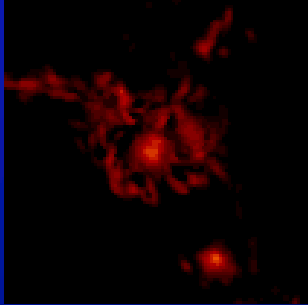


# Investigating Galaxy Clusters through $\gamma$ -rays



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Sept. 27 2002, Kashiwa

## Outline

- **Observational evidence for CRs in the ICM**
- **Properties of IGM shocks**
- **Novel numerical scheme to follow CRs evolution**
- **CR ionic pressure component in ICM**
- **Numerical model of  $\gamma$ -ray emission from GCs**
- **Contribution of cosmological CRs to the CGB**
- **Conclusions**

## Observations

Extended radio emission (eg Giovannini & Feretti 2000)

Halos: un-polarized, x-ray like morphology

Relics: polarized, peripheral

$$L_{1.4\text{GHz}} \propto T^\alpha \quad \text{with} \quad 4 \leq \alpha \leq 5 \quad \text{Liang et al. (2000)}$$

Radiation in excess to what expected from thermal emission of the intra-cluster medium:

Hard X-rays (Fusco-Femiano et al. 1999, Kaastra et al. 1999)

Extreme-UV (Lieu et al. 1996a,b; Bowyer et al. 1996,1999)

## Properties of cosmic shocks



Nbody+Hydro TVD cosmological code (Ryu et al. 1993)

Simulate LSS formation: gas + dark matter + B-field

+

COSMOCR: Cosmic Ray Code for Cosmological Applications (Miniati 2001, CPC 141, 17)

• The code allows an *explicit treatment* of the CR component.

injection at shocks

acceleration + losses

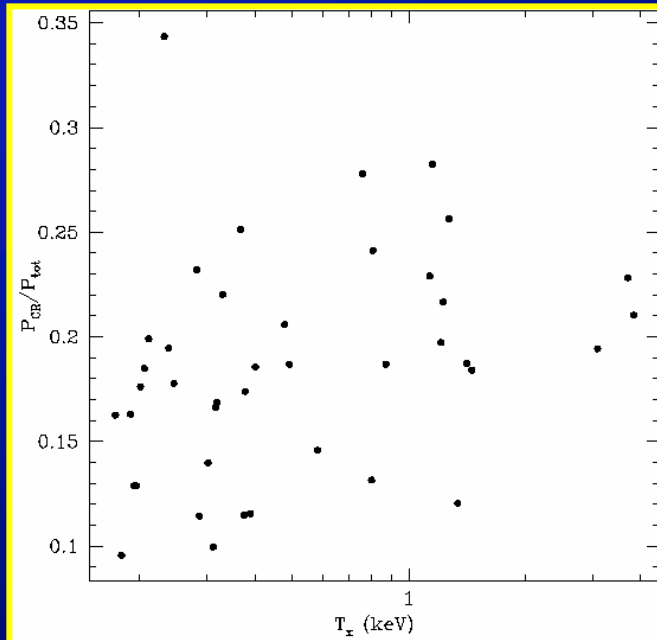
transport

⇒ spatial & spectral CR distribution

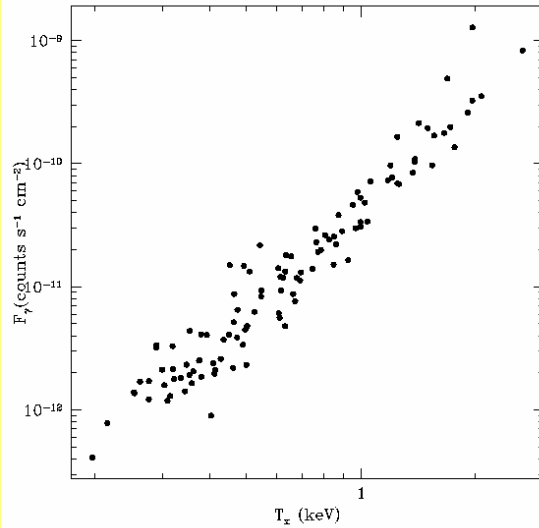
• Includes: ions, primary  $e^-$  and secondary  $e^\pm$

Ratio of CR  
to thermal  
pressure  
vs  
cluster  
temperature

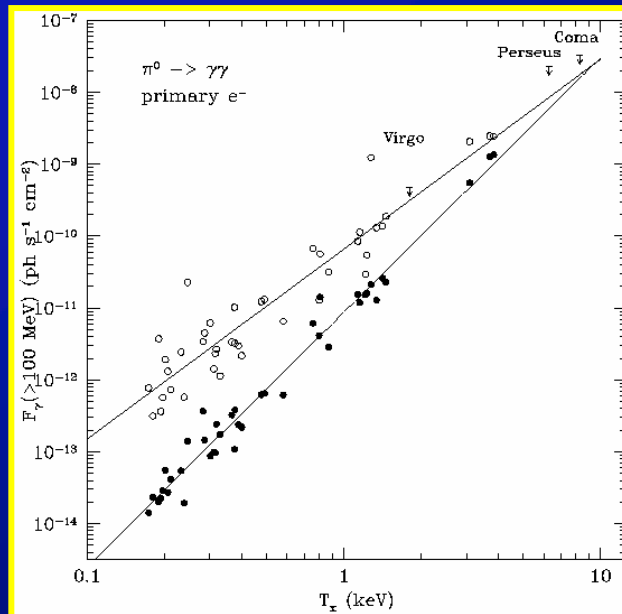
$\Lambda$ CDM  
 $N=512^3$   
 $\sigma_8=0.9$   
 $\Omega_\Lambda=0.7$   
 $\Omega_b=0.04$   
 $n_s=1$

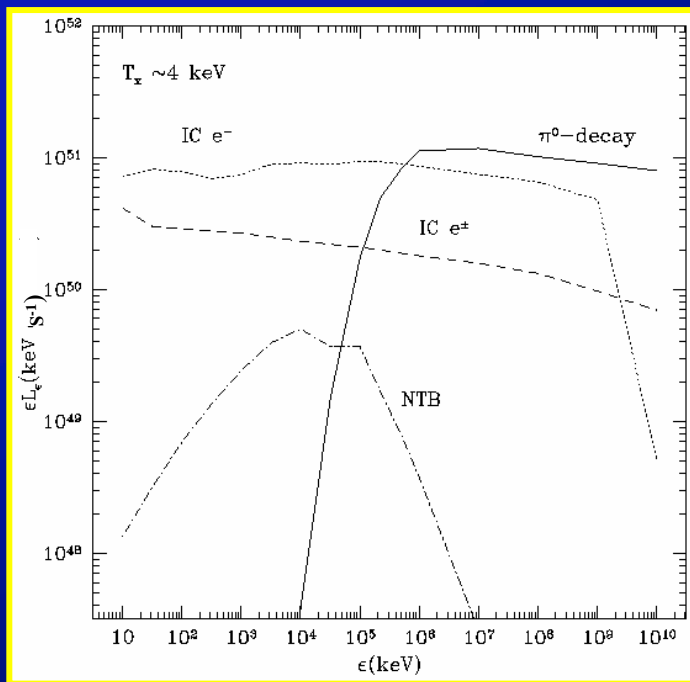


$\gamma$ -ray flux from  
 $\pi^0 \rightarrow 2\gamma$  decay  
 vs  
 cluster  
 temperature



$$F_\gamma(>100\text{MeV}) = 7.4 \cdot 10^{-9} \times \left( \frac{T_x}{6.7\text{keV}} \right)^3 \text{ counts } s^{-1} \text{ cm}^{-2}$$

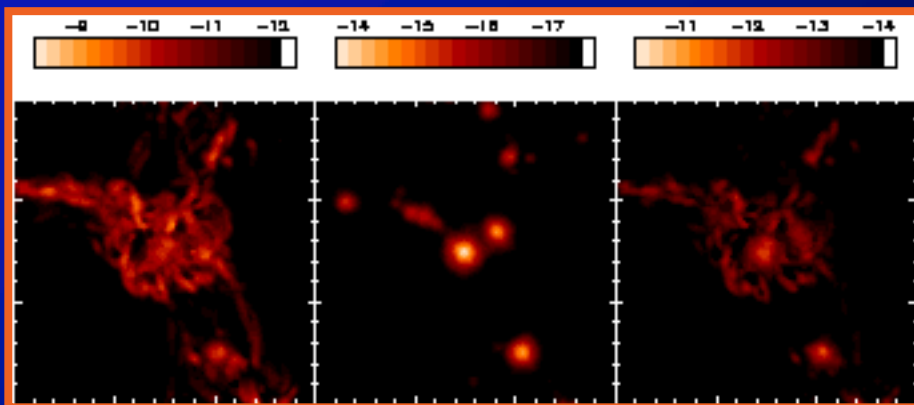


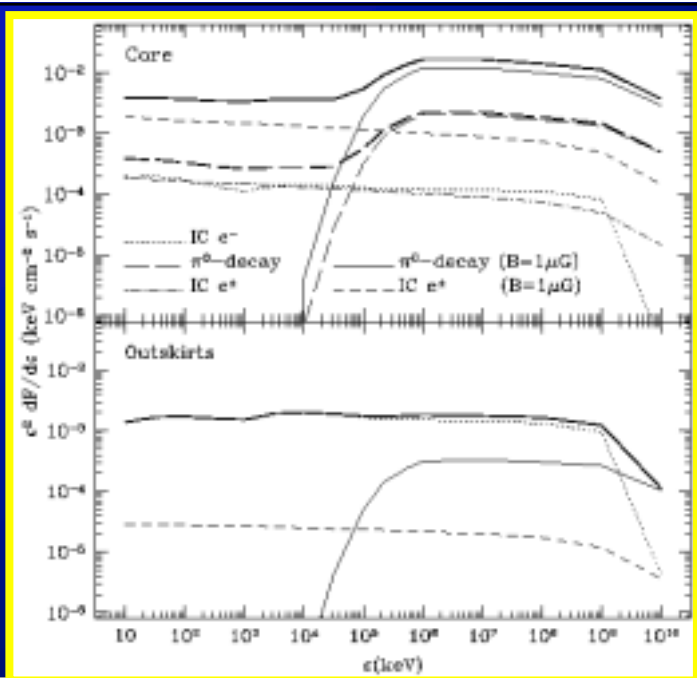


Hard X-ray  
 $F(>100 \text{ keV})$

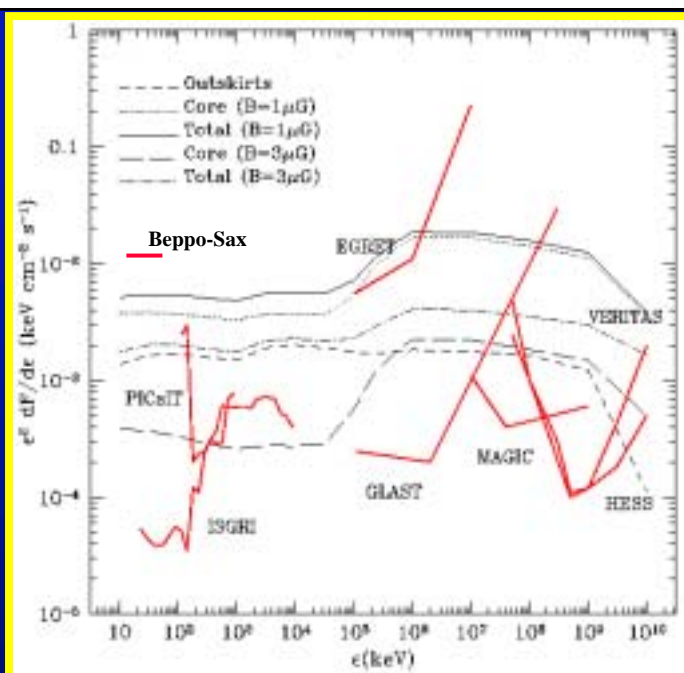
Thermal X-ray

$\gamma$ -ray  
 $F(>100 \text{ MeV})$





Miniati 2002, MN subm.

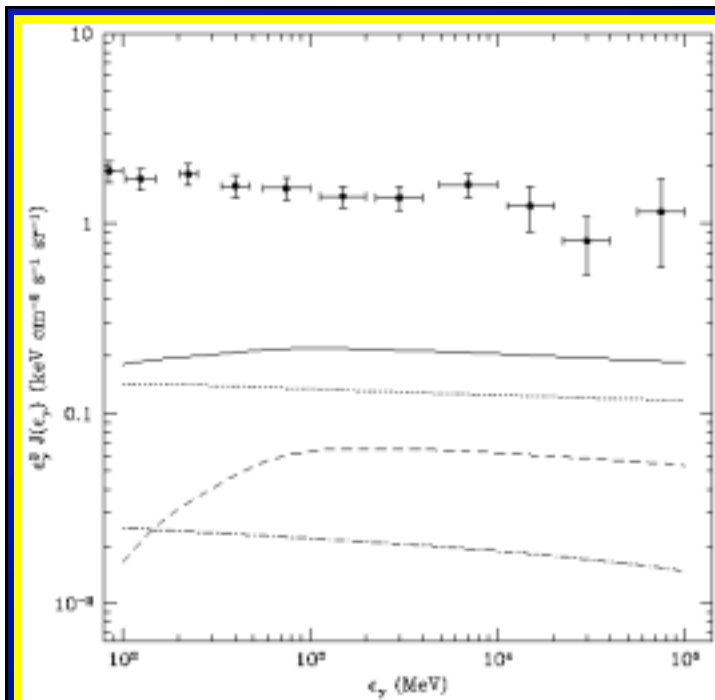


Miniati 2002, MN subm.

## Cosmic $\gamma$ -ray Background

$$\varepsilon^2 J(\varepsilon) \approx 5.4 (\varepsilon/\text{keV})^{-0.1} \text{ keV cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1} \text{ (Sreekumar et al. 1998)}$$

- Identified EGRET balzars ..... 7% (Mukherjee & Chiang 99)
- Undetected  $\gamma$ -ray loud blazars .. 100% (Padovani et al. 93, Steker & Salomon 96)
- Undetected  $\gamma$ -ray loud blazars ... 25% (Chiang & Mukherjee 98)
- IGM shock accelerated CR  $e^-$  ... 75% (Loeb & Waxman 00)



Data  
(Sreekumar et al. 98)

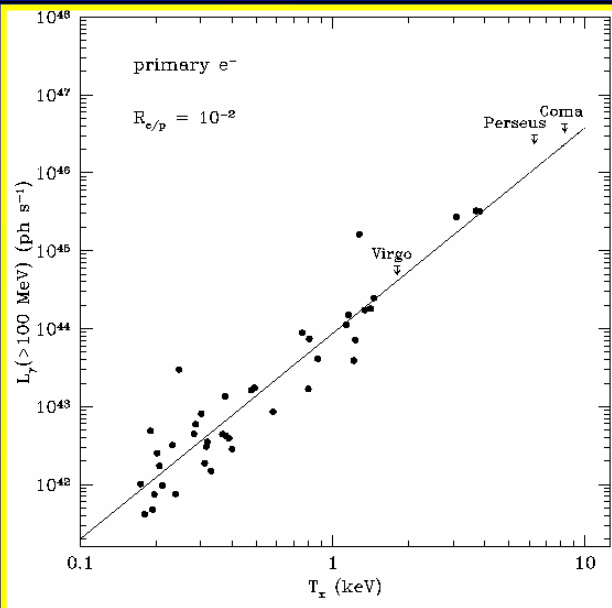
Total  
IC  $e^-$   
 $\pi^0$ -decay

IC  $e^\pm$

Miniati 2002, MN  
(astro-ph/0203014)

Miniati 2002, MN in press  
(astro-ph/0203014)

Upper limits from  
EGRET experiment  
(Reimer 1999;  
Sreekumar et al. 1996)



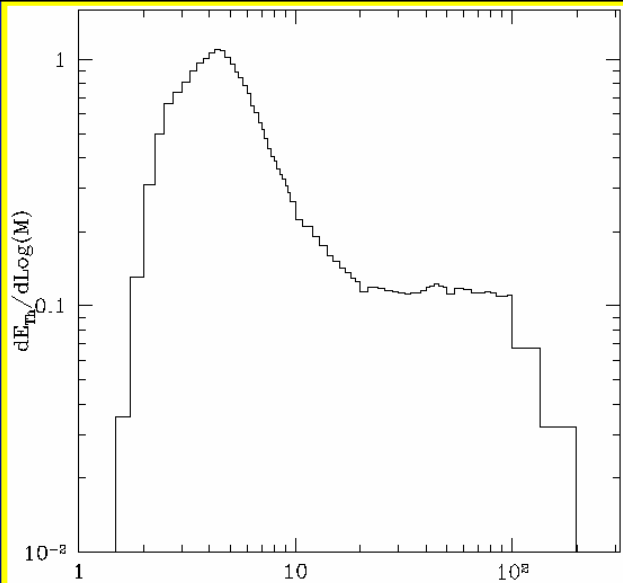
$$L_{\gamma}(>100\text{MeV}) = 8.7 \times 10^{43} \cdot \left(\frac{\eta}{4 \times 10^{-3}}\right) \cdot \left(\frac{T_x}{\text{keV}}\right)^{2.6} \text{ ph s}^{-1}$$

## Summary

- Most effective shocks have  $4 < M < 10$  or so
- CR ions might be dynamically important in the ICM
- HXR and  $\gamma$ -ray observations will help us better understand ICM physics but imaging capability is very important for a correct interpretation of the observational results.
- EGRET experiment implies:  $\eta_e \equiv \frac{P_{cre}}{\rho v_s^2} \leq 0.01$
- Cosmological CRs might contribute  $\sim 25\%$  of the CGB (70% from CR  $e^-$  and 30% from CR  $p + e^\pm$ )



$\Delta E_{th}(M)$   
 30%  $M < 4$   
 45%  $4 \leq M \leq 10$   
 25%  $M > 10$



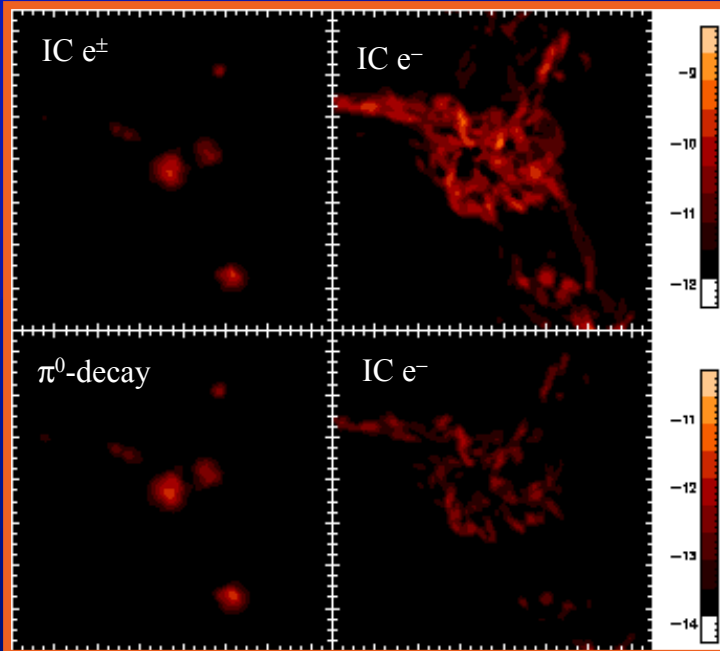
Miniati 2002, MN  
 (astro-ph/0203014)

$$\frac{\partial^2 \Delta E_{th}}{\partial \log M \partial \log T}; \quad \Delta E_{th} = k_B \int_{t_i}^{Mach(z=0)} dt \int \Delta T n v \cdot dS|_{shock}$$

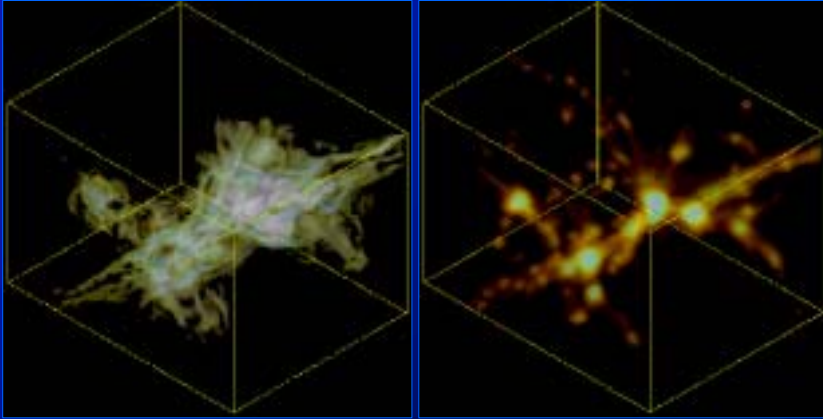
$[F] = \text{ph s}^{-1} \text{cm}^{-2}$

$F(>100 \text{ keV})$

$F(>100 \text{ MeV})$

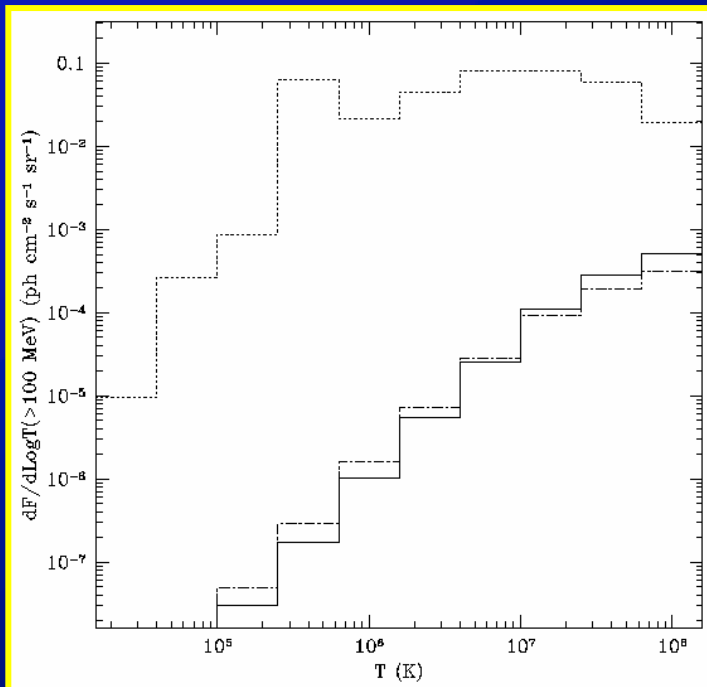


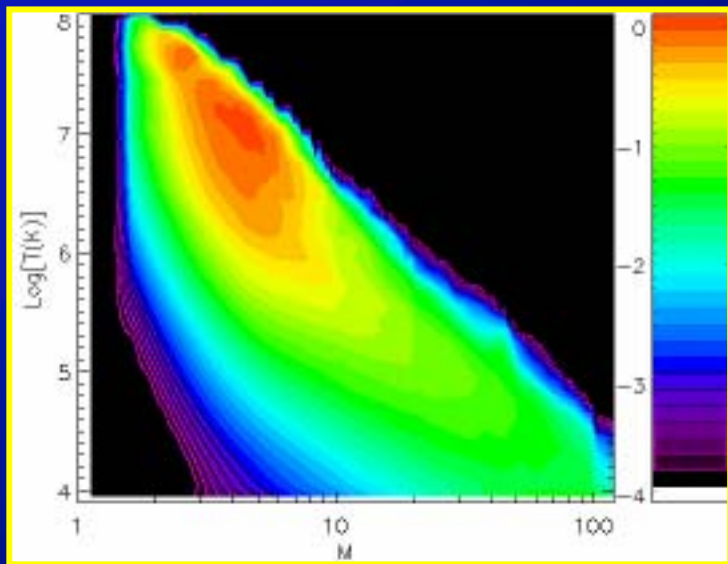
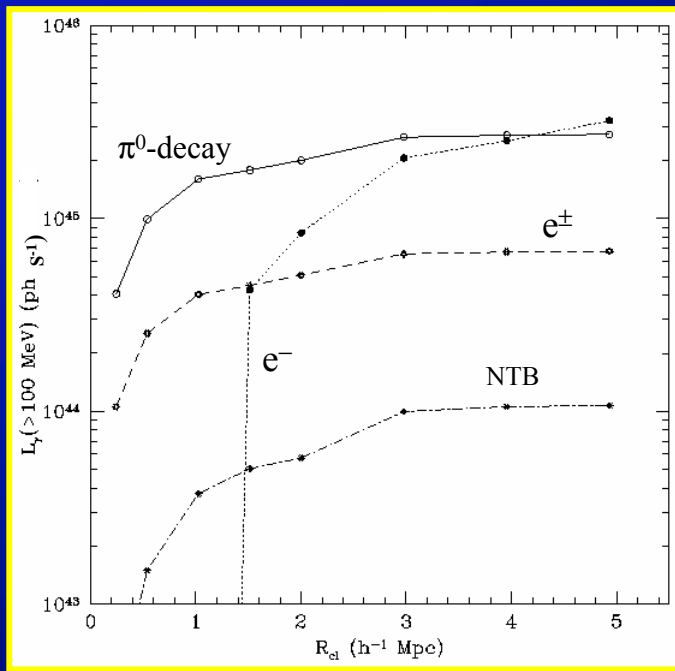
# Properties of cosmic shocks



SCDM,  $L=80$  Mpc,  $\sigma_8=1.05$ ,  $n=1$ ,  $N=270^3$

Miniati, Ryu, Jones, Kang, Cen & Ostriker (2000, ApJ 542, 608)

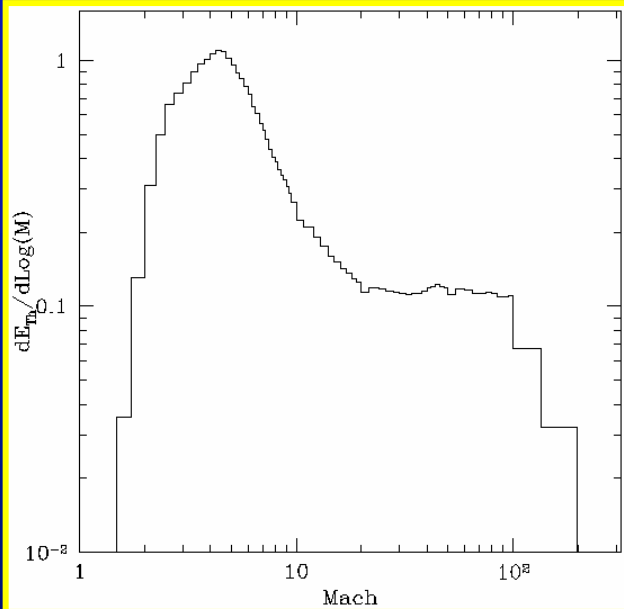




Miniati 2002, MN  
(astro-ph/0203014)

$$\frac{\partial^2 \Delta E_{\text{th}}}{\partial \log M \partial \log T}; \quad \Delta E_{\text{th}} = k_B \int_{t_i}^{t(z=0)} dt \int \Delta T n v \cdot dS|_{\text{shock}}$$

$\Delta E_{th}(M)$	
30%	$M < 4$
45%	$4 \leq M \leq 10$
25%	$M > 10$



$$\frac{\Delta T}{T} = \frac{[2\gamma_{gas} M^2 - (\gamma_{gas} - 1)][(\gamma_{gas} - 1)M^2 + 2]}{(\gamma_{gas} + 1)^2 M^2} - 1$$

## Injection

Based on the thermal leakage model of Kang & Jones (1995)

The post-shock gas is assumed to have nearly thermalized to a Maxwellian distribution with temperature  $T_{shock}$ . Then those particles with momentum above a threshold Value  $p_{inj}$  are injected in the acceleration mechanism.

$$p \geq p_{inj} \equiv c_1 2\sqrt{m_p k_B T_{shock}}$$

$$c_1 \approx 2.6$$

...injection

$$\eta \equiv \frac{n_{inj}}{n_2} \cong 8\sqrt{2/\pi} c_1^3 e^{-2c_1^2} \approx 10^{-4}$$

The particles are injected as a power-law distribution in accord to the diffusive shock acceleration theory.

$$f(p) = f_{\text{maxwell}}(p_{inj}) \left( \frac{p}{p_{inj}} \right)^{-q} \quad q = \frac{3r}{r-1} = \frac{4}{1 - \frac{1}{M^2}}$$

$$\frac{P_{cr}}{\rho_1 u_1^2} \approx 10 - 30 \% \quad \text{and} \leq 40 \%$$

## Motivation

Shock heating is an important process in cosmic structure formation.

- Most of the baryons in the local universe have been shock-heated to  $10^5$ - $10^7$  K
- LSS most prominent in thermal X-ray from hot ICM

Astrophysical shocks are collision-less and can generate cosmic-ray (CR) distributions via 1<sup>st</sup> order Fermi acceleration mechanism (Axford et al. 1977, Krimsky 1977, Bell 1978a,b, Blandford & Ostriker 1978).

$$f(p) \propto p^{-q}; \quad q = \frac{3r}{r-1} = \frac{4}{1 - \frac{1}{M^2}} \quad \text{Test particle limit}$$

e.g. ISM hosts copious CRs

- $\gamma$ -ray emission across the disc +  $e^\pm$
- Galactic radio halo from synchrotron emission
- CRs in rough equipartition with thermal and magnetic energy

$$F_{IC}^\pm \propto N_{e^\pm} \propto \frac{S_{\text{syn}}}{\langle B \rangle^{1+\alpha}}$$

$$N_{e^\pm} \propto \frac{N_{\text{CRI}} \langle \rho \rangle_{\text{gas}}}{1 + U_B / U_{\text{CMB}}}$$

$$F_{\pi^0 \rightarrow \gamma\gamma} \propto N_{\text{CRI}} \langle \rho \rangle_{\text{gas}} \propto S_{\text{syn}} \frac{1 + U_B / U_{\text{CMB}}}{\langle B \rangle^{1+\alpha}}$$

Time-scale for energy losses

$$\tau_{Coul} \approx 5 \times 10^9 \left( \frac{\hat{p}}{0.1} \right) \left( \frac{n}{10^{-3} \text{cm}^{-3}} \right)^{-1} \text{ yr}$$

$$\tau_{p-p} \approx 5.5 \times 10^{10} \left( \frac{n}{10^{-3} \text{cm}^{-3}} \right)^{-1} \text{ yr}$$

Time-scale for diffusive escape

$$\tau_{esc} \approx 9 \times 10^9 \left( \frac{R_{cl}}{2 \text{ Mpc}} \right)^2 \left( \frac{E}{10^{15} \text{ eV}} \right)^{-1/3} \text{ yr}$$

...electrons injection

$$\frac{P_{cr}}{\rho_1 u_1^2} \leq 40 \% ; R_{e/p} = 10^{-2}$$

$$f(p) = f_{\text{maxwell}}(p_{inj}) \left( \frac{p}{p_{inj}} \right)^{-q} \quad q = \frac{3r}{r-1} = \frac{4}{1 - \frac{1}{M^2}}$$

$\gamma$ -ray background

$$\mathcal{E}^2 J(\mathcal{E}) = \frac{\mathcal{E}}{4\pi H_0} \int_0^{z_{\text{max}}} \frac{e^{-\tau_{\gamma}}}{[\Omega_m (1+z)^3 + \Omega_{\Lambda}]^{1/2}} \frac{j[\mathcal{E}(1+z), z]}{(1+z)^4} dz$$