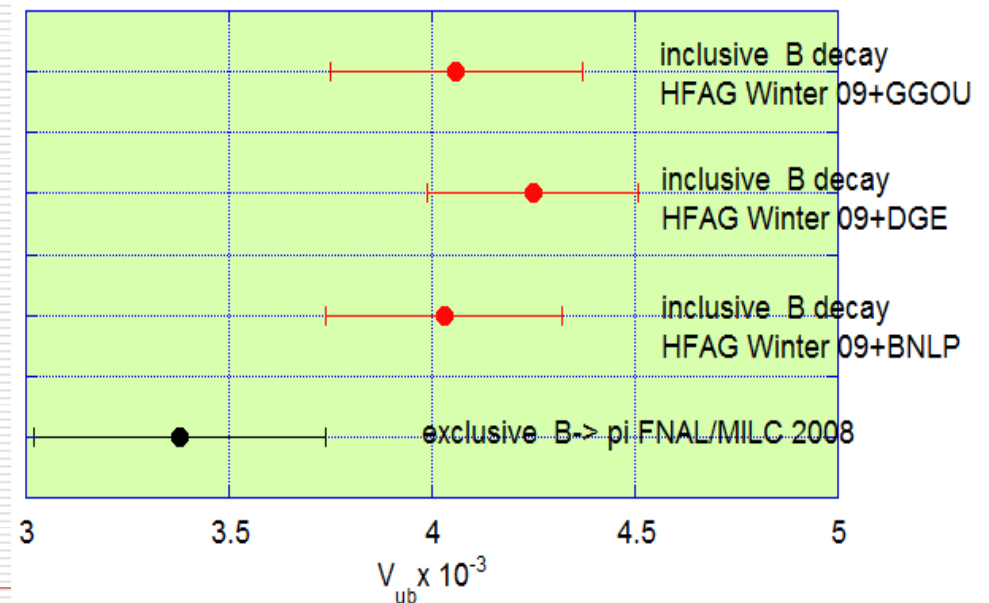
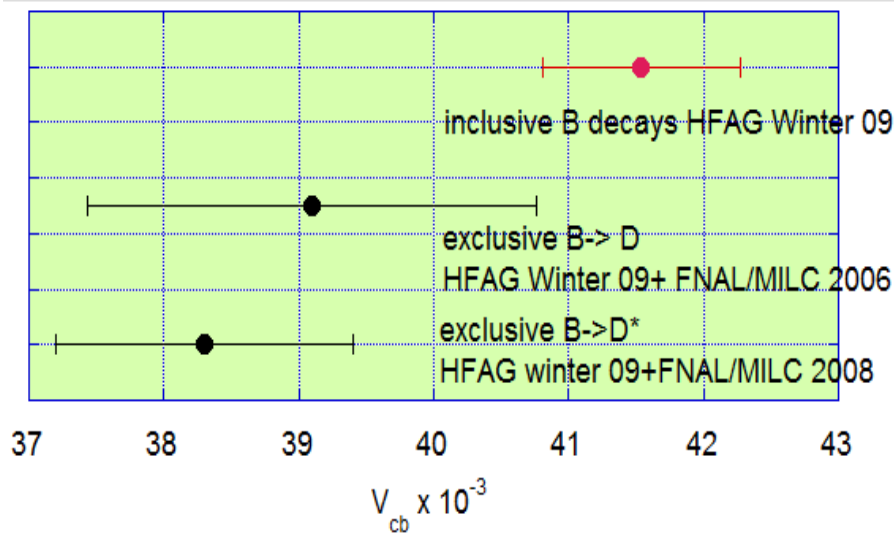


Inclusive vs exclusive determination of V_{cb} , V_{ub}

- No change of lattice numbers since 2008
- Inclusive (exp + non-lattice theory) values still differ from exclusive (exp + lattice) values at 2 sigma level

V_{cb}

V_{ub}





Status of D decay constants

Experiment 2009

CLEO ArXiv 0901.1216

$$f_D = 205.8 \pm 8.9 \text{ MeV}$$

$$f_{D_s} = 259.5 \pm 7.3 \text{ MeV}$$

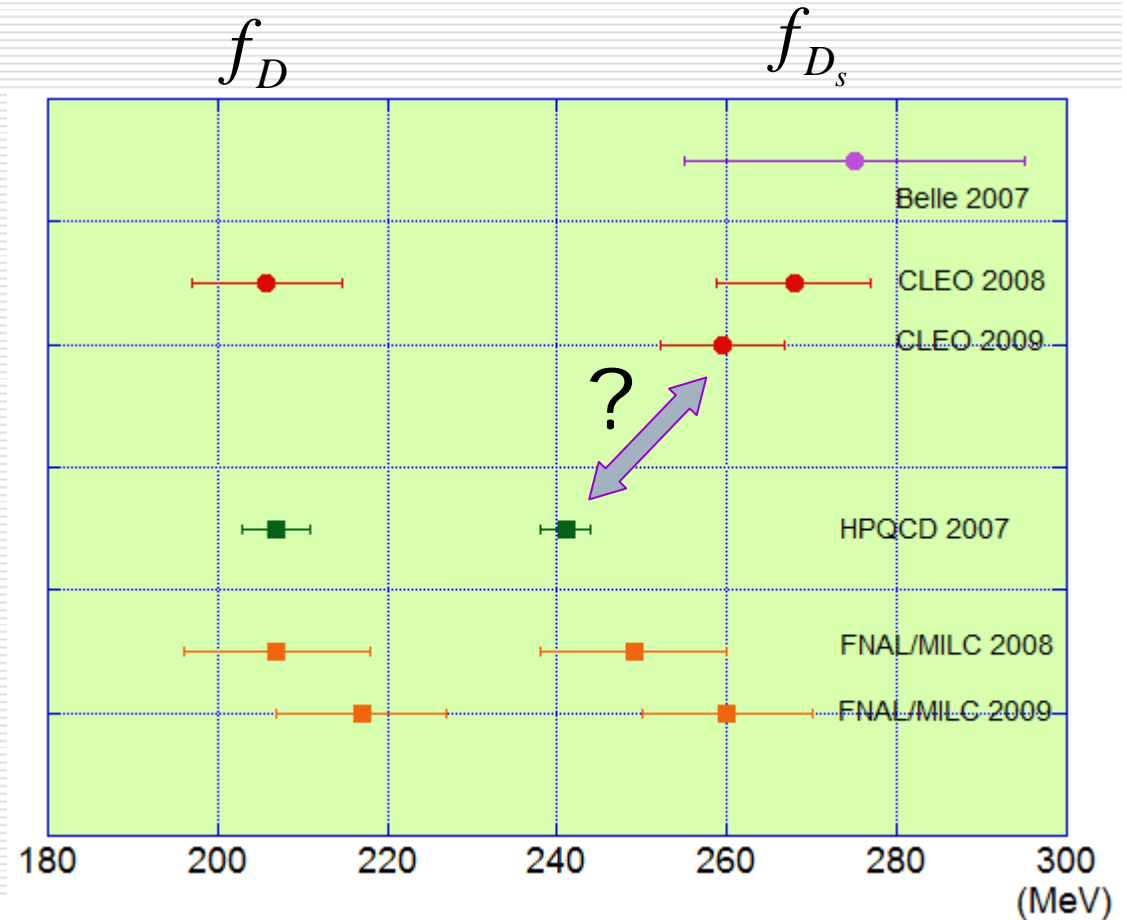
Lattice QCD

HPQCD

no change since the 2007 value

FNAL/MILC

moved up, but large error



Lattice QCD has to resolve the systematic uncertainty:

- HISQ action for HPQCD
- Clover action for FNAL/MILC