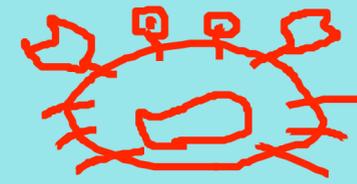


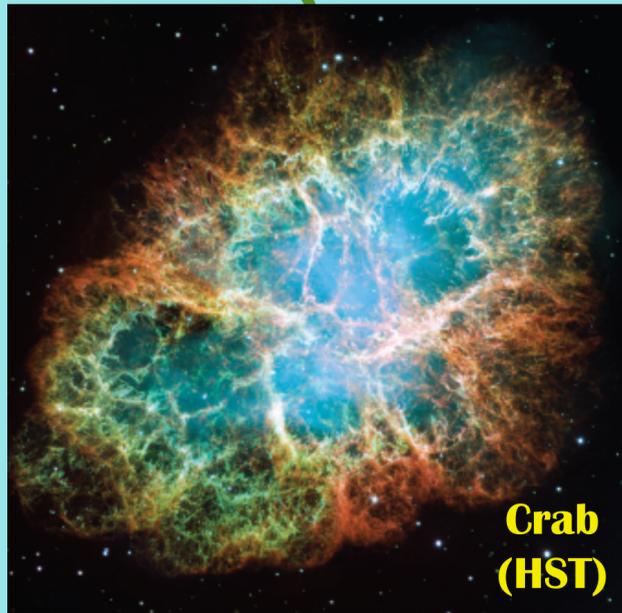
パルサー星雲からの
高エネルギー放射の
一次元モデル

田中 周太

Introduction to



PWN (Pulsar wind nebula)



Crab
(HST)

Line emission
(SN ejecta)



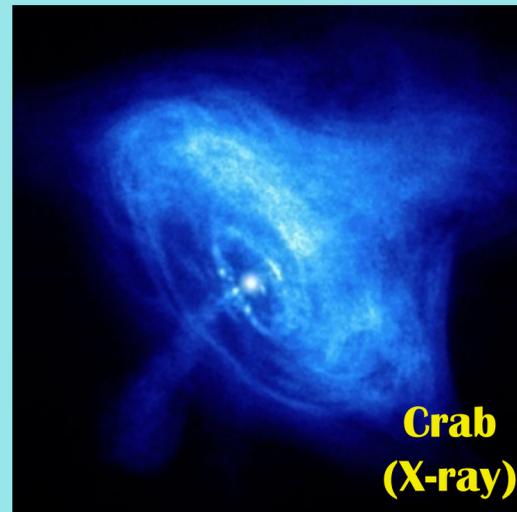
Crab
(radio - X-ray)

Non-thermal
(Pulsar wind)



G21.5-0.9
(X-ray)

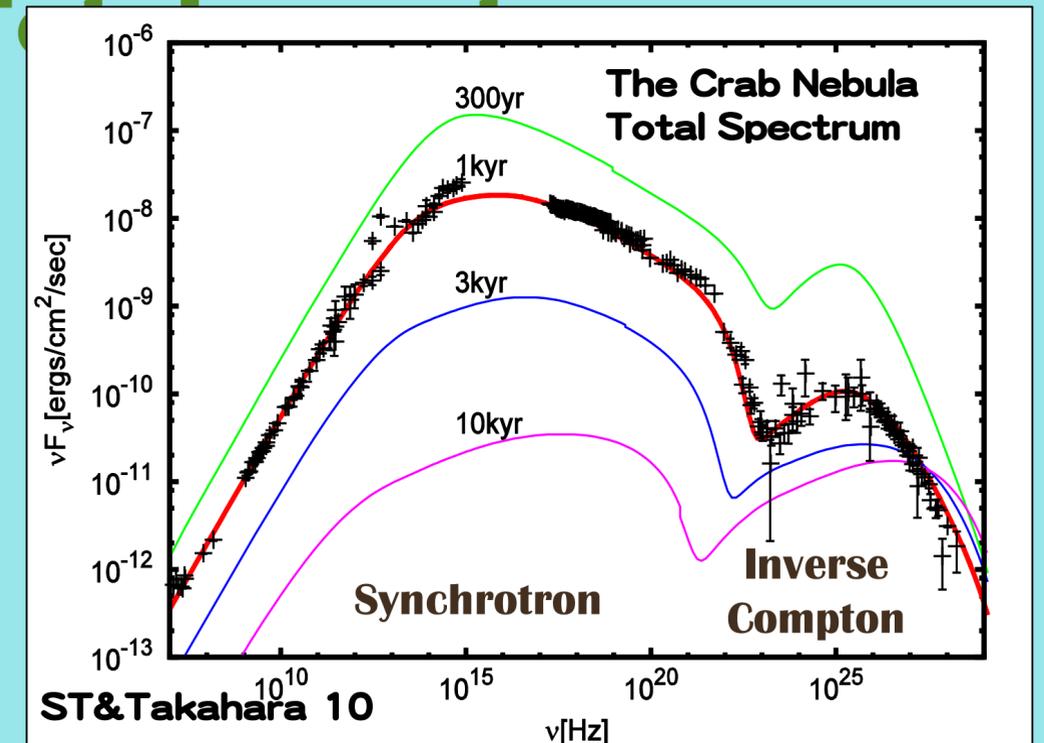
Confined by SNR



Crab
(X-ray)

Powered by PSR

Total Spectrum



We reproduced

- Total spectrum @ 1kyr
- Flux decrease rate $\sim 0.2\%$ /year @ radio

We obtain

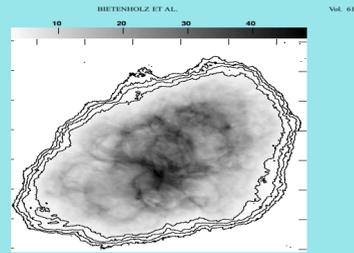
- mean B-field inside PWN.
- particle energy & number inside PWN.
- magnetization at injection.
- particle injection spectrum (broken PL).
- spin evolution of the central PSR.

Virtue of

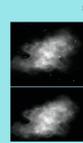
Bright and extended object

- **Brightness map in different frequencies of the Crab.**

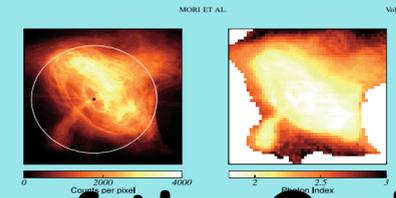
**Bietenholz+04ApJ
(radio)**



**Temim+06ApJ
(IR)**

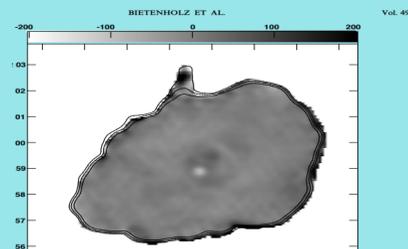


Mori+04ApJ (X-ray)

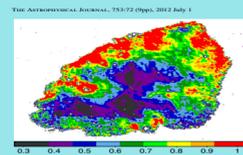


- **Spectral index map in different frequencies of the Crab.**

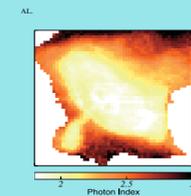
**Bietenholz+97ApJ
(radio)**



**Temim+12ApJ
(IR)**



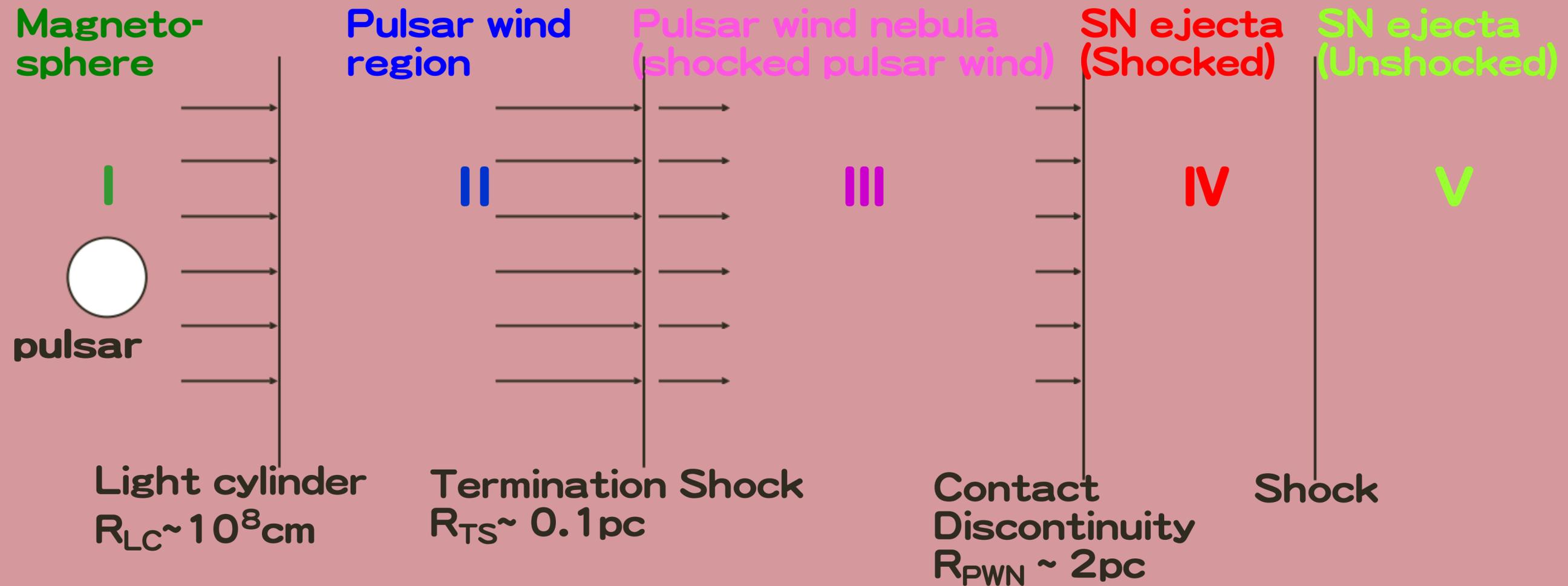
Mori+04ApJ (X-ray)



PWN emission reflects B-field & particle distribution of PWN

2-1: Qualitative Description

PWN is powered by central pulsar



Young PWNe are almost spherical.
(Whole particles are confined within SNR)

Motivation

What do we learn from observations of the Crab Nebula?

PSR & PWN properties

- Magnetization of pulsar wind. \rightarrow σ -problem
- Injection spectrum of particles. \rightarrow κ -problem & particle acceleration
- e^\pm escaping process from PWN. \rightarrow PAMELA anomaly

Physics

- Particle transport mechanism inside PWN. \rightarrow advection? diffusion?
- Particle acceleration process inside PWN. \rightarrow second-order acceleration?
- Expansion of PWN. \rightarrow Interaction between PWN & SNR at outer boundary.

Model

Transport Equation

$$\frac{\partial f}{\partial t} = -\nabla \cdot [(\mathbf{v} - K\nabla)f] + \frac{1}{p^2} \frac{\partial}{\partial p} \left[\left(\frac{1}{3} \nabla \cdot \mathbf{v} + (\beta_{\text{syn}} + \beta_{\text{IC}}) p^3 f \right) p^3 \right] + Q_{\text{inj}}$$

Radial flow & toroidal B-field + induction equation

$$\mathbf{v}(r) = v_0 \left(\frac{r}{r_0} \right)^{-\alpha_v} \mathbf{e}_r, \quad B(r, t_{\text{inj}}) = B(t_{\text{inj}}) \left(\frac{r}{r_0} \right)^{-\alpha_B} \mathbf{e}_\varphi,$$

$$\frac{\partial B}{\partial t} = \nabla \times (\mathbf{v} \times B)$$

Bohm-type diffusion & particle injection

$$K = \frac{1}{3} \xi c r_{\text{gyro}} \quad Q_{\text{inj}}(r, t, p) = q(t) \frac{\delta(r - r_0)}{4\pi r_0^2} \times \begin{cases} (p/p_b)^{2-p_1} & \text{for } p_{\text{min}} \leq p \leq p_b, \\ (p/p_b)^{2-p_2} & \text{for } p_b \leq p \leq p_{\text{max}}, \end{cases}$$

We take into account the evolution of the spin-down luminosity (injection) & the R_{PWN} (expansion).

Application to the Crab Nebula

We consider that the particles and the magnetic field are injected at the inner radius $r = r_0$ at a time $t = t_{inj}$.

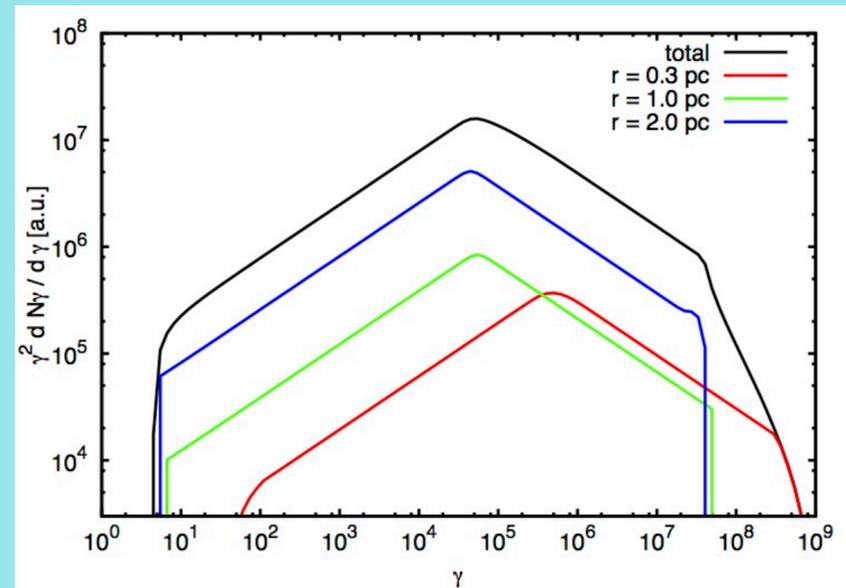
Main parameters are

1. gyro-factor ξ ,
2. inner radius r_0 or velocity profile index α_v ,
 $R_{PWN} = 2.0\text{pc}$, $V_{PWN} = 1.500\text{km/s}$ @ $t_{age} = 1\text{kyr}$
($r_0 = 0.1\text{pc}$ & $v_0 = c/3$ are too much.)
3. injected B-field or magnetization σ
 $B_0 = 300\mu\text{G}$
4. broken power-law injection with
 $p_1 = 1.5$, $p_2 = 2.5$, $\gamma_{min} = 10^2$, $\gamma = 10^6$, $\gamma_{max} = 10^9$

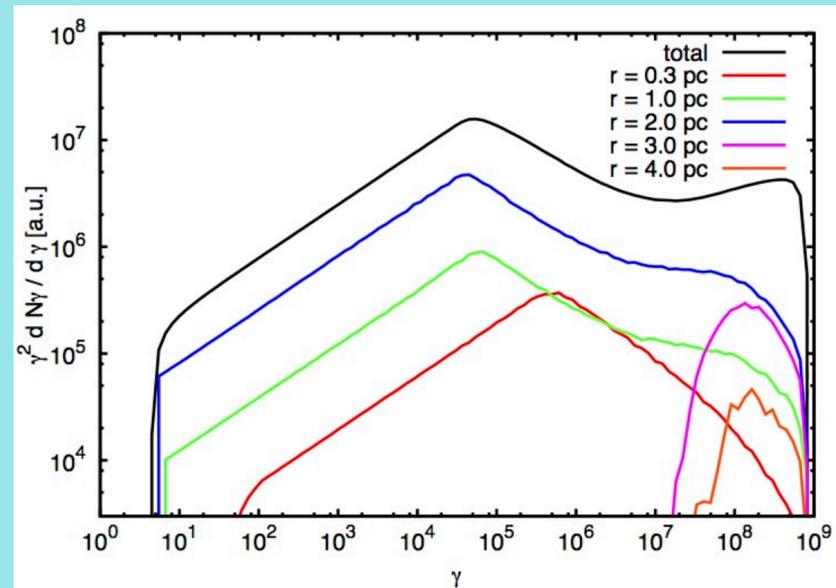
Results: Particle

$r_0 = 0.007 \text{ pc}$

without diffusion



with diffusion ($\xi = 1$)



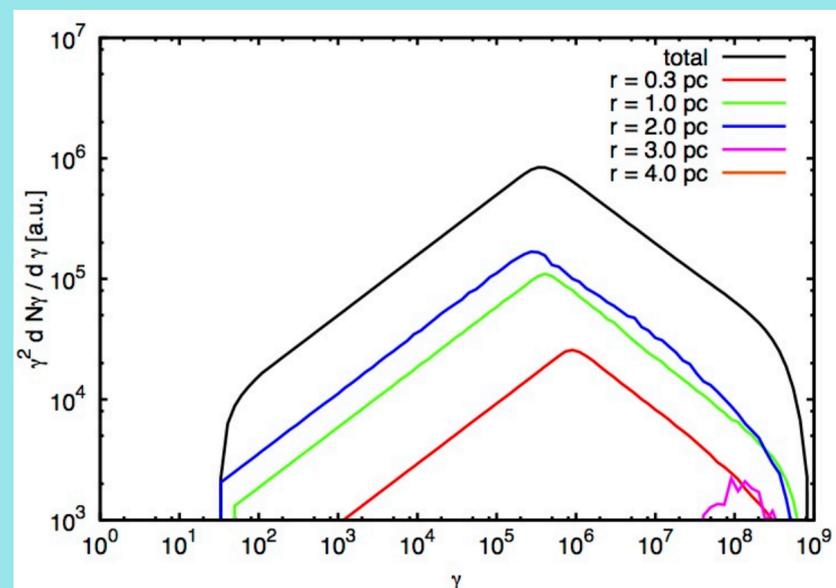
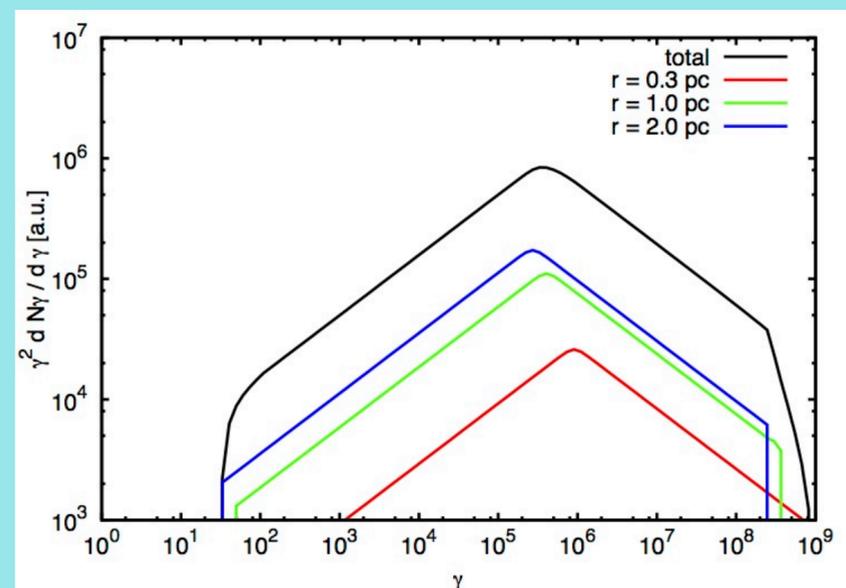
1. strong adiabatic cooling for $r_0 = 0.007 \text{ pc}$.

2. When ξ is non-zero, e^\pm can escape from PWN ($r \gg R_{\text{PWN}}$).

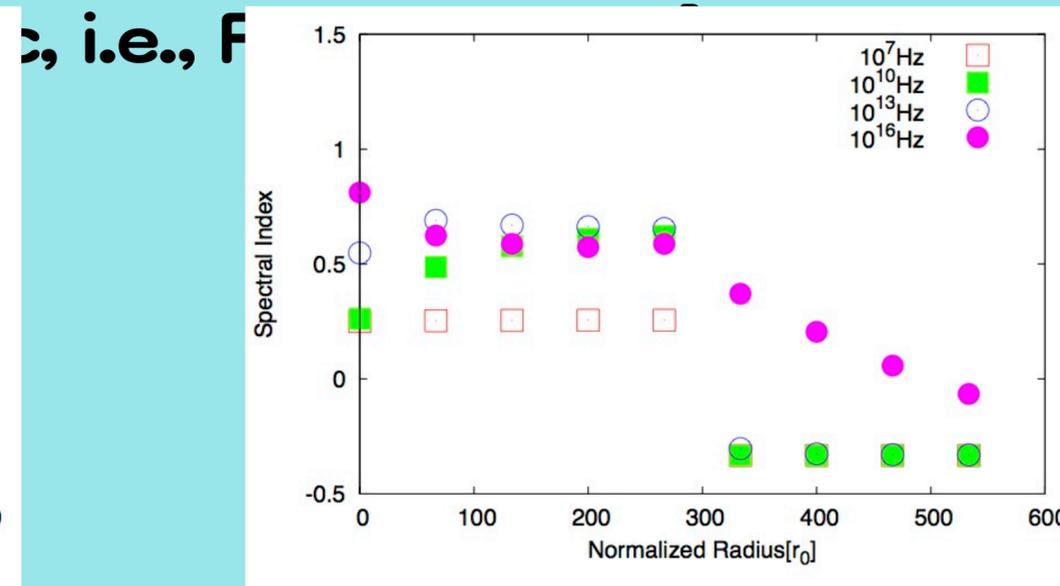
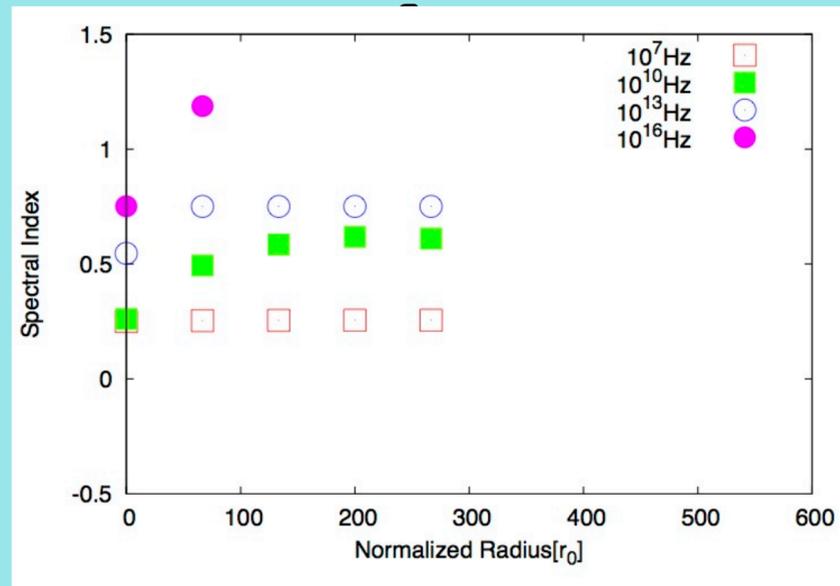
3. Escaped high energy e^\pm are cooled slowly.

4. For $r_0 = 0.007 \text{ pc}$, because r_0 is smaller than diffusion length for high energy e^\pm , more particles can escape than for $r_0 = 0.2 \text{ pc}$, mostly in an early phase of evolution.

r



Results: Spectral Index Profile



without diffusion with diffusion ($\xi=1$)
 $p_1 = 1.5$ & $p_2 = 2.5$ corresponds $\alpha = 0.25$ & 0.75

1. Without diffusion, we find synchrotron cooling effect increases spectral index at $10^{16}\text{Hz} > 10$ at $100r_0 \sim 0.7\text{pc}$.
2. Diffusion effect smooths spectral index profile and we do not find synchrotron cooling hardening of the spectrum for $E_0 = 100 \mu\text{G}$. We need to select appropriate parameter set of (ξ, B_0) to fit the spectral index distribution.
3. $\alpha = -1/3$ corresponds to low freq. side of synchrotron spectra and is appeared at $r > R_{\text{PWN}} @ < 10^{16}\text{Hz}$.

Conclusion

- Convection-Diffusion transport equation is considered to study the broadband emission structure of PWNe.
- For the escaping of e^{\pm} from PWNe, we should compare diffusion length & r_0 .
- γ -ray distribution almost traces the particle distribution for inverse Compton scattering off CMB.
- Synchrotron Self-Compton process is considered to be dominated in the Crab Nebula in γ -rays and should be considered the future work.
- We need larger value of B-field at injection point, to obtain softer spectra observed for many PWNe in X-rays ($\alpha > 2$).