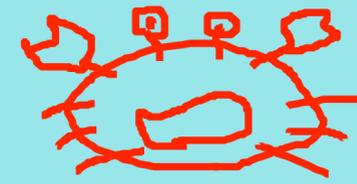


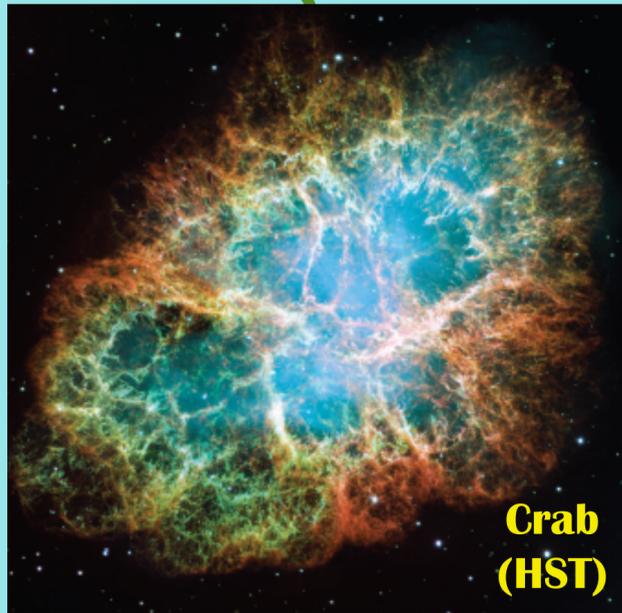
パルサー星雲からの
高エネルギー放射の
一次元モデル

田中 周太

Introduction to



PWN (Pulsar wind nebula)



Crab
(HST)

Line emission
(SN ejecta)



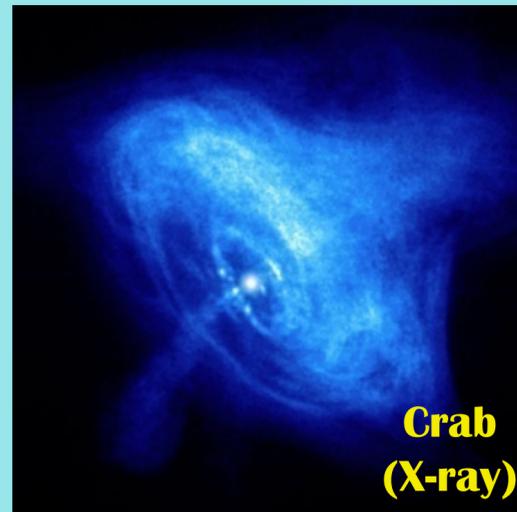
Crab
(radio - X-ray)

Non-thermal
(Pulsar wind)



G21.5-0.9
(X-ray)

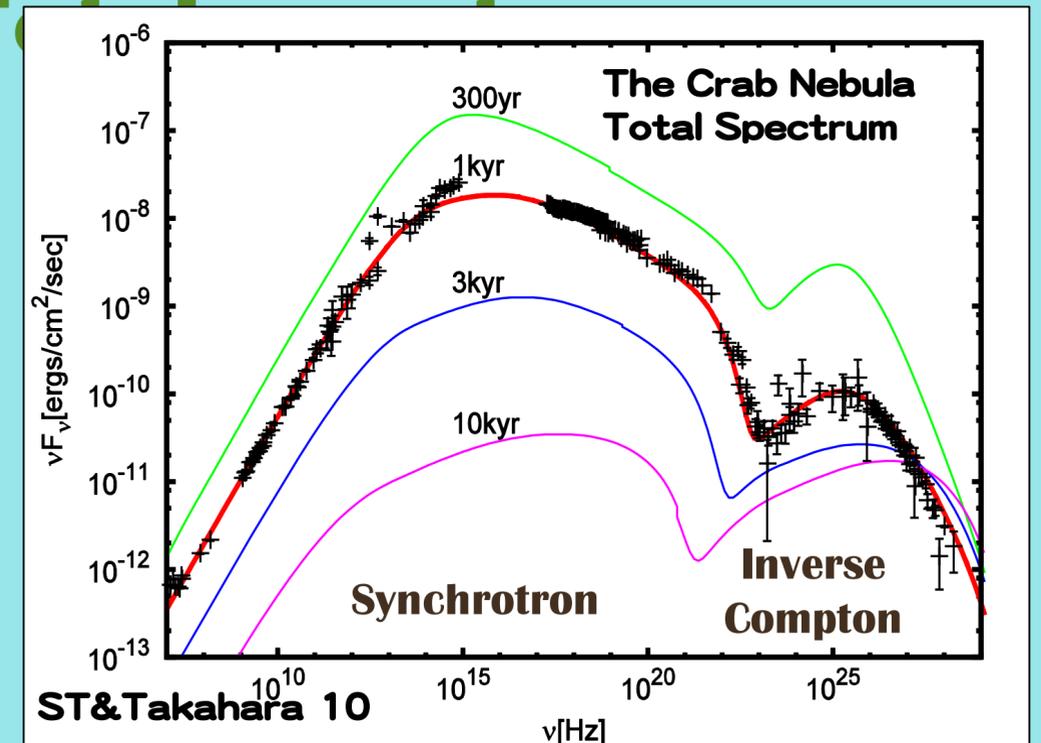
Confined by SNR



Crab
(X-ray)

Powered by PSR

Total Spectrum



We reproduced

- Total spectrum @ 1kyr
- Flux decrease rate $\sim 0.2\%$ /year @ radio

We obtain

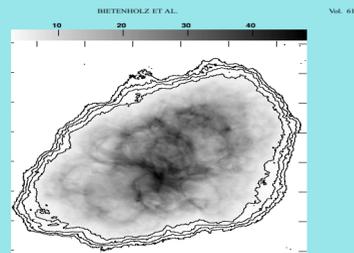
- mean B-field inside PWN.
- particle energy & number inside PWN.
- magnetization at injection.
- particle injection spectrum (broken PL).
- spin evolution of the central PSR.

Virtue of

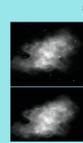
Bright and extended object

- **Brightness map in different frequencies of the Crab.**

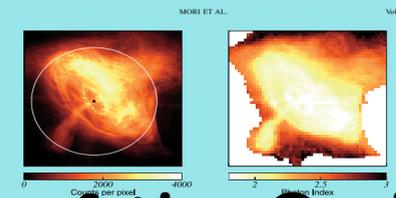
**Bietenholz+04ApJ
(radio)**



**Temim+06ApJ
(IR)**

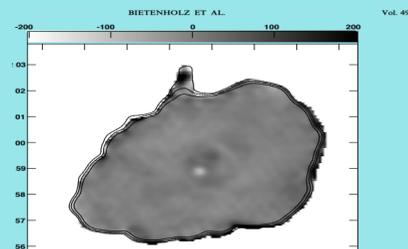


Mori+04ApJ (X-ray)

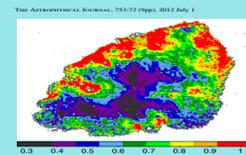


- **Spectral index map in different frequencies of the Crab.**

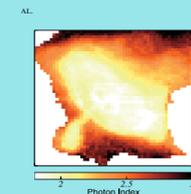
**Bietenholz+97ApJ
(radio)**



**Temim+12ApJ
(IR)**



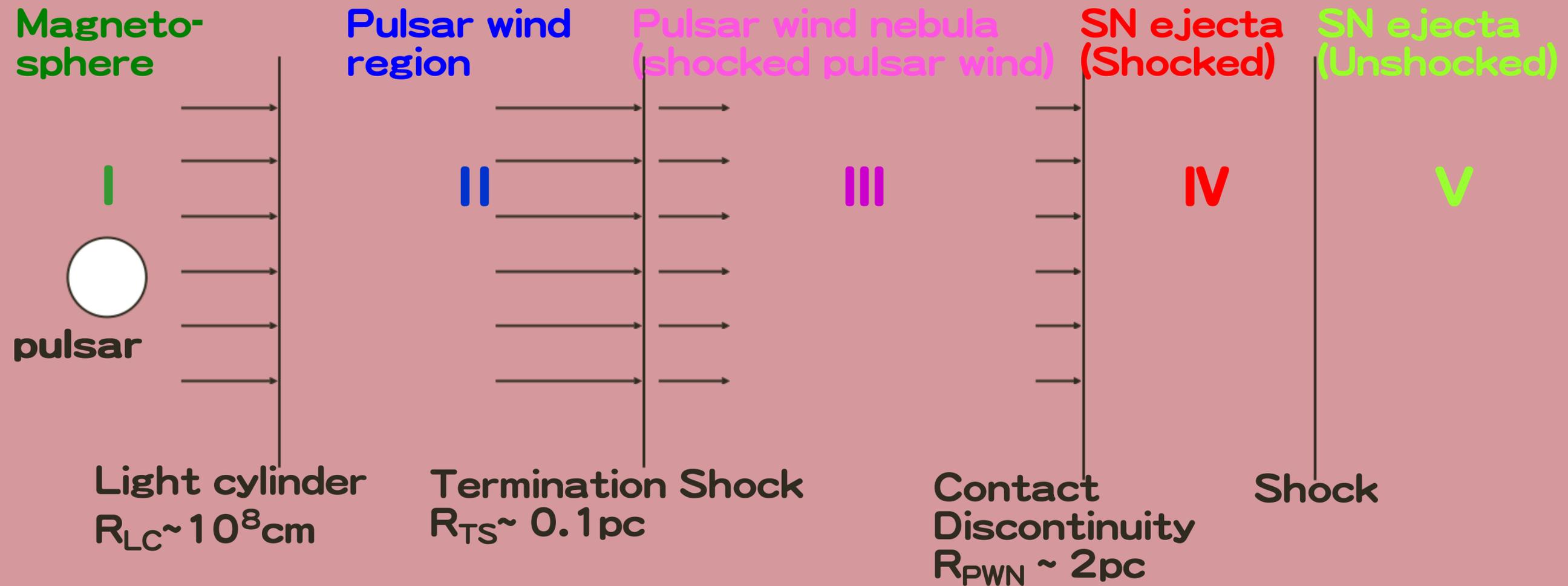
Mori+04ApJ (X-ray)



PWN emission reflects B-field & particle distribution of PWN

2-1: Qualitative Description

PWN is powered by central pulsar



**Young PWNe are almost spherical.
(Whole particles are confined within SNR)**

Motivation

What do we learn from observations of the Crab Nebula?

PSR & PWN properties

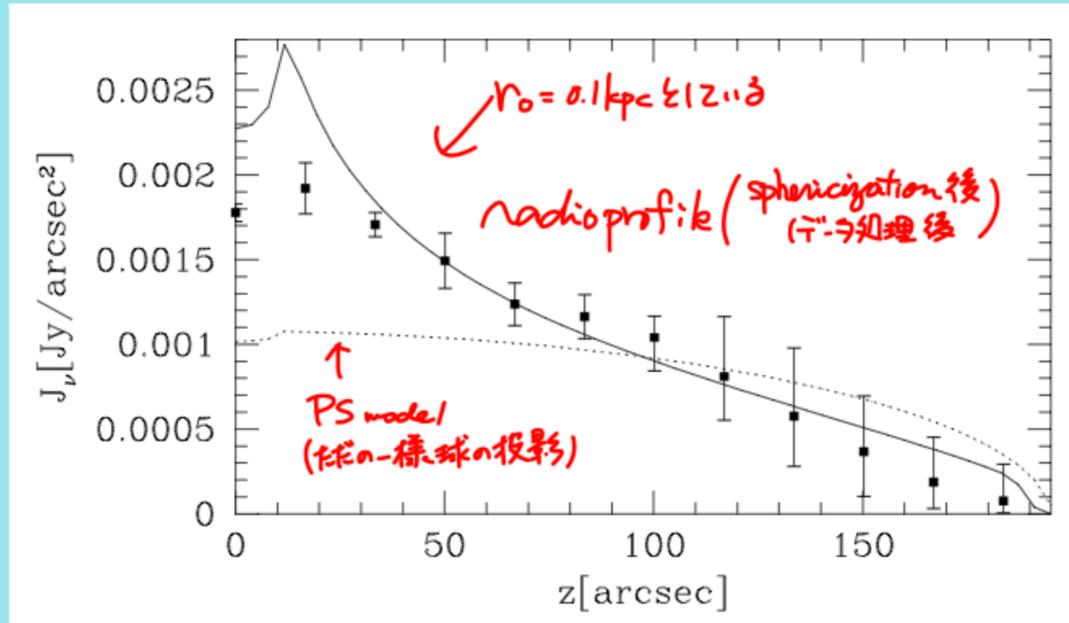
- Magnetization of pulsar wind. \rightarrow σ -problem
- Injection spectrum of particles. \rightarrow κ -problem & particle acceleration
- e^\pm escaping process from PWN. \rightarrow PAMELA anomaly

Physics

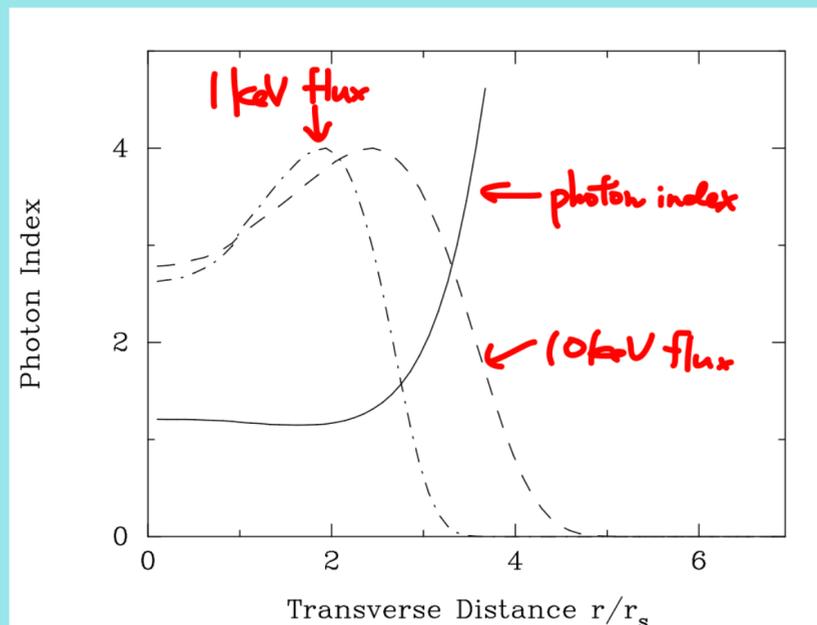
- Particle transport mechanism inside PWN. \rightarrow advection? diffusion?
- Particle acceleration process inside PWN. \rightarrow second-order acceleration?
- Expansion of PWN. \rightarrow Interaction between PWN & SNR at outer boundary.

Past Studies

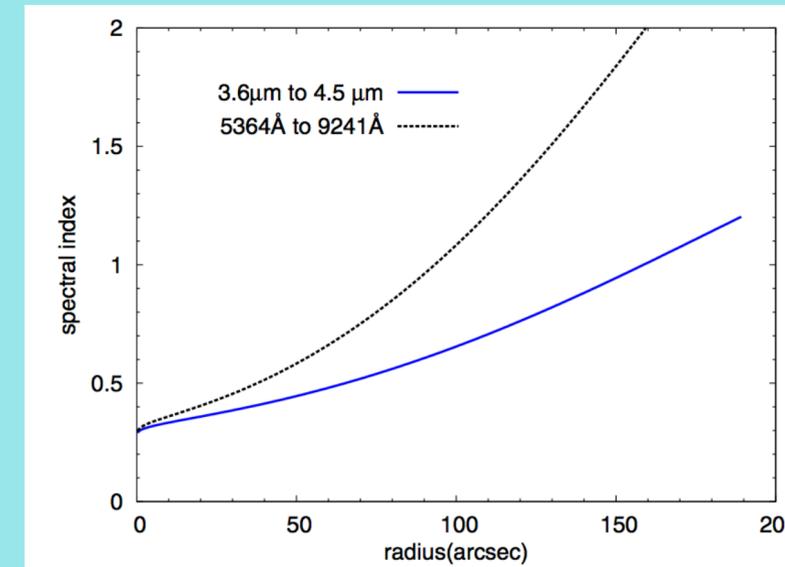
Amato+00, advection only Radio brightness map



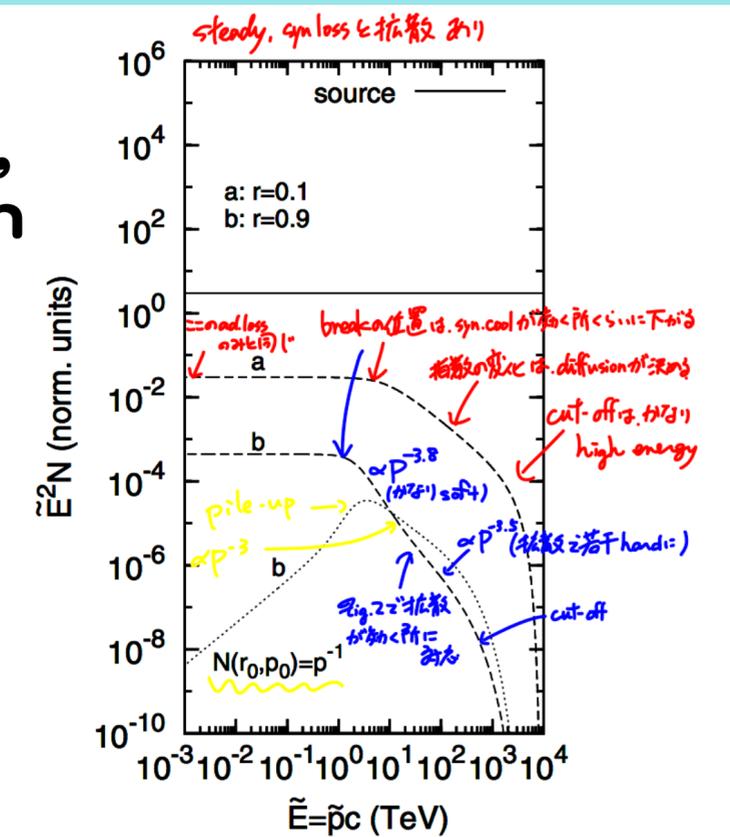
Reynolds03, advection only X-ray spectral index map



Tang & Chevalier 12, diffusion optical spectral index map



Vorstar & Moraal 13, diffusion & advection particle spectrum



Model

Transport Equation

$$\frac{\partial f}{\partial t} = -\nabla \cdot [(\mathbf{v} - K\nabla)f] + \frac{1}{p^2} \frac{\partial}{\partial p} \left[\left(\frac{1}{3} \nabla \cdot \mathbf{v} + (\beta_{\text{syn}} + \beta_{\text{IC}}) p^3 f \right) p^3 f \right] + Q_{\text{inj}}$$

Radial flow & toroidal B-field + induction equation

$$\mathbf{v}(r) = v_0 \left(\frac{r}{r_0} \right)^{-\alpha_v} \mathbf{e}_r, \quad B(r, t_{\text{inj}}) = B(t_{\text{inj}}) \left(\frac{r}{r_0} \right)^{-\alpha_B} \mathbf{e}_\varphi,$$

$$\frac{\partial B}{\partial t} = \nabla \times (\mathbf{v} \times B)$$

Bohm-type diffusion & particle injection

$$K = \frac{1}{3} \xi c r_{\text{gyro}} \quad Q_{\text{inj}}(r, t, p) = q(t) \frac{\delta(r - r_0)}{4\pi r_0^2} \times \begin{cases} (p/p_b)^{2-p_1} & \text{for } p_{\text{min}} \leq p \leq p_b, \\ (p/p_b)^{2-p_2} & \text{for } p_b \leq p \leq p_{\text{max}}, \end{cases}$$

We take into account the evolution of the spin-down luminosity (injection) & the R_{PWN} (expansion).

Application to the Crab Nebula

We consider that the particles and the magnetic field are injected at the inner radius $r = r_0$ at a time $t = t_{inj}$.

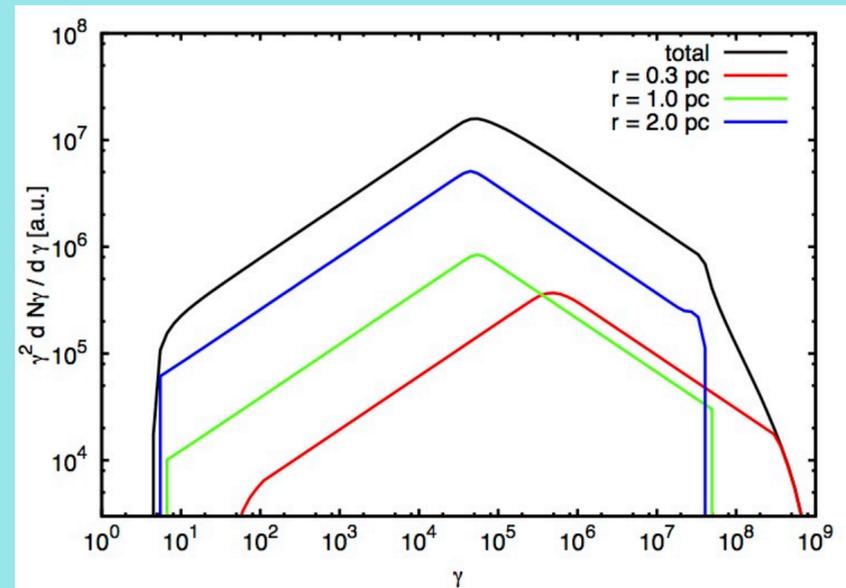
Main parameters are

1. gyro-factor ξ ,
2. inner radius r_0 or velocity profile index α_v ,
 $R_{PWN} = 2.0\text{pc}$, $V_{PWN} = 1.500\text{km/s}$ @ $t_{age} = 1\text{kyr}$
($r_0 = 0.1\text{pc}$ & $v_0 = c/3$ are too much.)
3. injected B-field or magnetization σ
 $B_0 = 300\mu\text{G}$
4. broken power-law injection with
 $p_1 = 1.5$, $p_2 = 2.5$, $\gamma_{min} = 10^2$, $\gamma = 10^6$, $\gamma_{max} = 10^9$

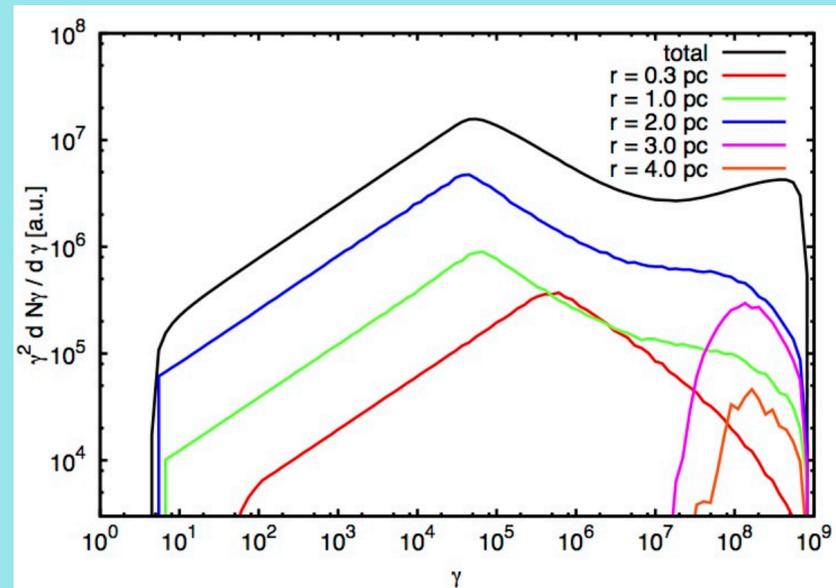
Results: Particle

$r_0 = 0.007 \text{ pc}$

without diffusion



with diffusion ($\xi = 1$)



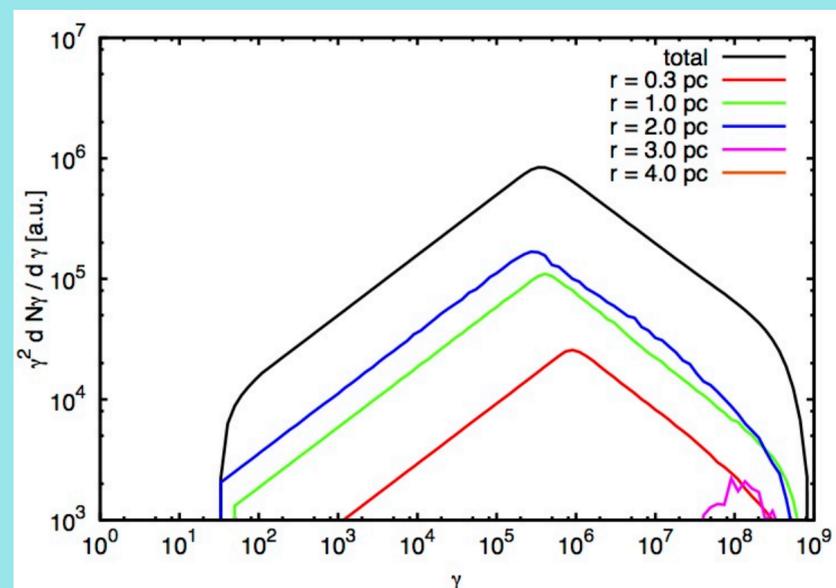
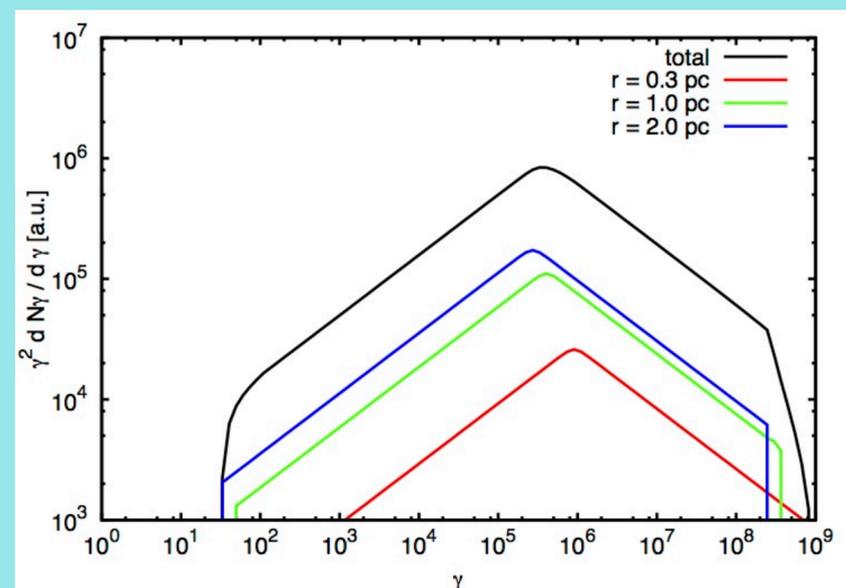
1. strong adiabatic cooling for $r_0 = 0.007 \text{ pc}$.

2. When ξ is non-zero, e^\pm can escape from PWN ($r \gg R_{\text{PWN}}$).

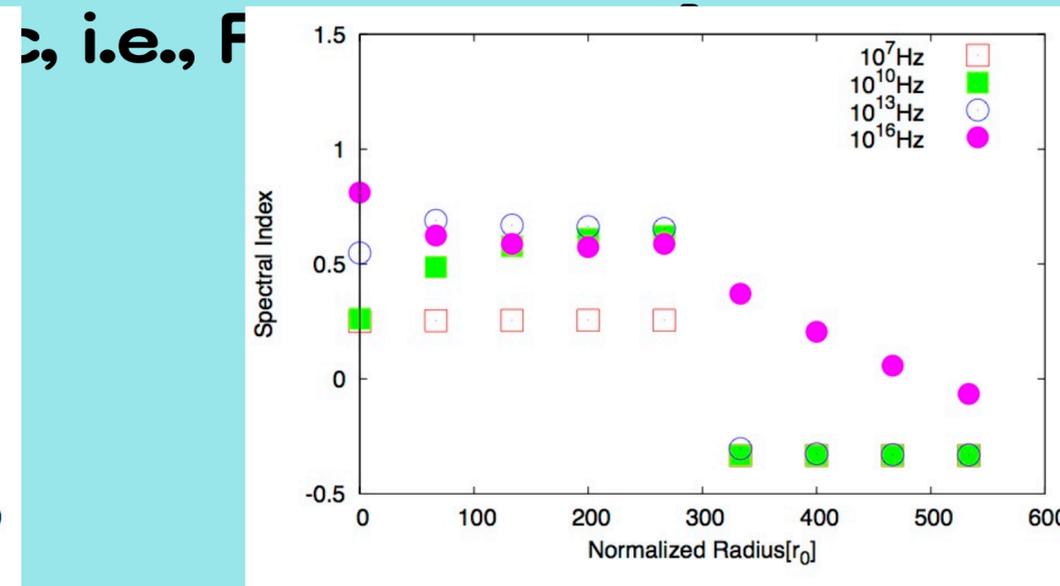
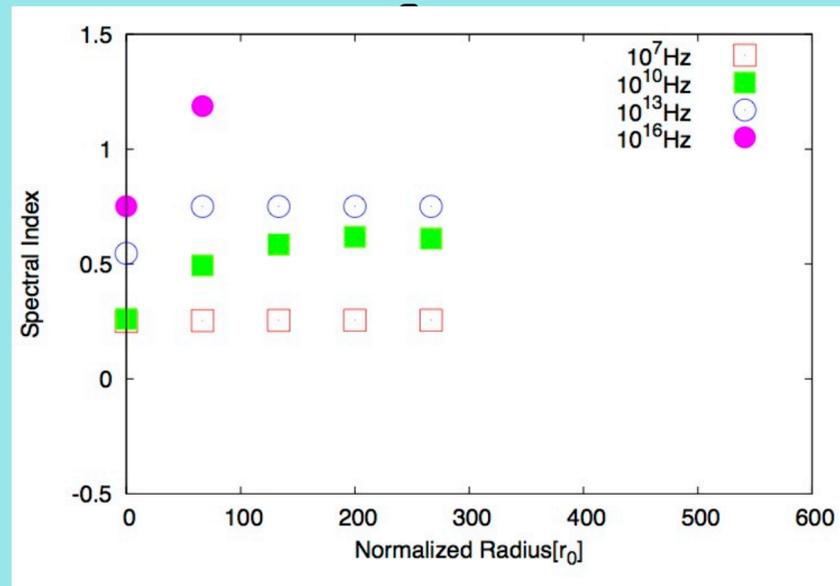
3. Escaped high energy e^\pm are cooled slowly.

4. For $r_0 = 0.007 \text{ pc}$, because r_0 is smaller than diffusion length for high energy e^\pm , more particles can escape than for $r_0 = 0.2 \text{ pc}$, mostly in an early phase of evolution.

r



Results: Spectral Index Profile



without diffusion
 $p_1 = 1.5$ & $p_2 = 2.5$ corresponds $\alpha = 0.25$ & 0.75

1. Without diffusion, we find synchrotron cooling effect increases spectral index at $10^{16}\text{Hz} > 10$ at $100r_0 \sim 0.7\text{pc}$.
2. Diffusion effect smooths spectral index profile and we do not find synchrotron cooling hardening of the spectrum for $E_0 = 100 \mu\text{G}$. We need to select appropriate parameter set of (ξ, B_0) to fit the spectral index distribution.
3. $\alpha = -1/3$ corresponds to low freq. side of synchrotron spectra and is appeared at $r > R_{\text{PWN}} @ < 10^{16}\text{Hz}$.

Conclusion

- Convection-Diffusion transport equation is considered to study the broadband emission structure of PWNe.
- For the escaping of e^{\pm} from PWNe, we should compare diffusion length & r_0 .
- γ -ray distribution almost traces the particle distribution for inverse Compton scattering off CMB.
- Synchrotron Self-Compton process is considered to be dominated in the Crab Nebula in γ -rays and should be considered the future work.
- We need larger value of B-field at injection point, to obtain softer spectra observed for many PWNe in X-rays ($\alpha > 2$).